

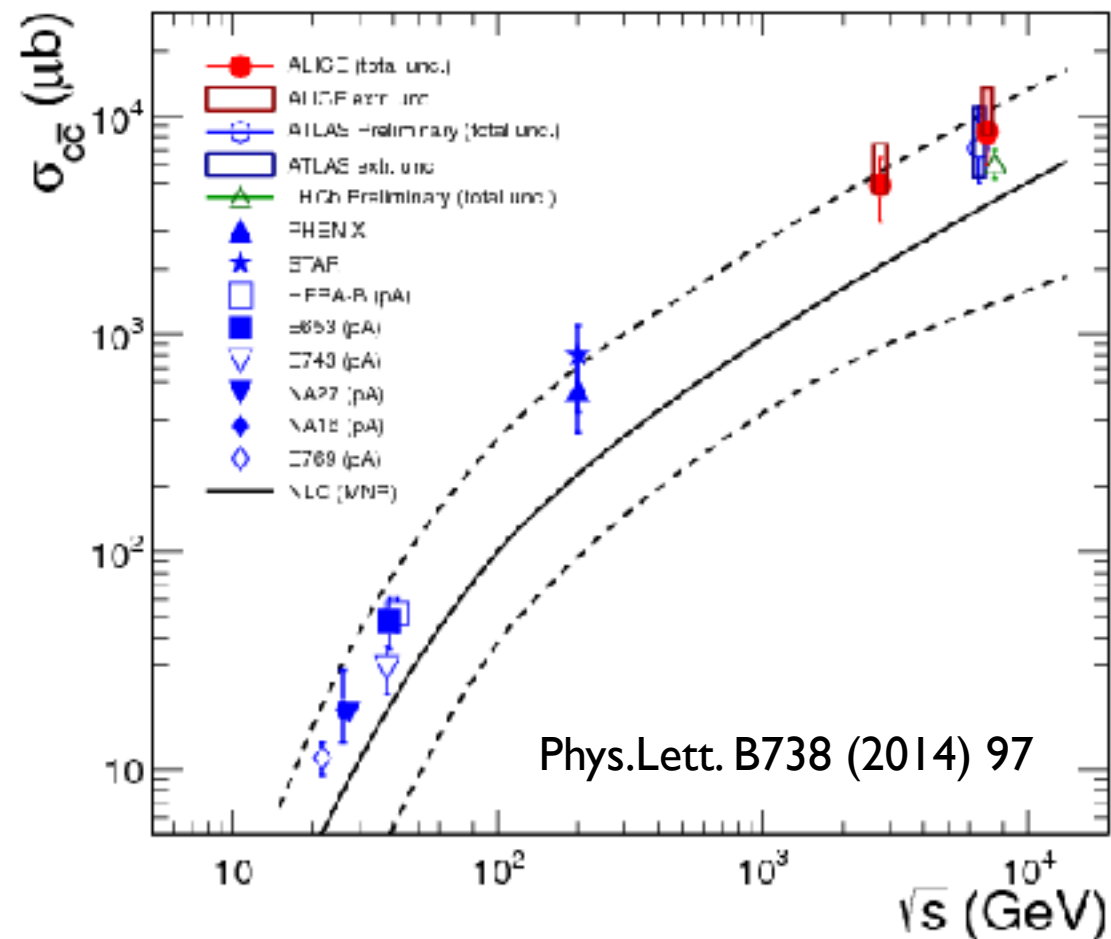
CMS HF measurements and implication to sPHENIX studies

Gian Michele Innocenti , Yen-Jie Lee
Massachusetts Institute of Technology (MIT)

***Brainstorming meeting for MVTX detector
physics deliverable***

Introduction to heavy flavour physics

Heavy quarks produced in high Q^2 processes at early stages of the collisions



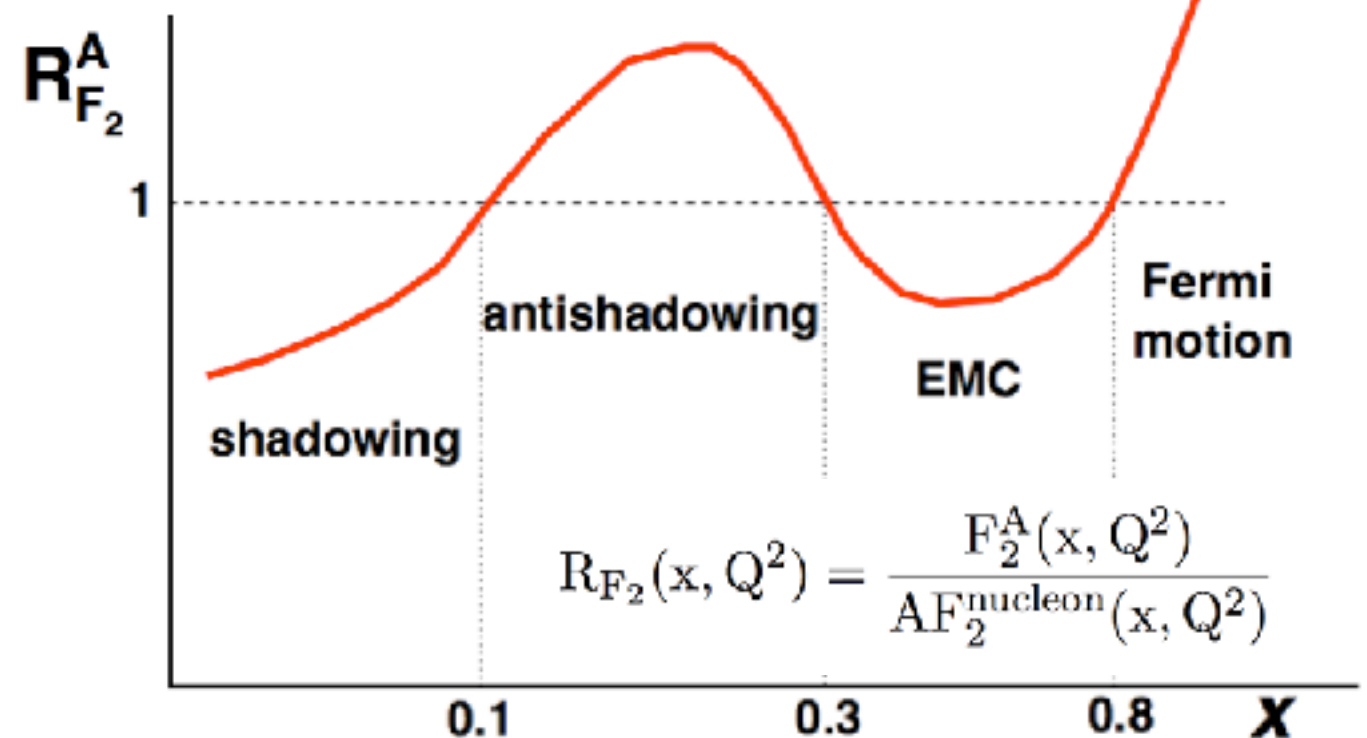
pp:

- test of pQCD calculations
- reference for pA and AA measurements
- role of MPI interactions

pPb:

- test of cold nuclear matter effects
 - PDF modifications
 - saturation
 - final state effects
- collective evolution (hydro?)

N.Armento, arXiv:hep-ph/0604108v2.



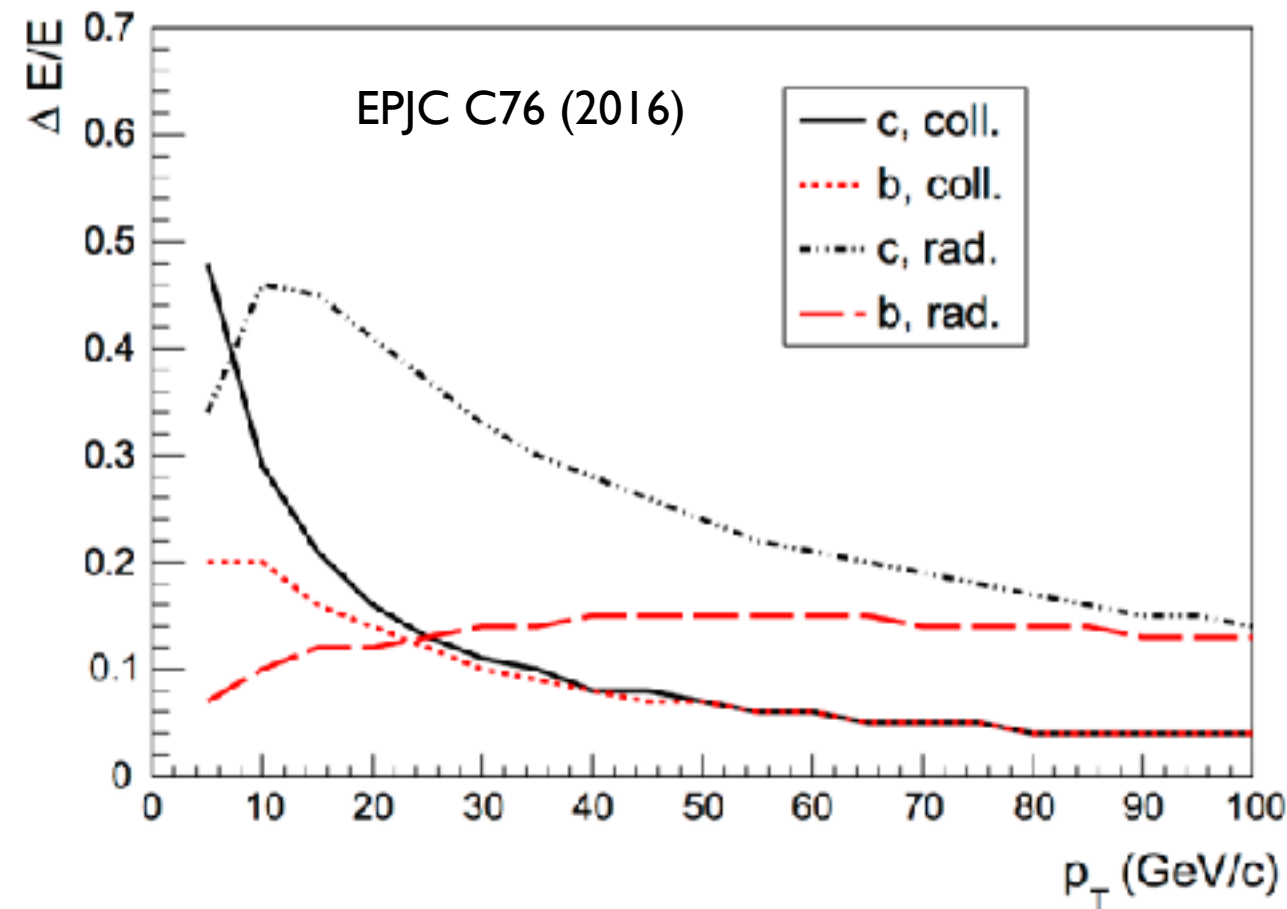
Introduction to heavy flavour physics

Heavy quark energy loss in PbPb:

- **collisional** vs **radiative** component

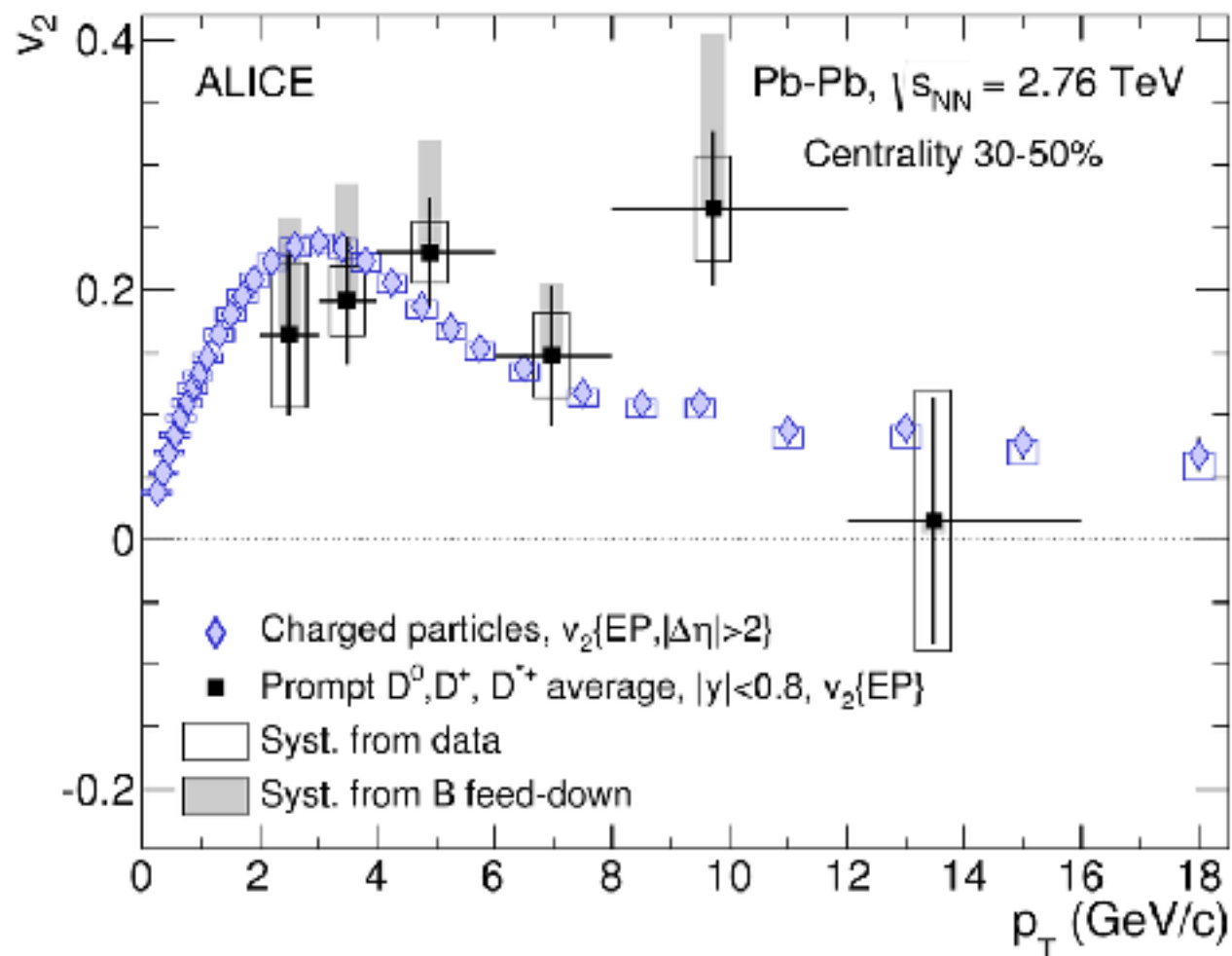
Flavour dependence energy loss:

- $\langle \Delta E \rangle \propto \alpha_s C_R q L^2$
- **Dead cone effect**: gluon radiation suppressed at small angles for massive quarks

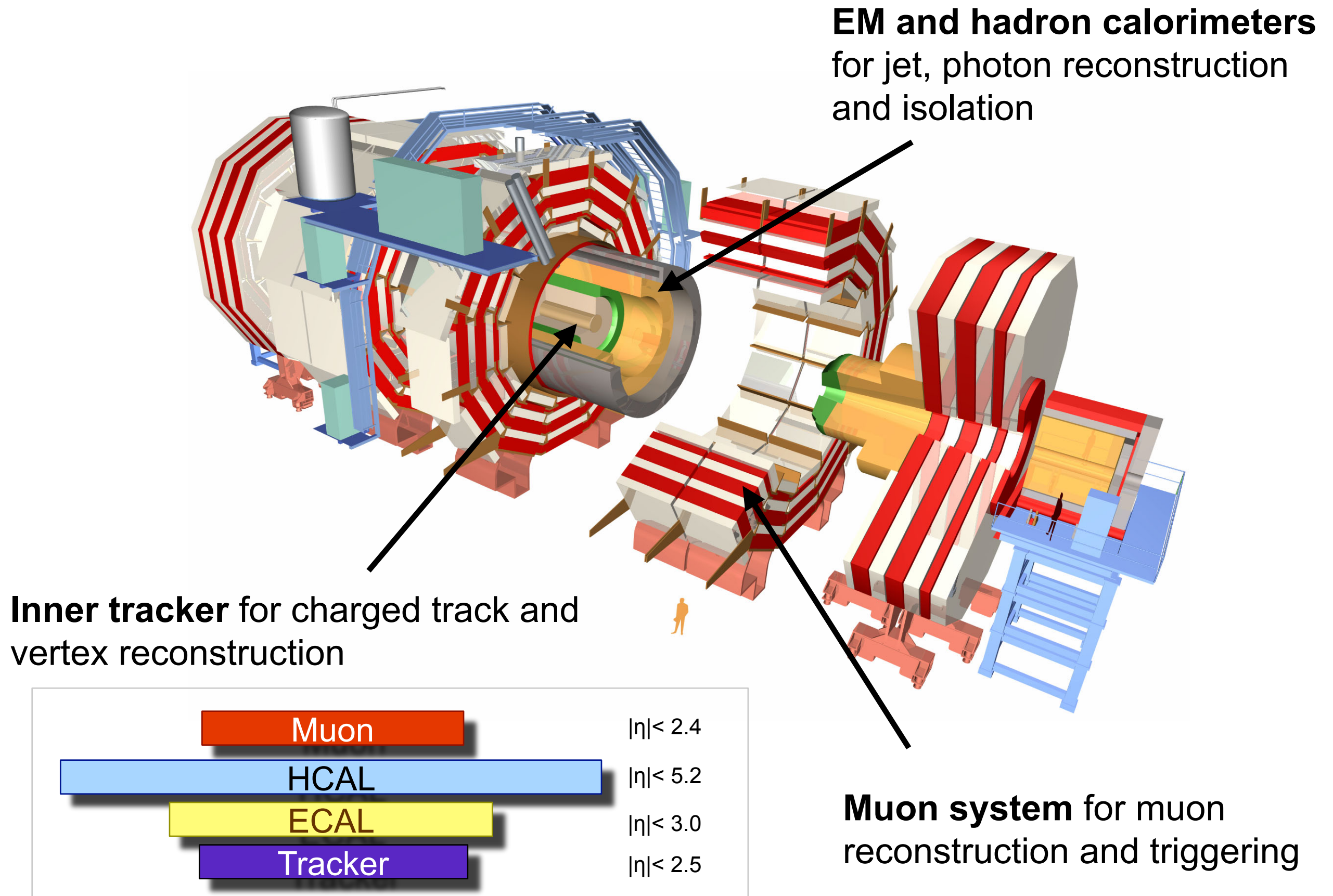


Collective behaviour:

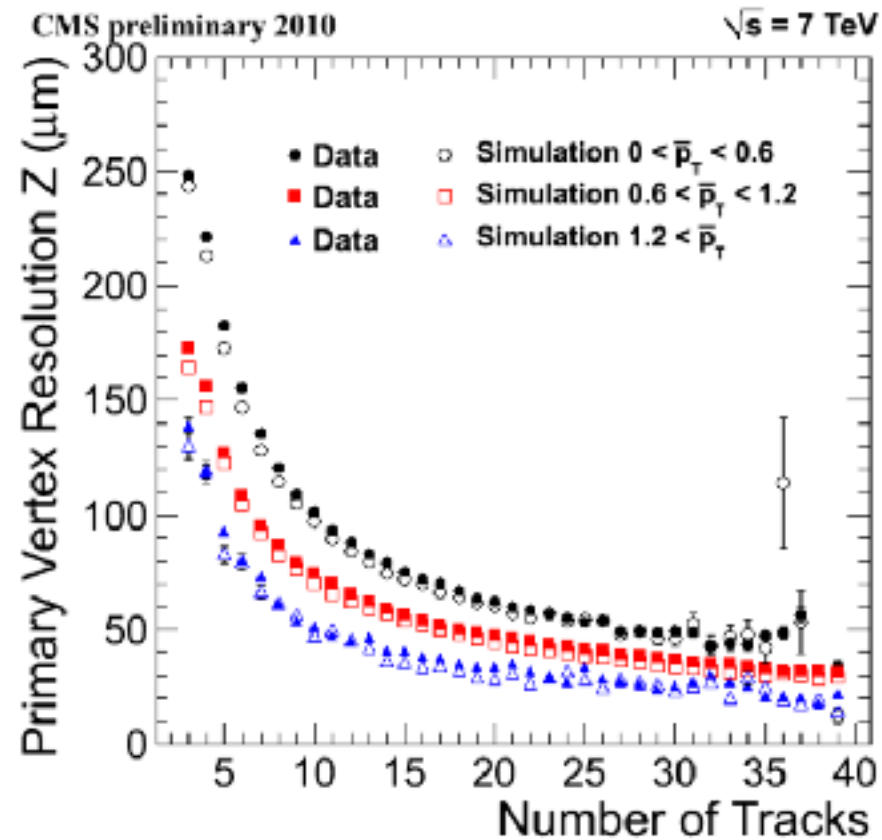
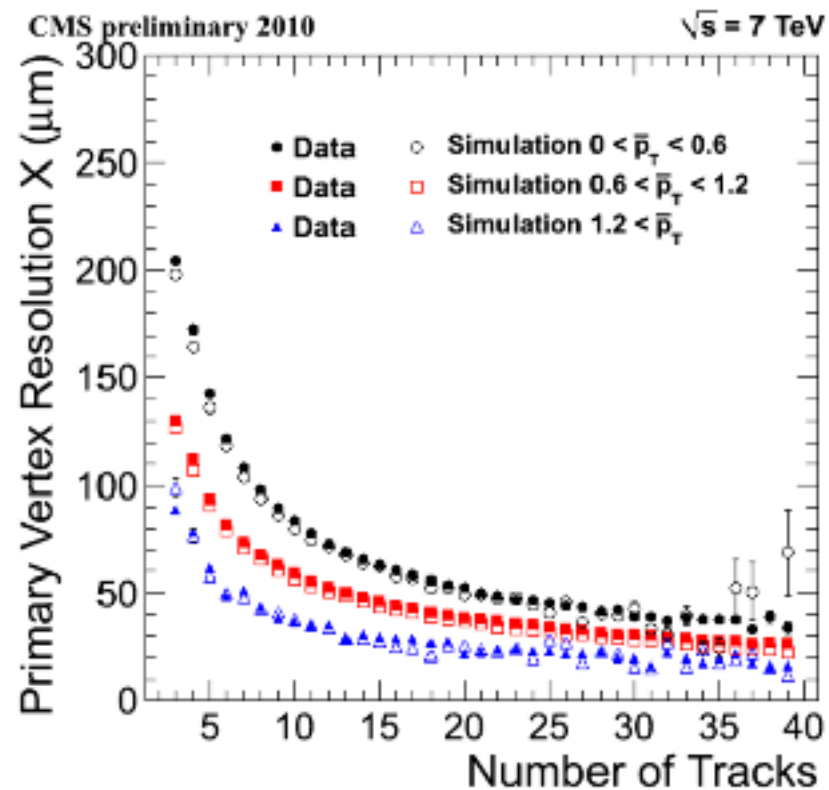
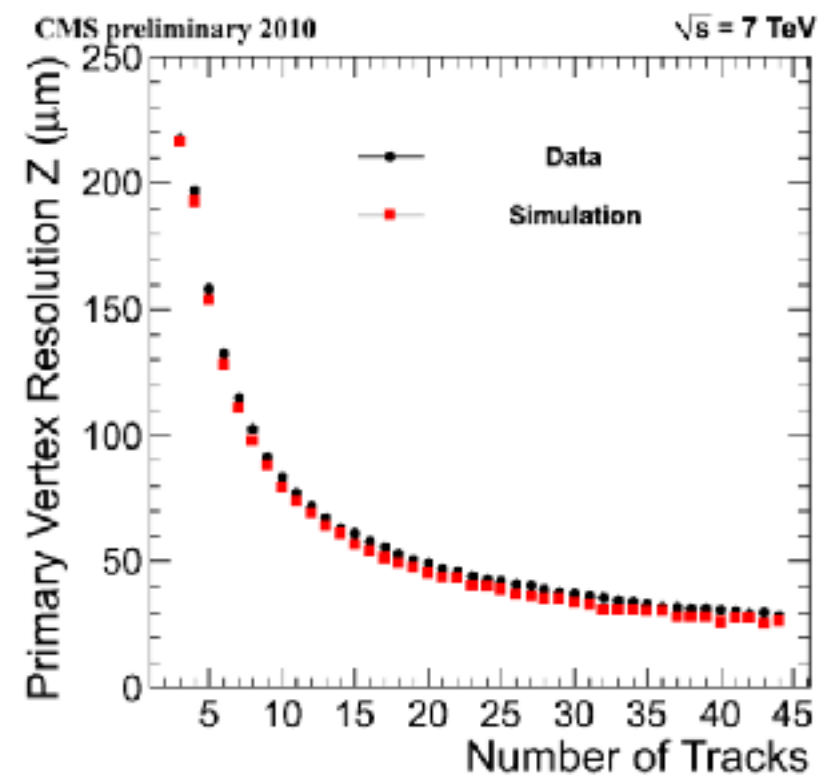
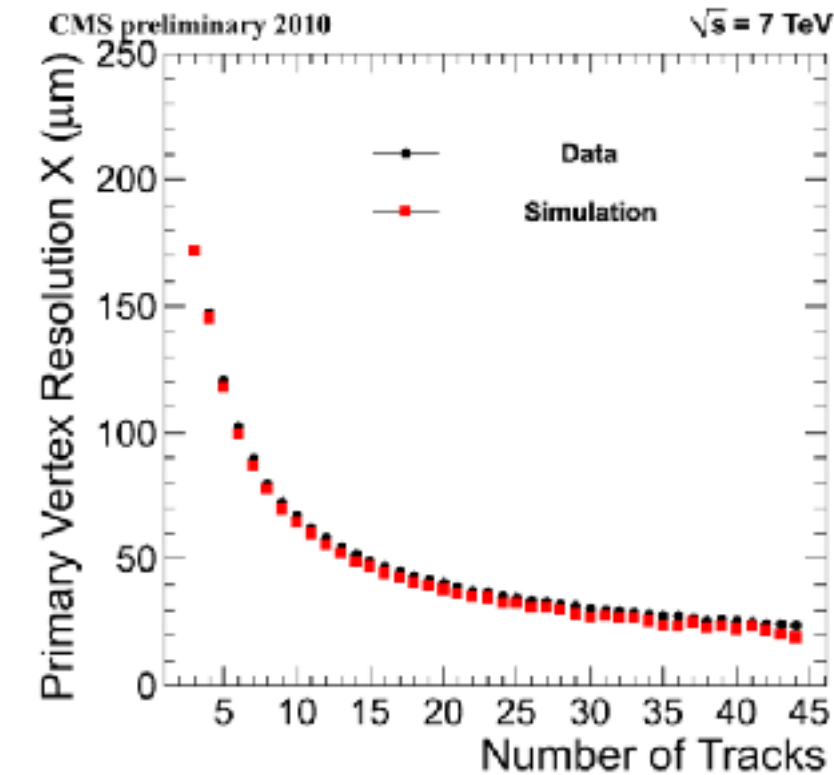
- v_n measurements to study collective behaviour of heavy quarks
- charm recombination in medium?



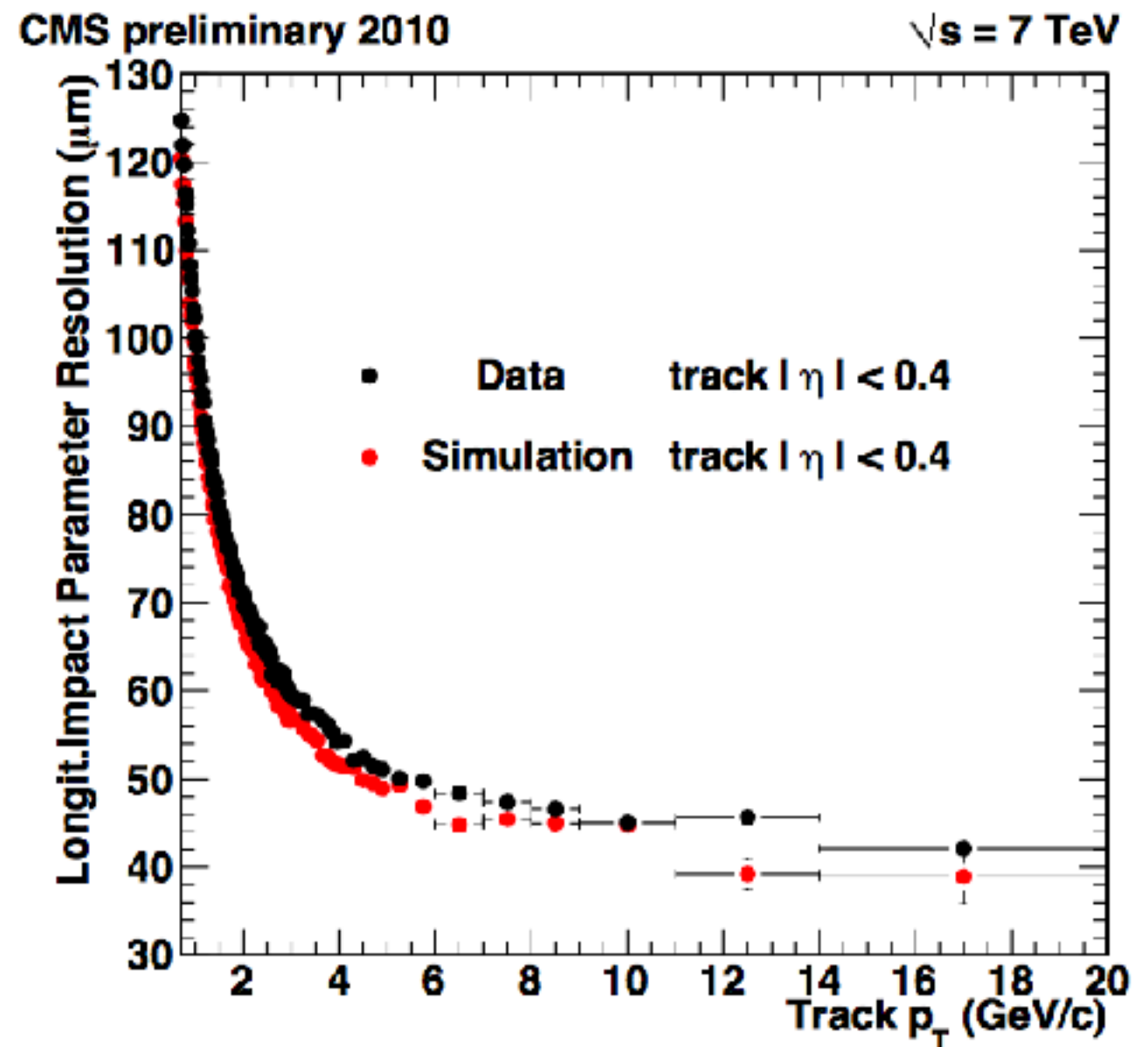
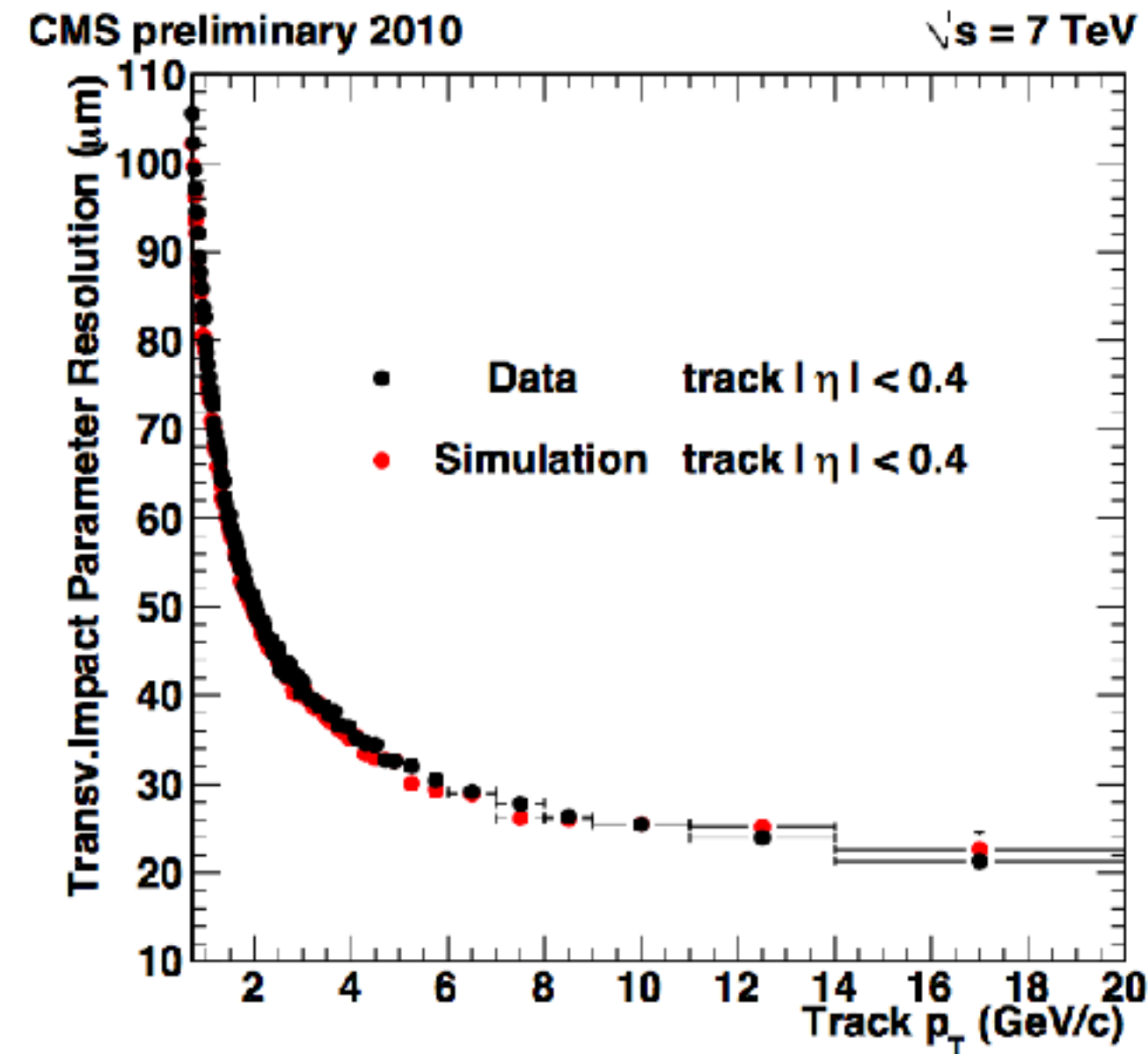
Overview of the CMS detector



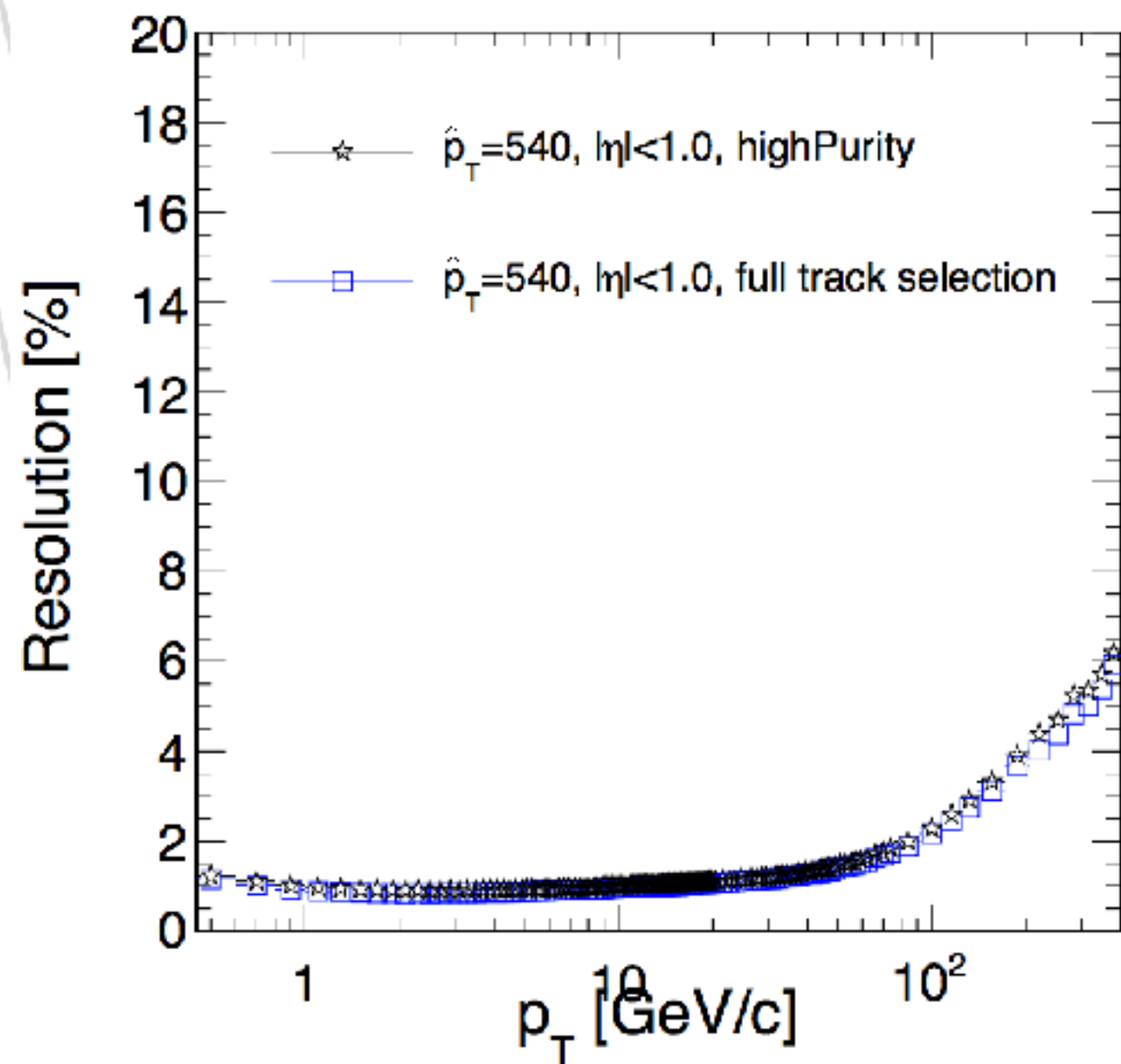
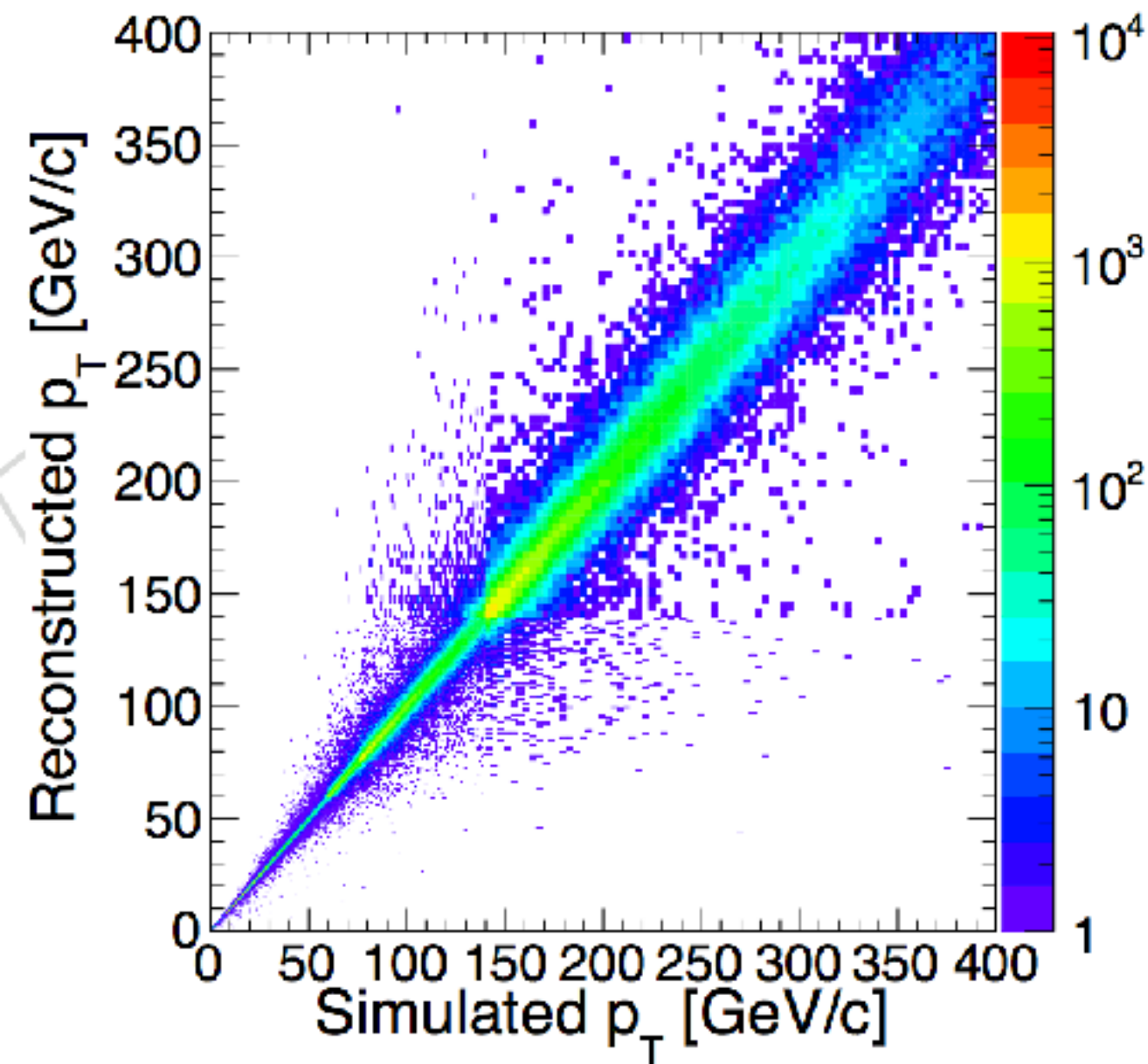
Primary vertex resolution in pp collisions



Impact parameter resolution in pp collisions



p_T resolution in PbPb collisions at 5 TeV

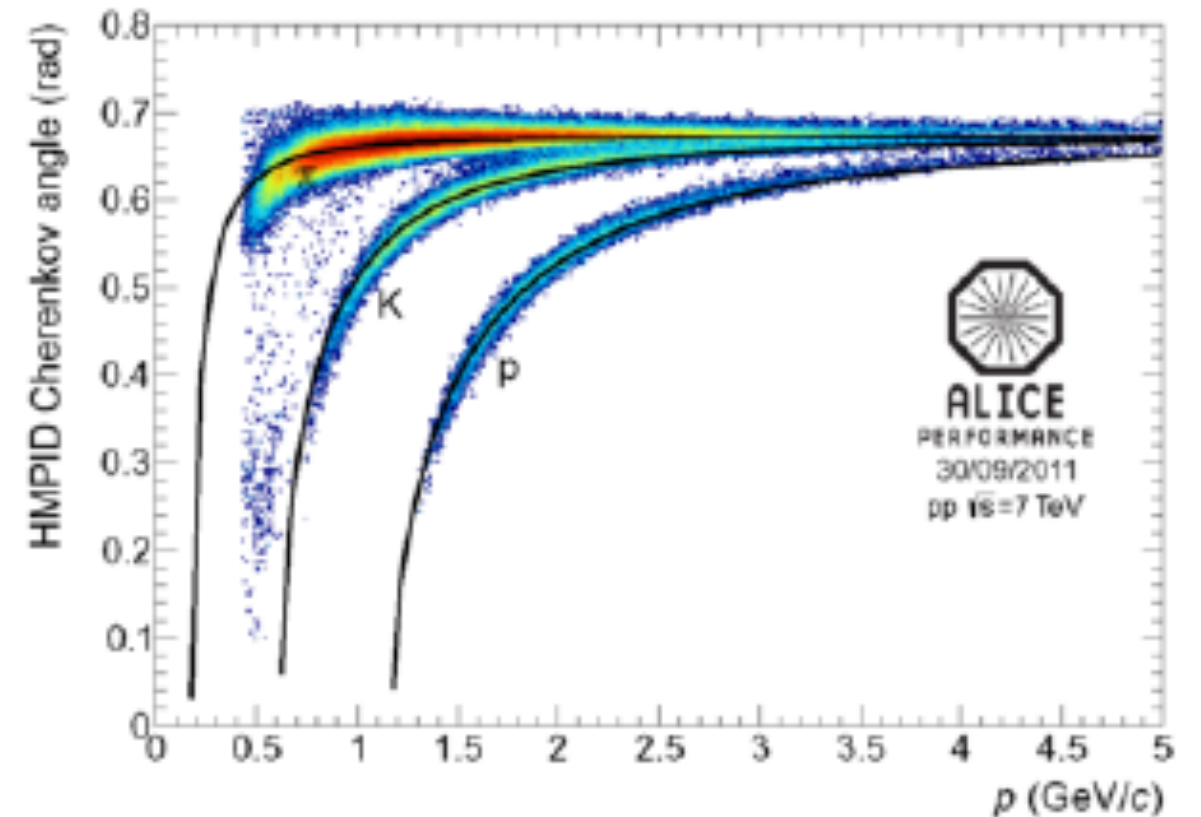
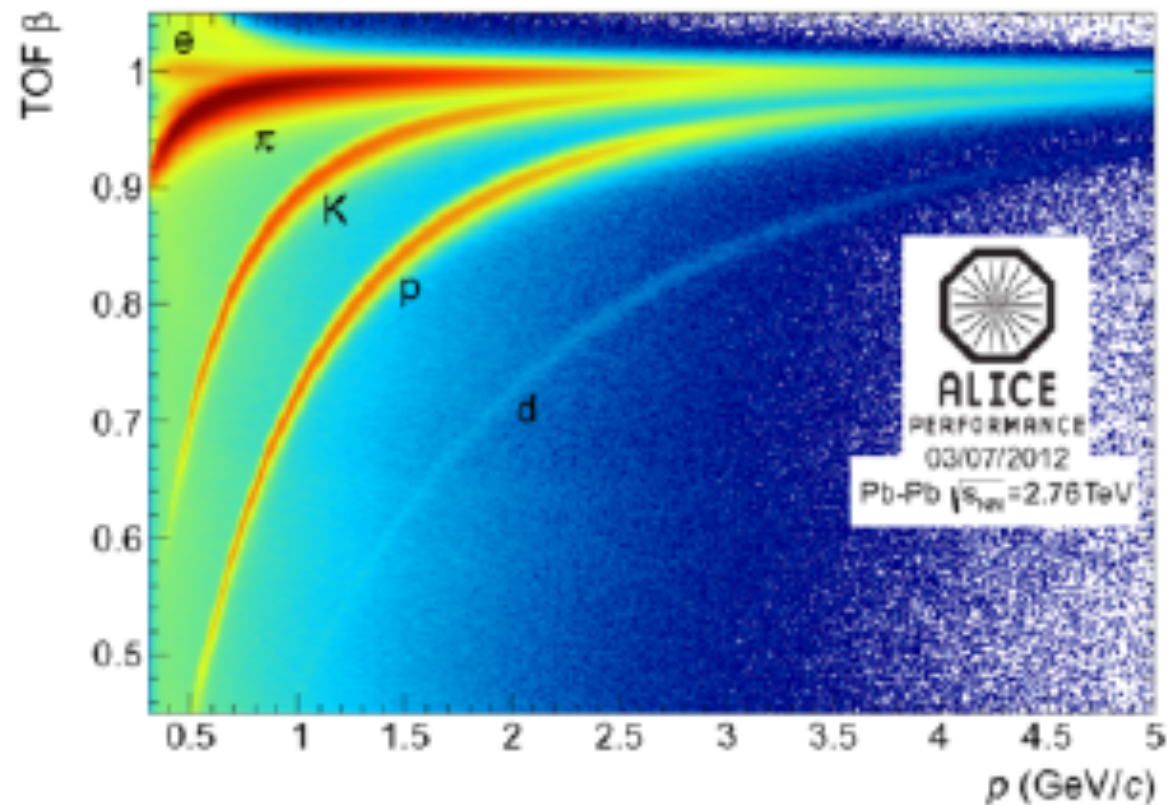
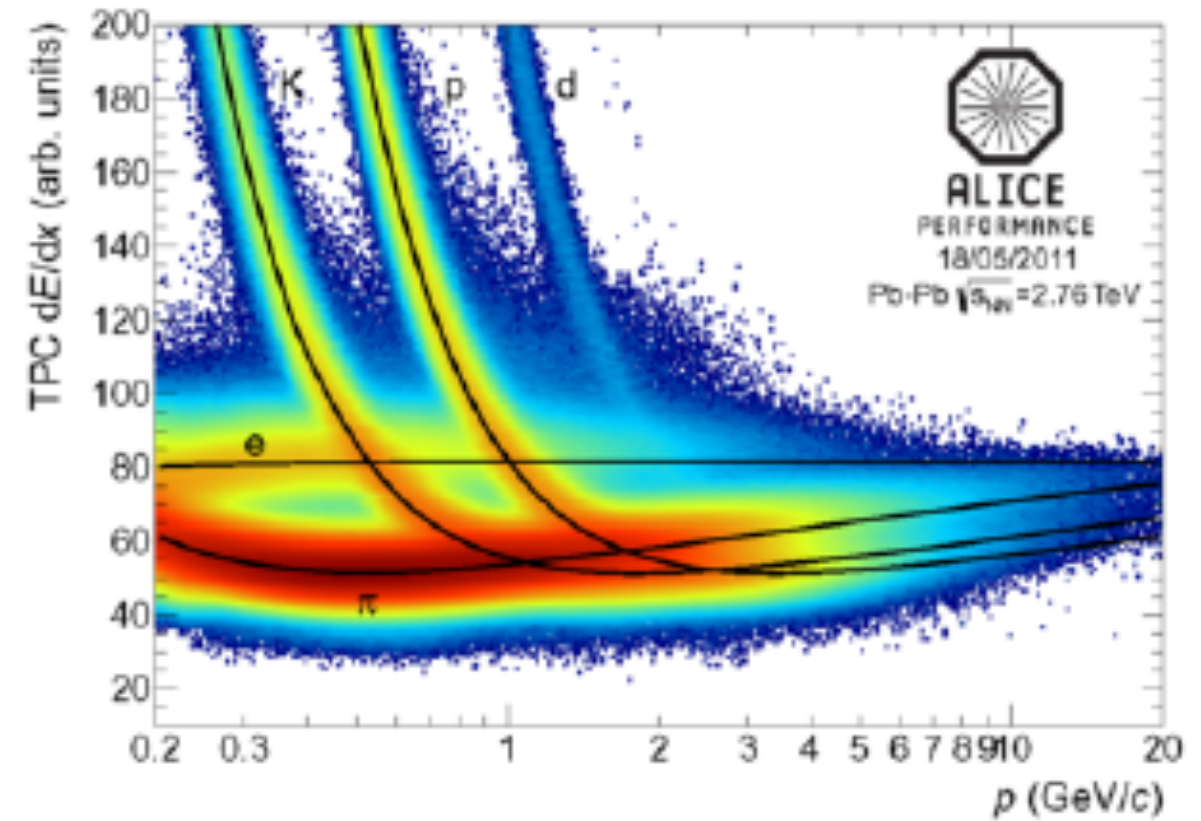
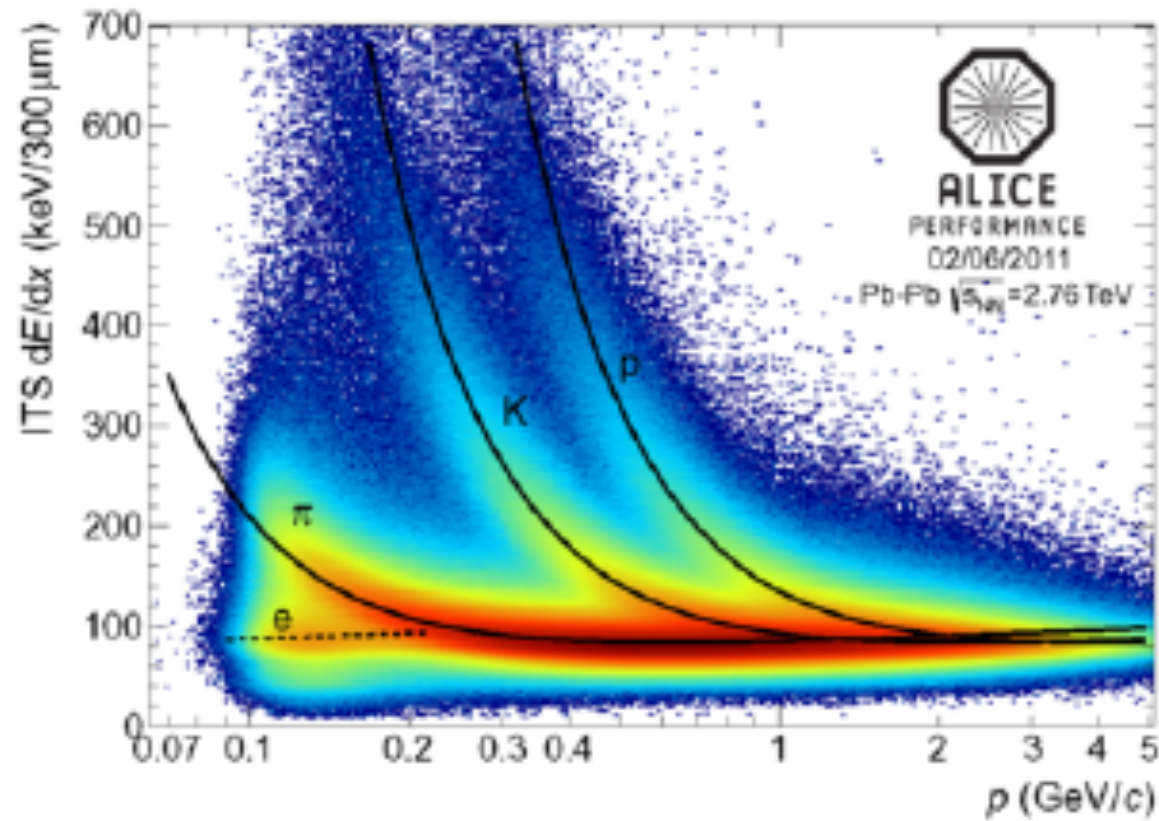


key features for a successful HF experiment

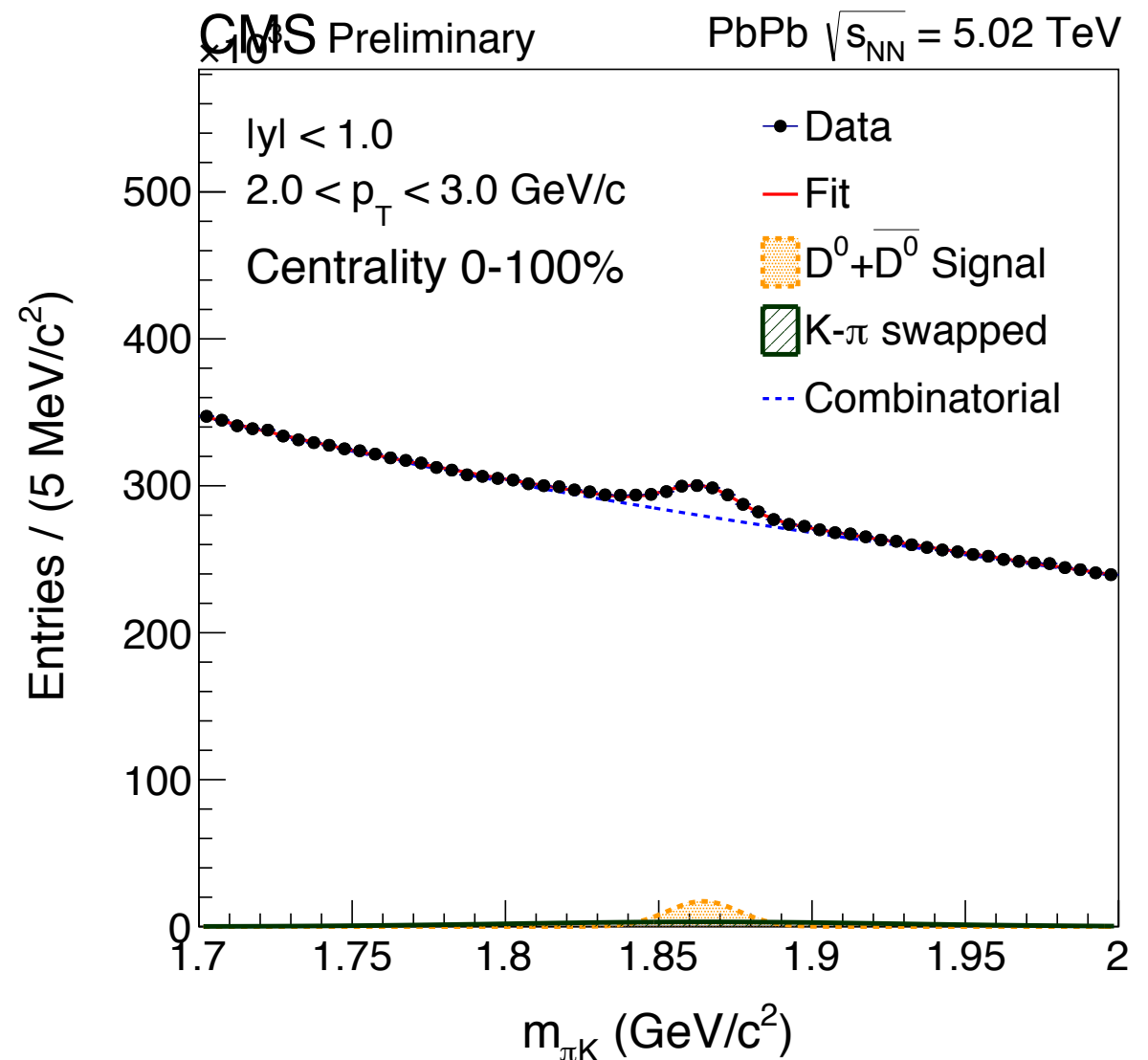
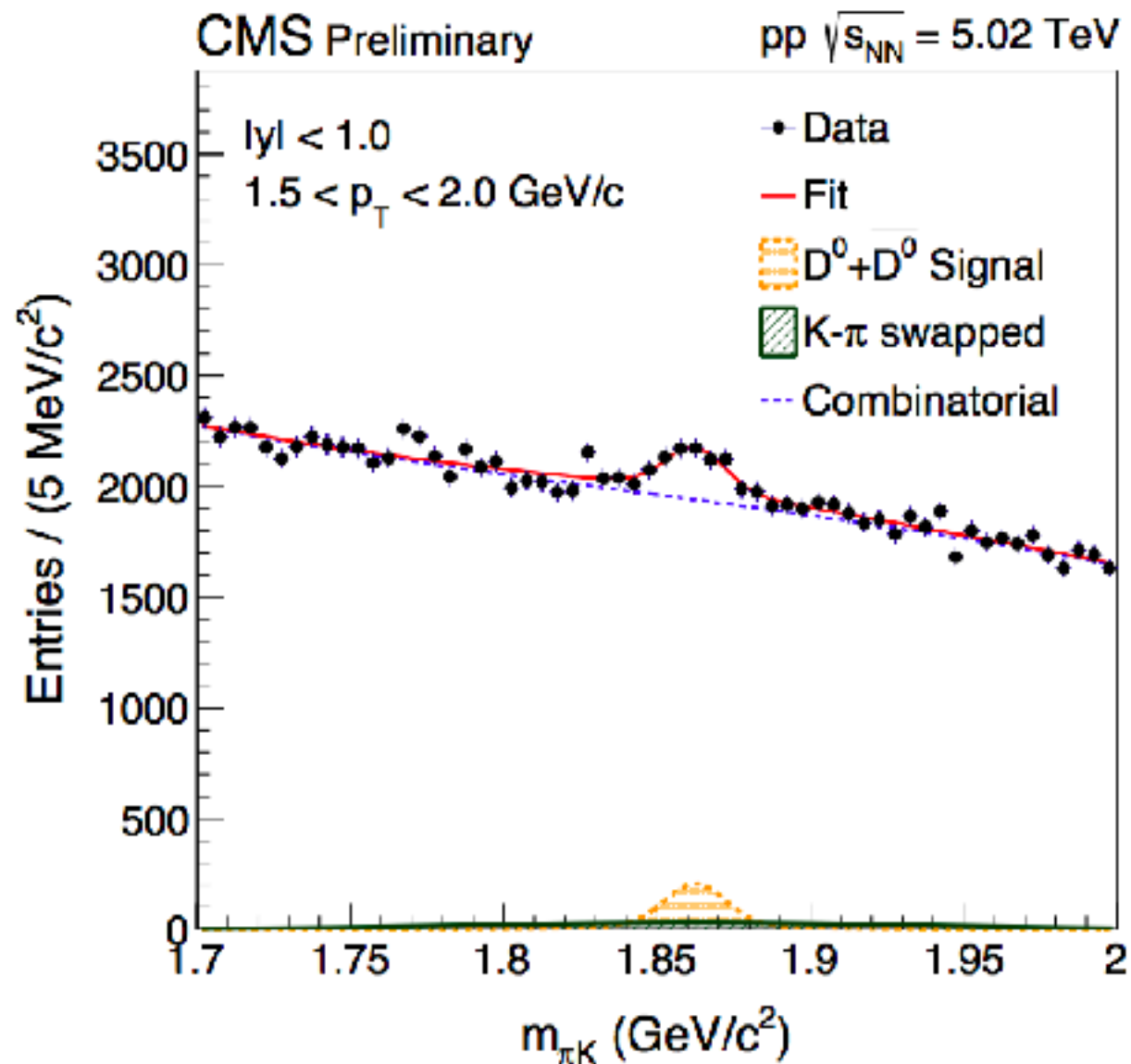
From our experience with CMS and LHC it is fundamental to have:

- primary and secondary vertex resolution
- good impact parameter resolution down to low p_T
- **Is PID fundamental for low p_T D mesons?** CMS low p_T results demonstrated that PID is important but is not strictly necessary
- **For low p_T physics, very big MB samples are clearly more important as shown by the comparison between CMS and ALICE results...**

PID ALICE performances ITS/TPC/TOF



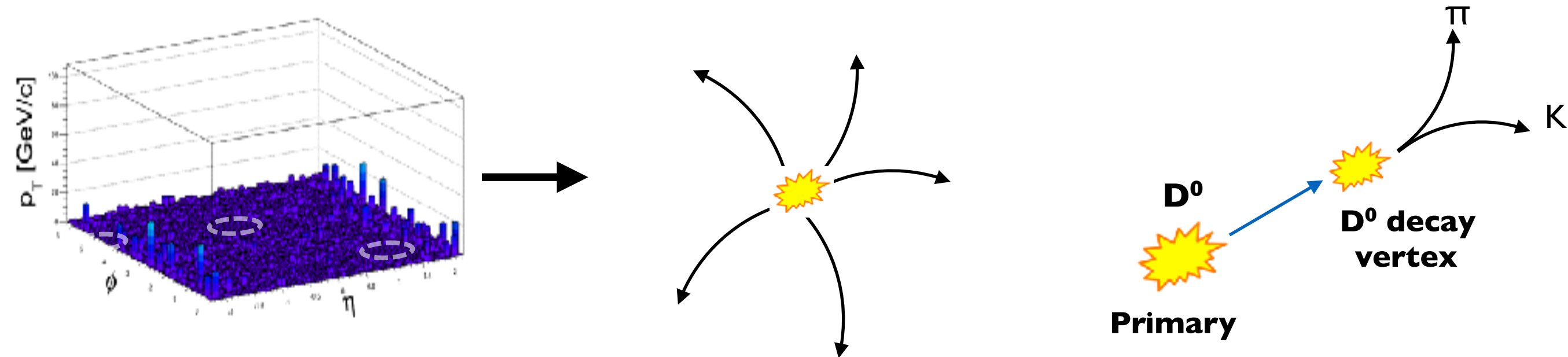
CMS low p_T results without PID



With very large samples (**~2.5 billion MB events in pp and ~300M events in 0-100% PbPb**) we can go down to very low p_T .

We can actually do better (down to 1 GeV or lower) with more cut optimisation (machine learning techniques...)

D^0 triggers at High-Level-Trigger (HLT)



Events firing hardware jet triggers (Level-1) are selected

- L1 jet algorithm with online background subtraction

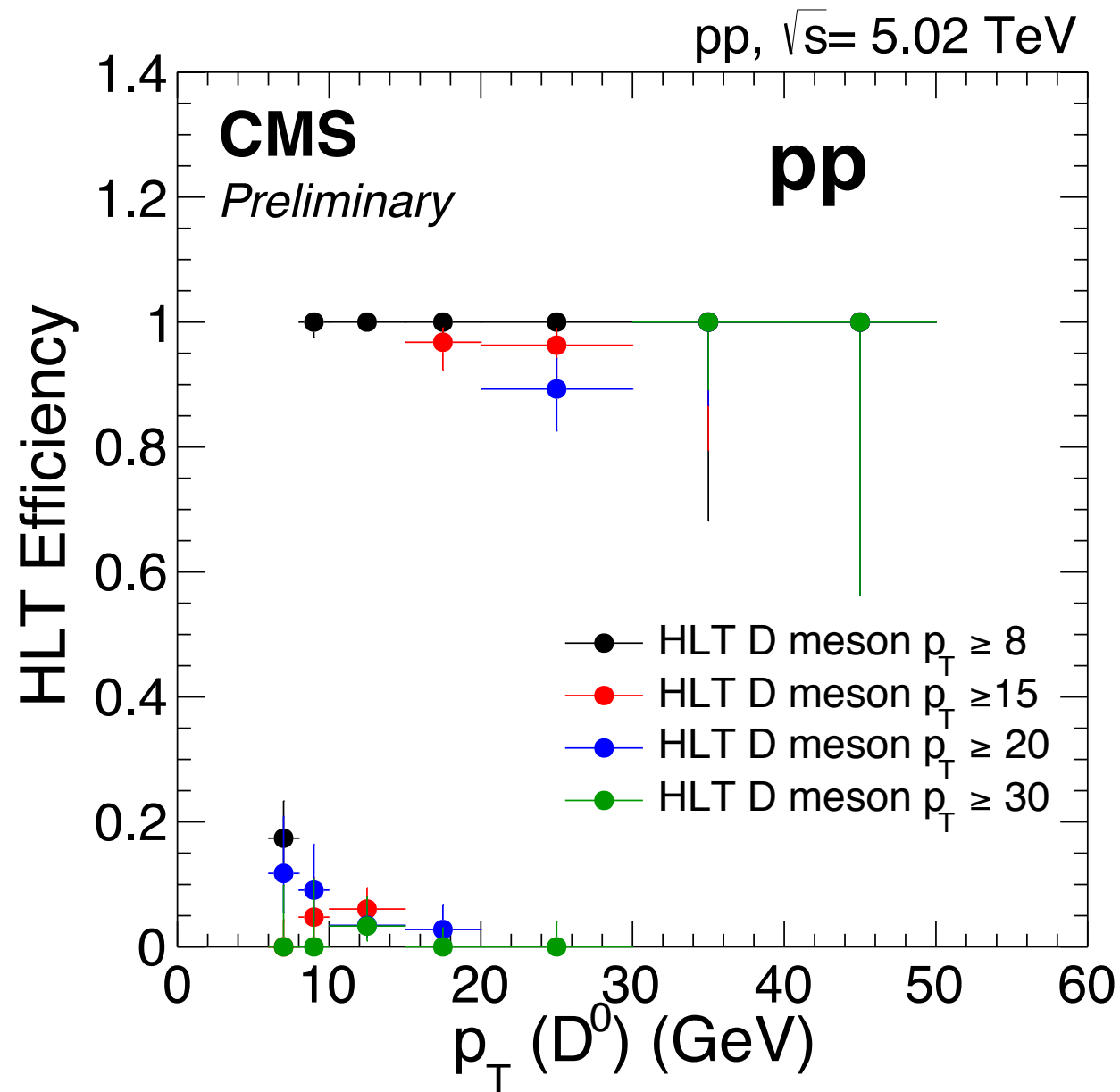
Tracks are reconstructed in software trigger system (HLT) for selected events

- Track seed p_T cut applied:
- $p_T > 2$ GeV for pp
 - $p_T > 8$ GeV for PbPb

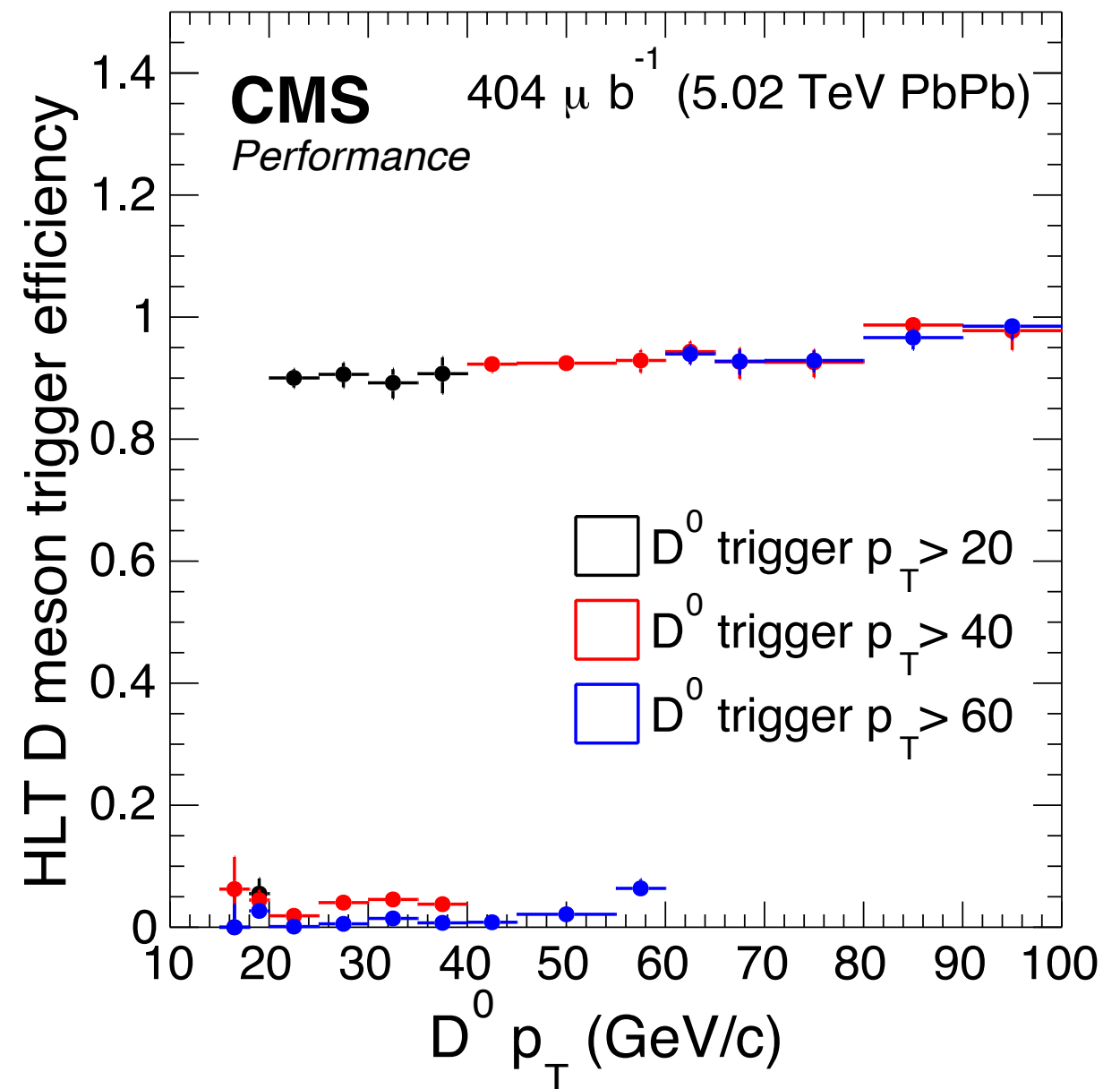
D^0 meson are reconstructed

- Online D^0 reconstruction
- loose selection to reduce the rates based on D^0 vertex displacement

Performances of D^0 triggers



→ pp efficiency reaches 100%
right above its D^0 p_T threshold

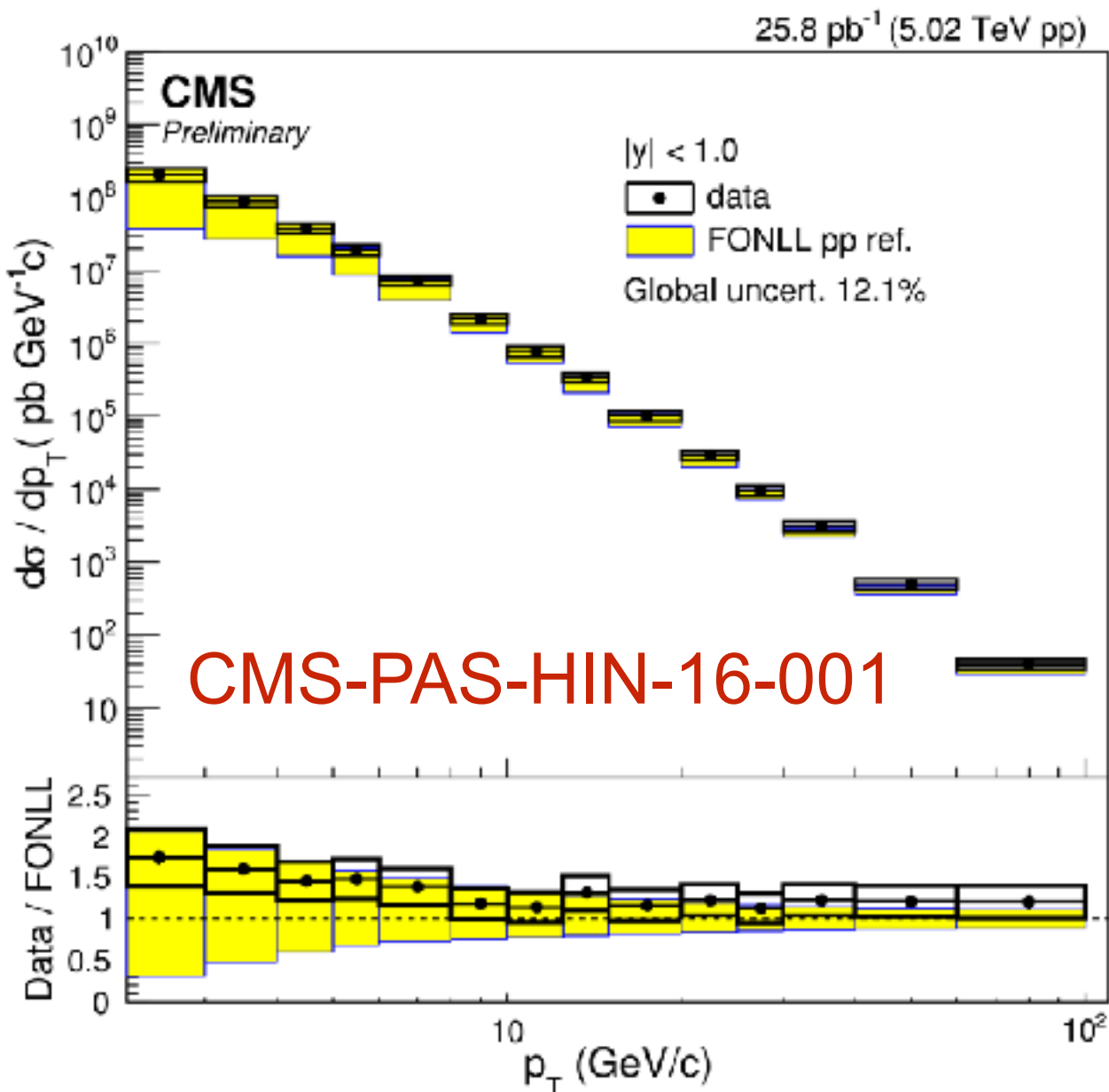


→ PbPb efficiency goes from
~90 to 100% depending on p_T

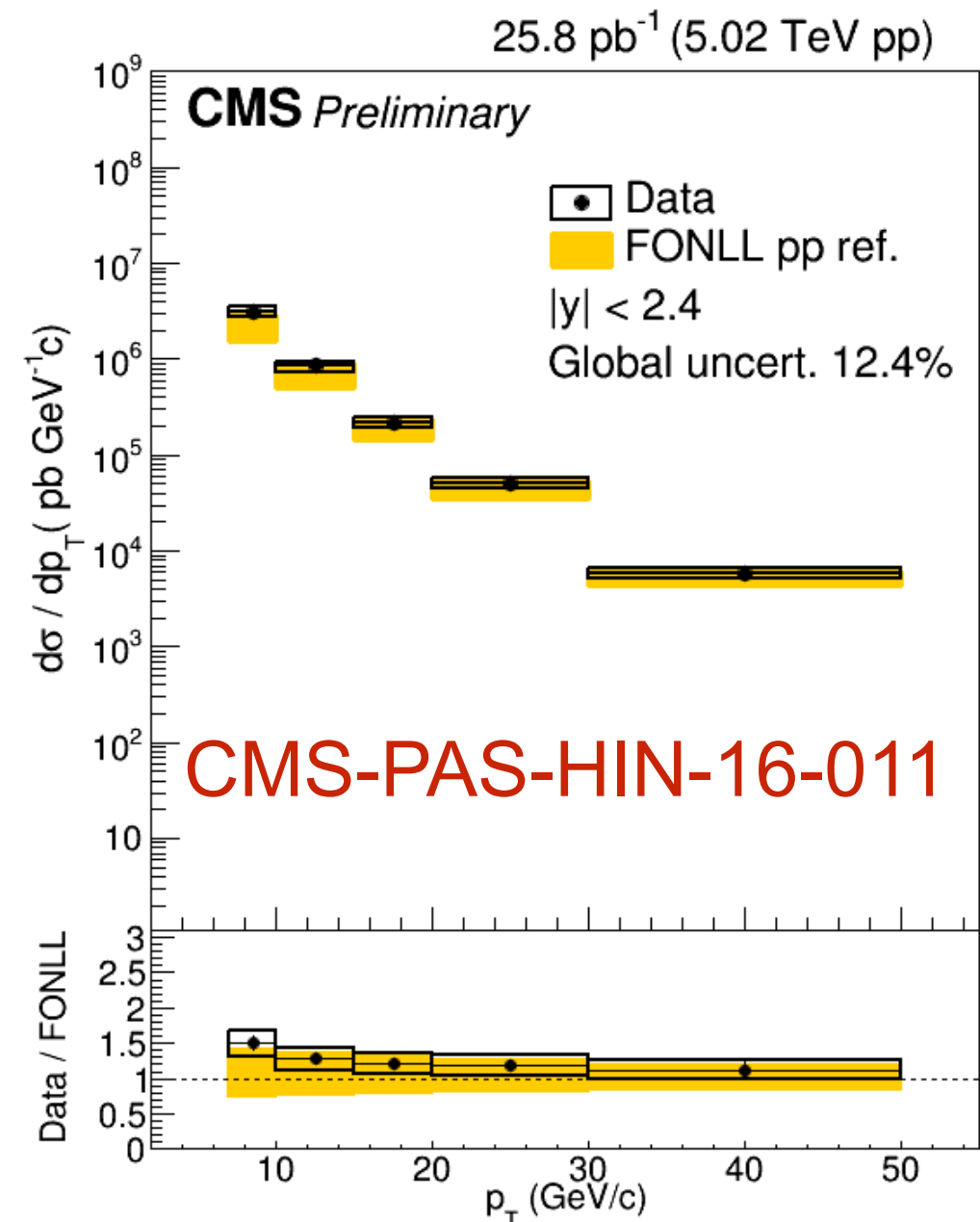
A quick summary of our results

D and B cross sections in pp collisions

D^0 at 5.02 TeV, $|y| < 1.0$



B^+ measurement at 5.02 TeV, $|y| < 2.4$



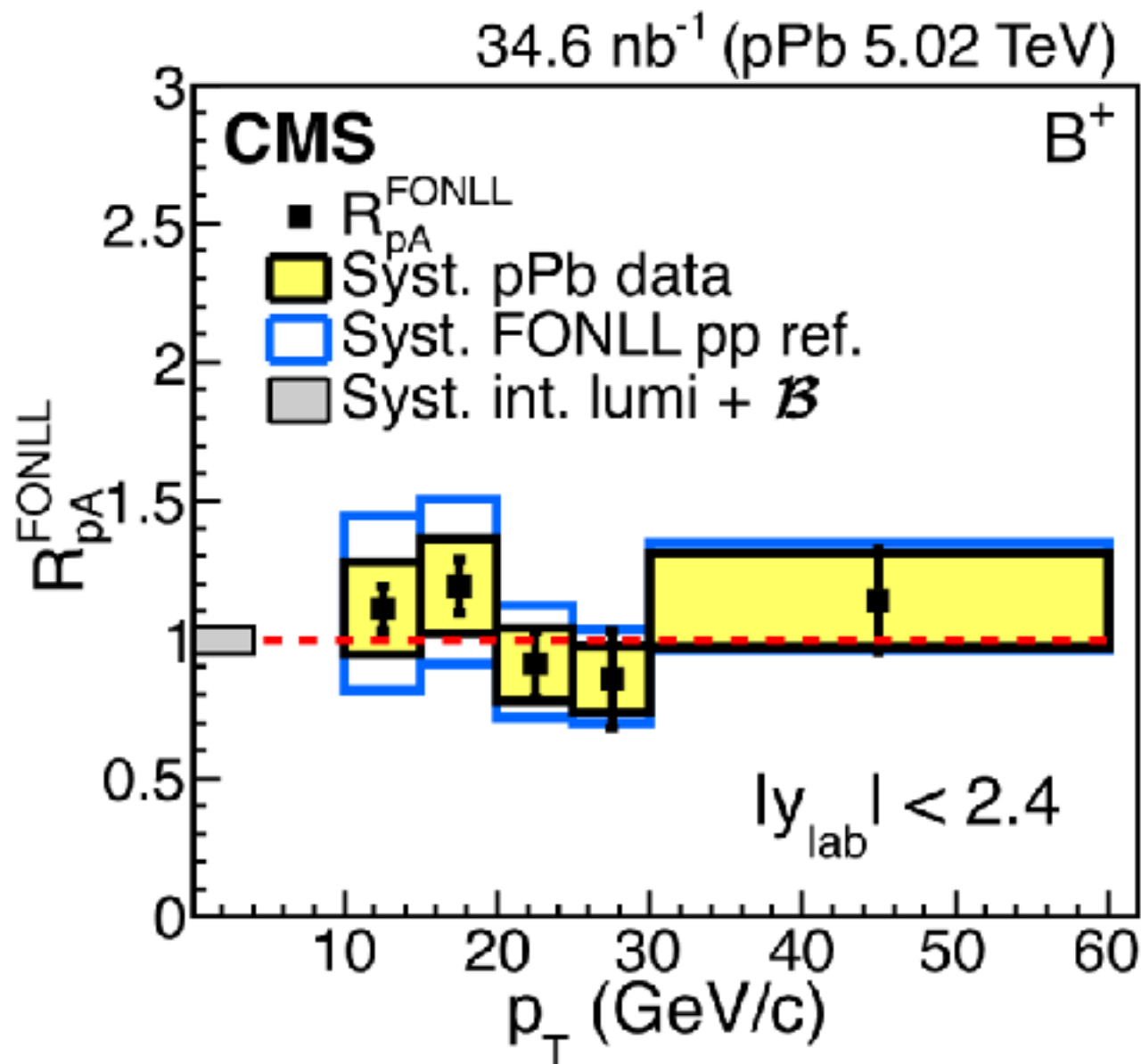
D and B meson production cross sections well described by NLO calculations:

→ D meson upper edge of FONLL calculations

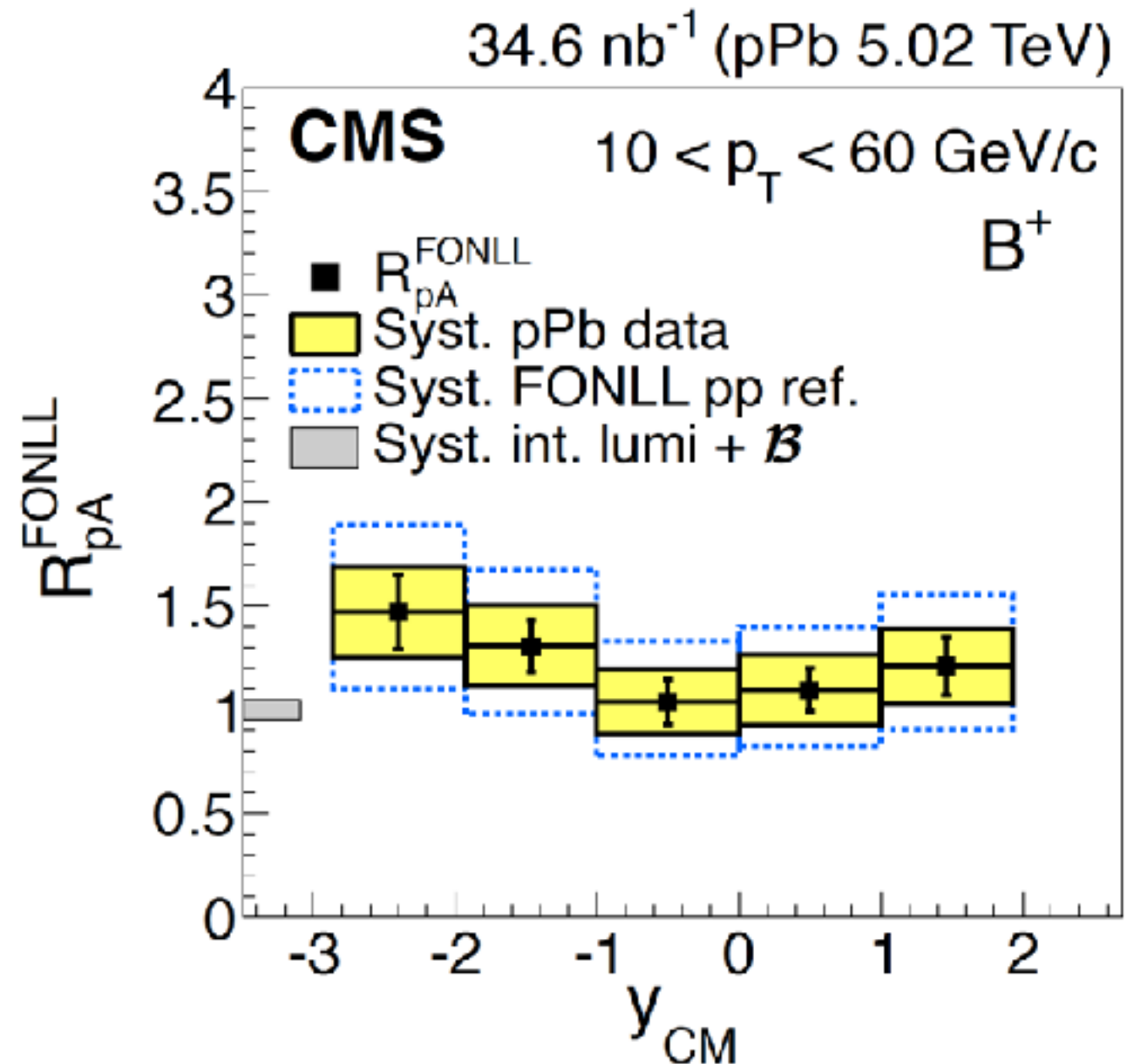
→ B meson consistent with central values of FONLL at higher p_T ,
 slightly higher for $p_T < 15$ GeV

B meson production in pPb collisions

PRL 116 (2016) 032301



FONLL R_{pA} fully compatible with unity

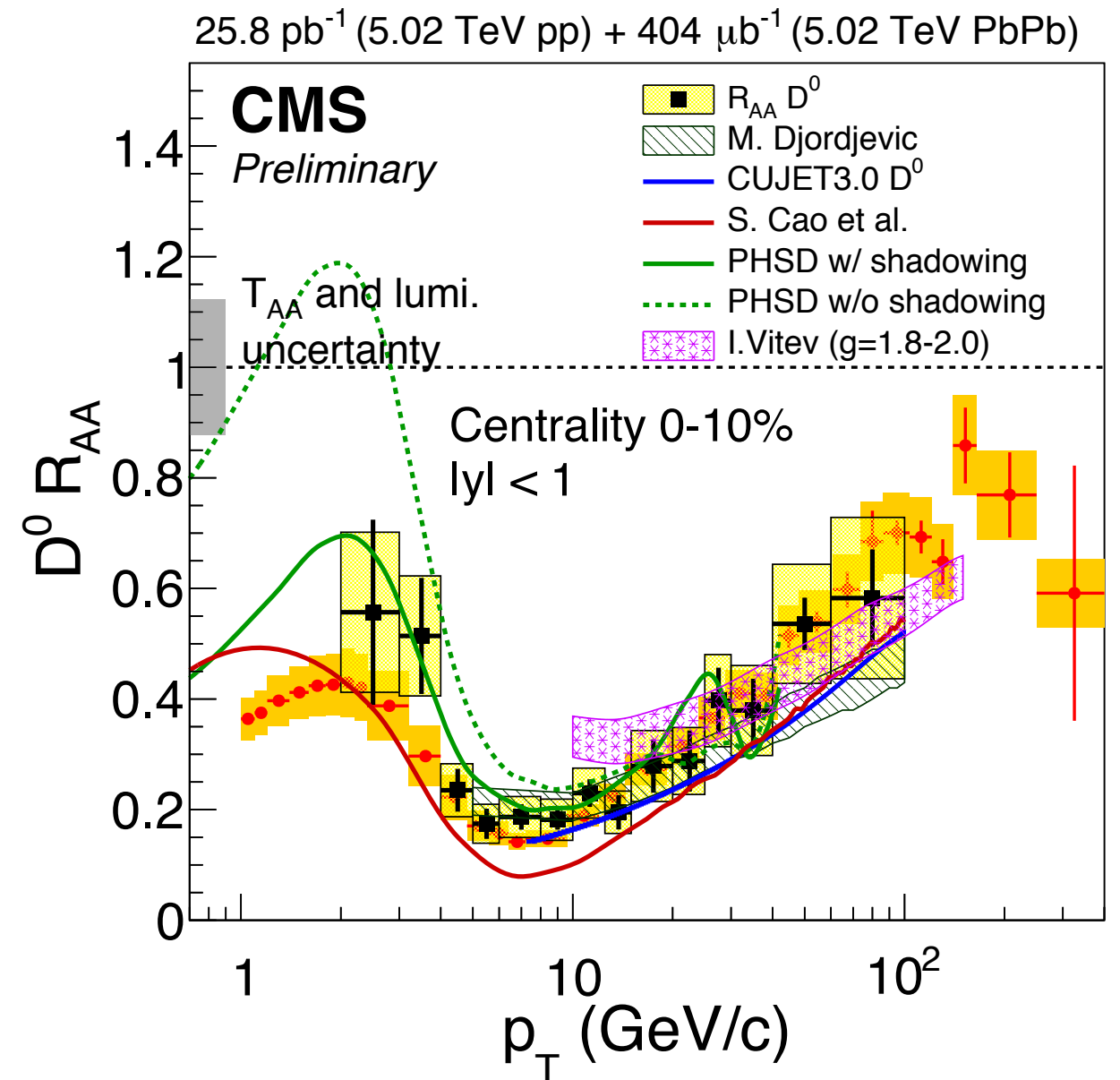
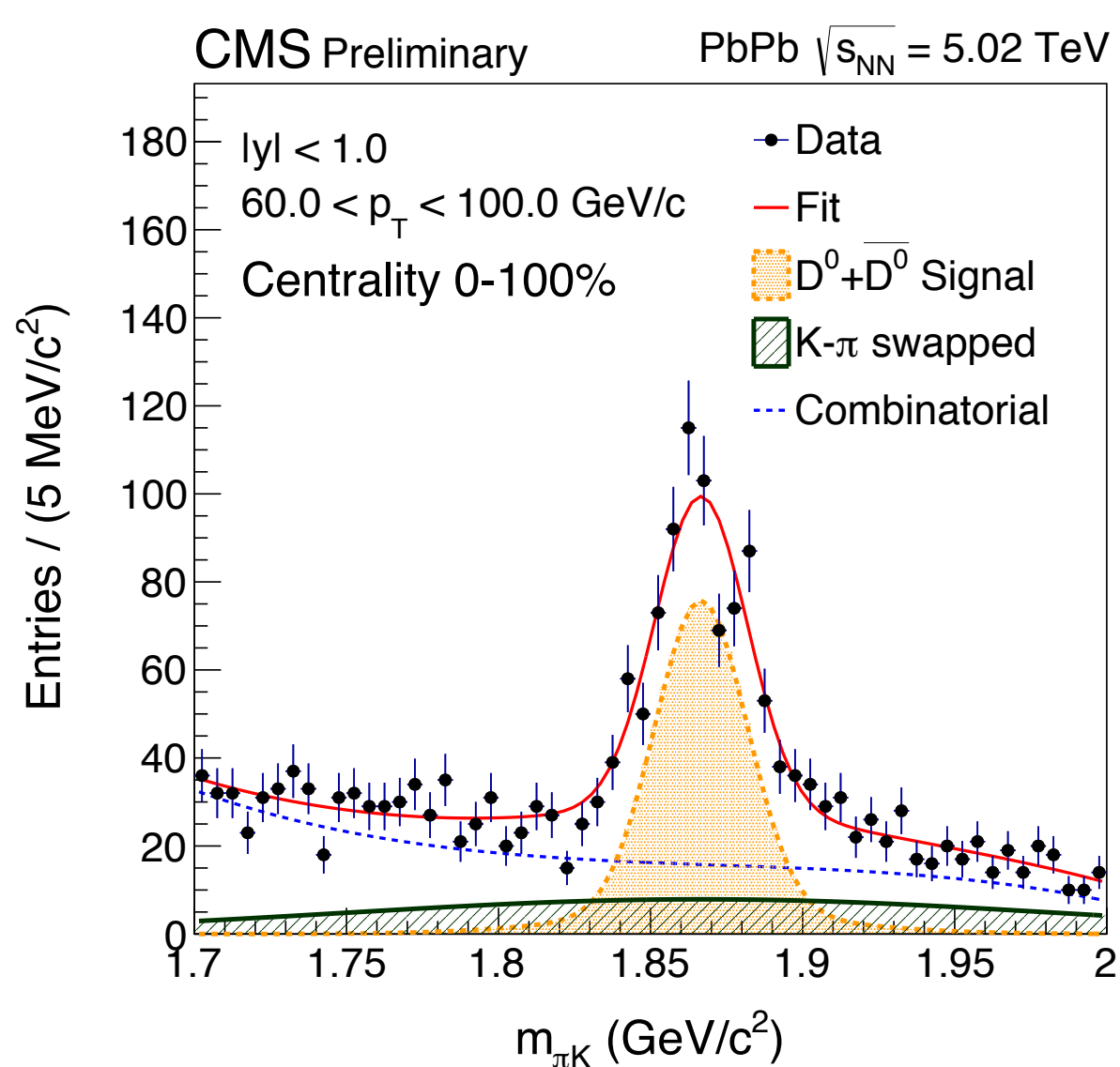


No sizeable modification as a function of rapidity

→ **D** R_{pA} and HF electron studies can give also complementary information!

D⁰ meson R_{AA} at 5.02 TeV

D⁰ R_{AA} measured from 2 to 100 GeV/c at central rapidity in 0-100%

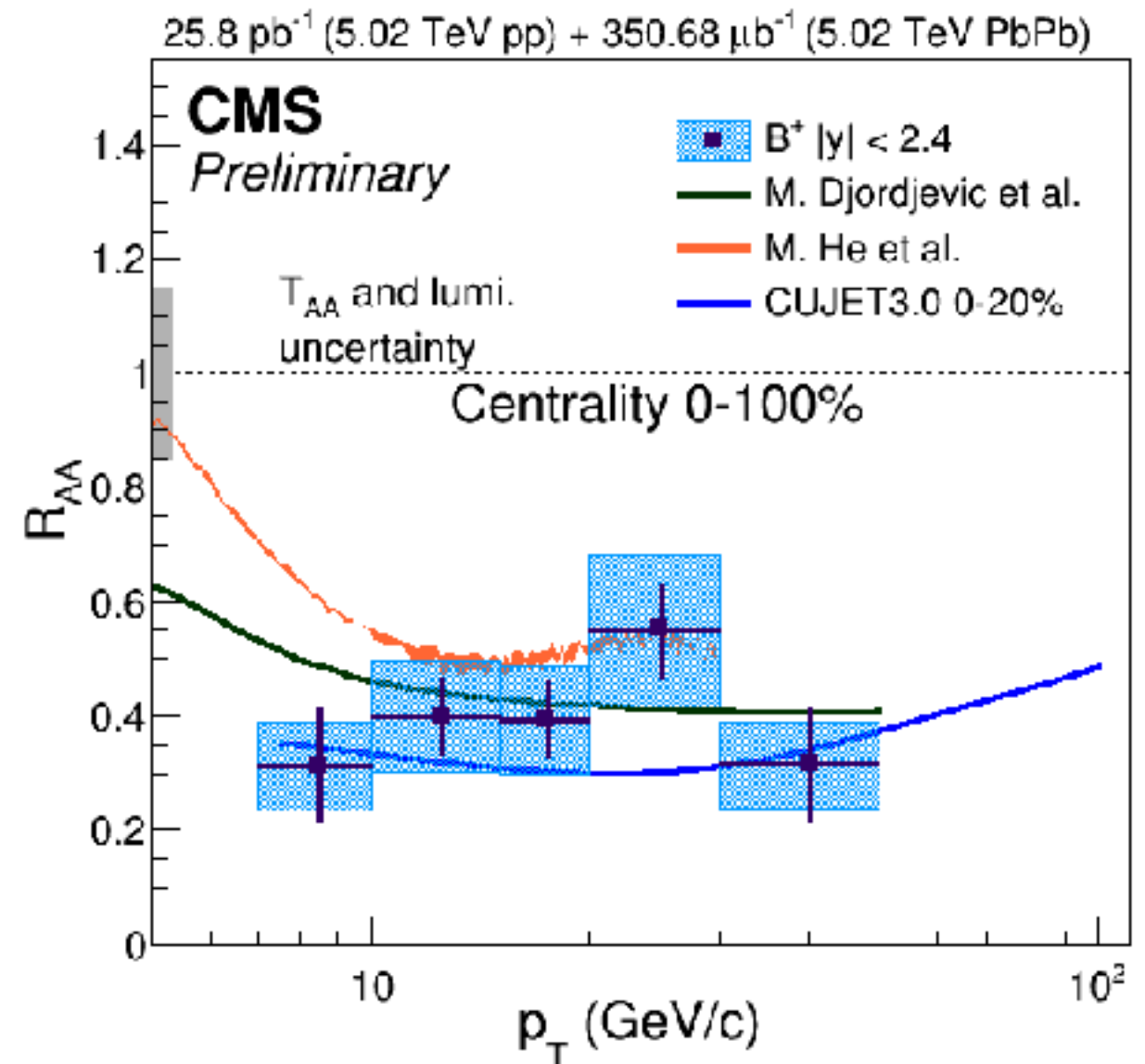
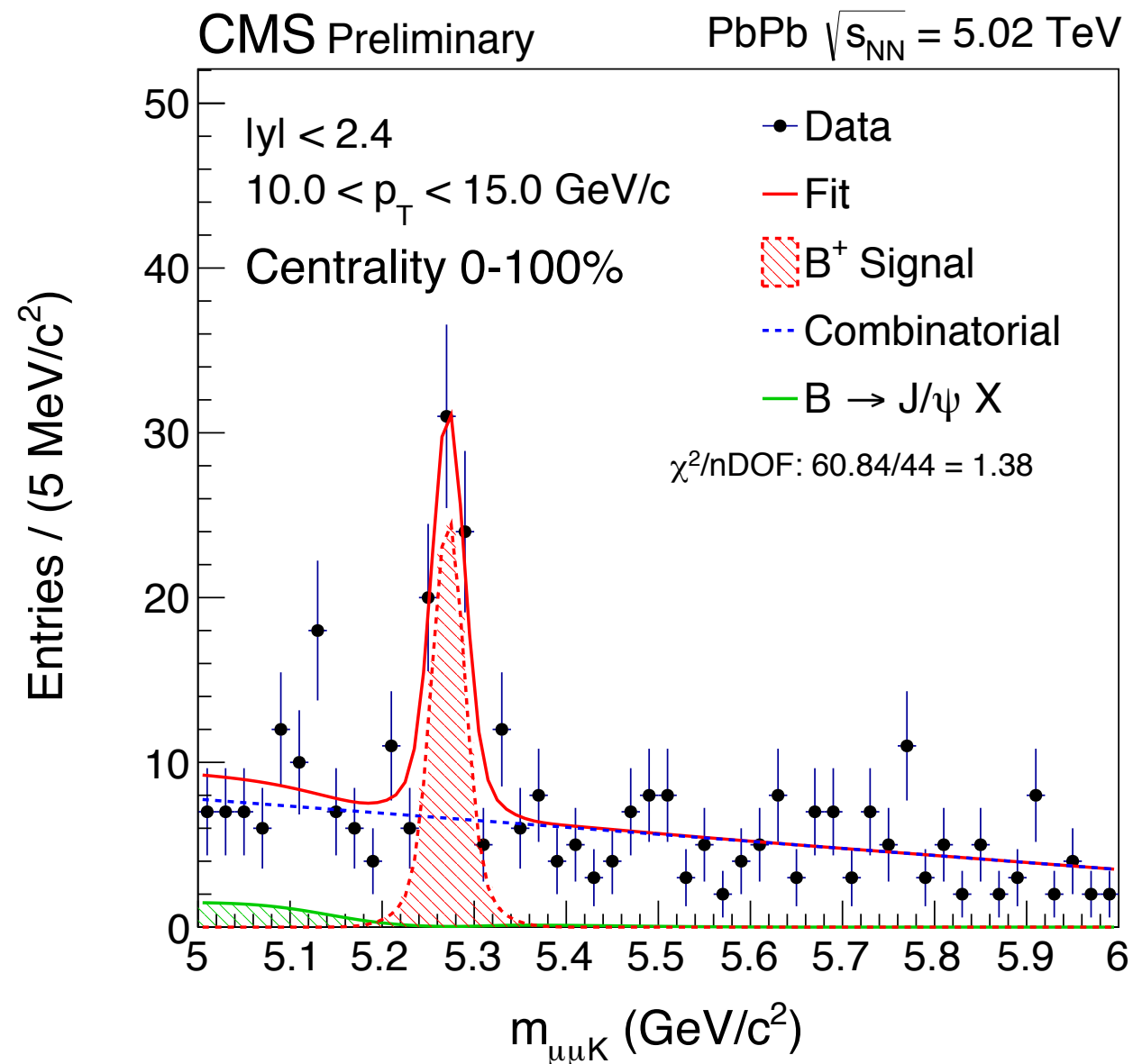


- ~300 MB events recorded for low p_T analysis
- Dedicated triggers designed to perform full track reconstruction at HLT and to reconstruct and identify D⁰ meson events at HLT

Exclusive B^+ meson measurement in PbPb

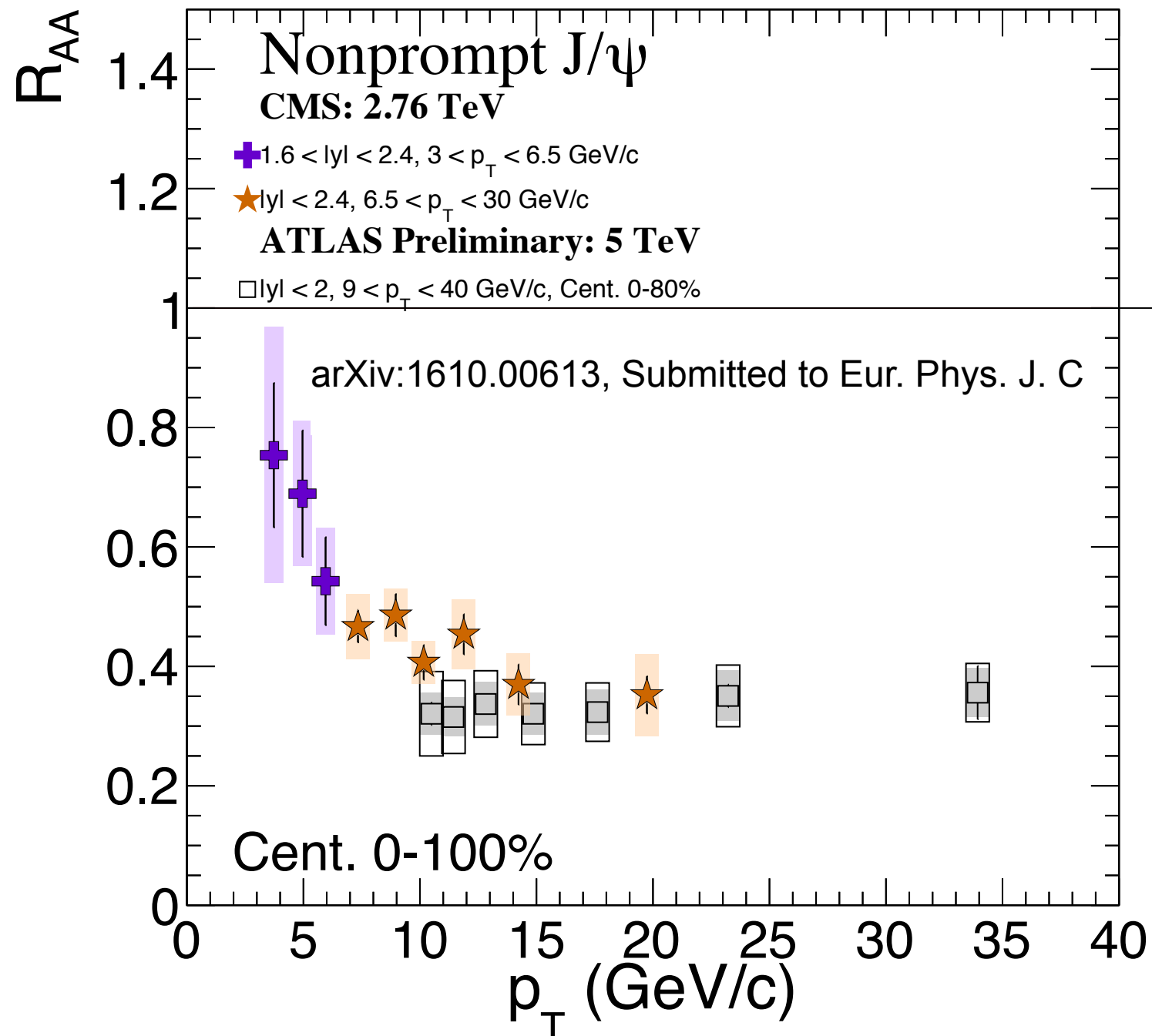
CMS B^+ production in PbPb at central rapidity $|y| < 2.4$

CMS-PAS-HIN-16-011



Strong suppression ($R_{AA} \sim 0.4$) observed in 0-100% PbPb collision for $p_T > 7$ GeV/c
Well described by theoretical calculations that include radiative energy loss

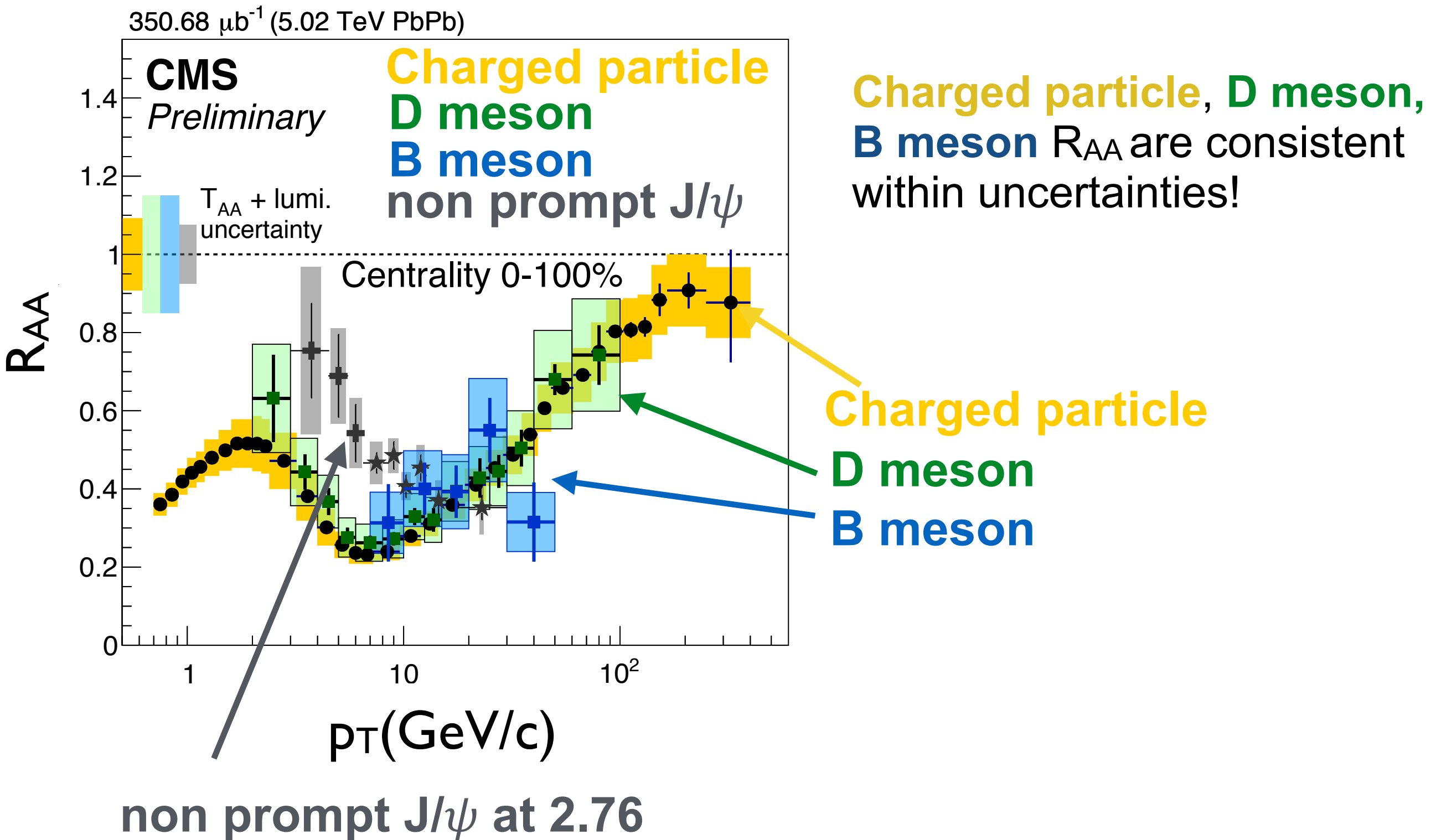
R_{AA} of non prompt J/ψ at 2.76 TeV



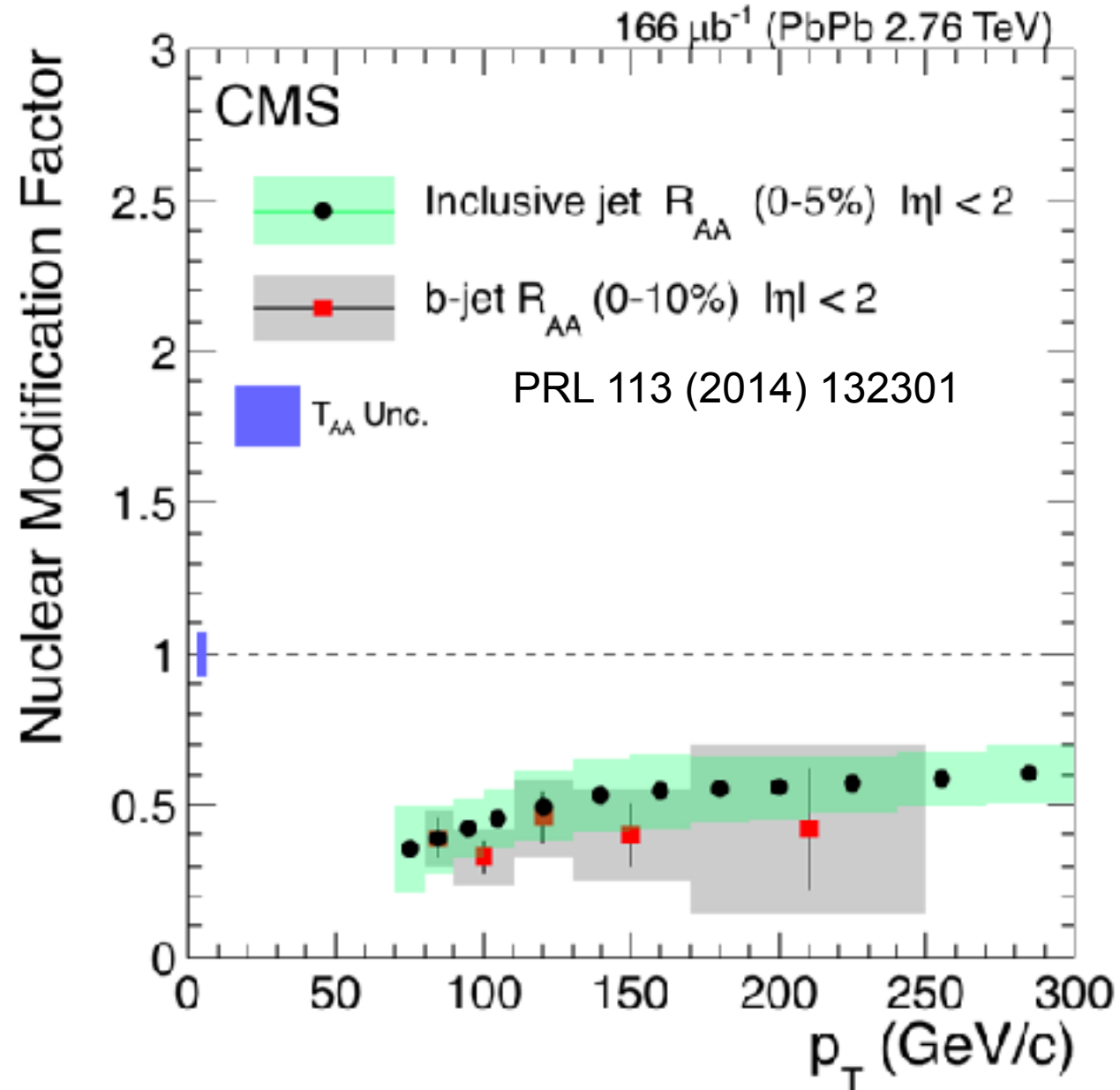
CMS non prompt $1.6 < |y| < 2.4$
CMS non prompt $|y| < 2.4$
ATLAS non prompt $|y| < 2.9$

Strong suppression observed for non prompt J/ψ in PbPb collisions
→ **similar studies can be done in the c,b→electrons channel**

Flavour dependence at 5.02 TeV



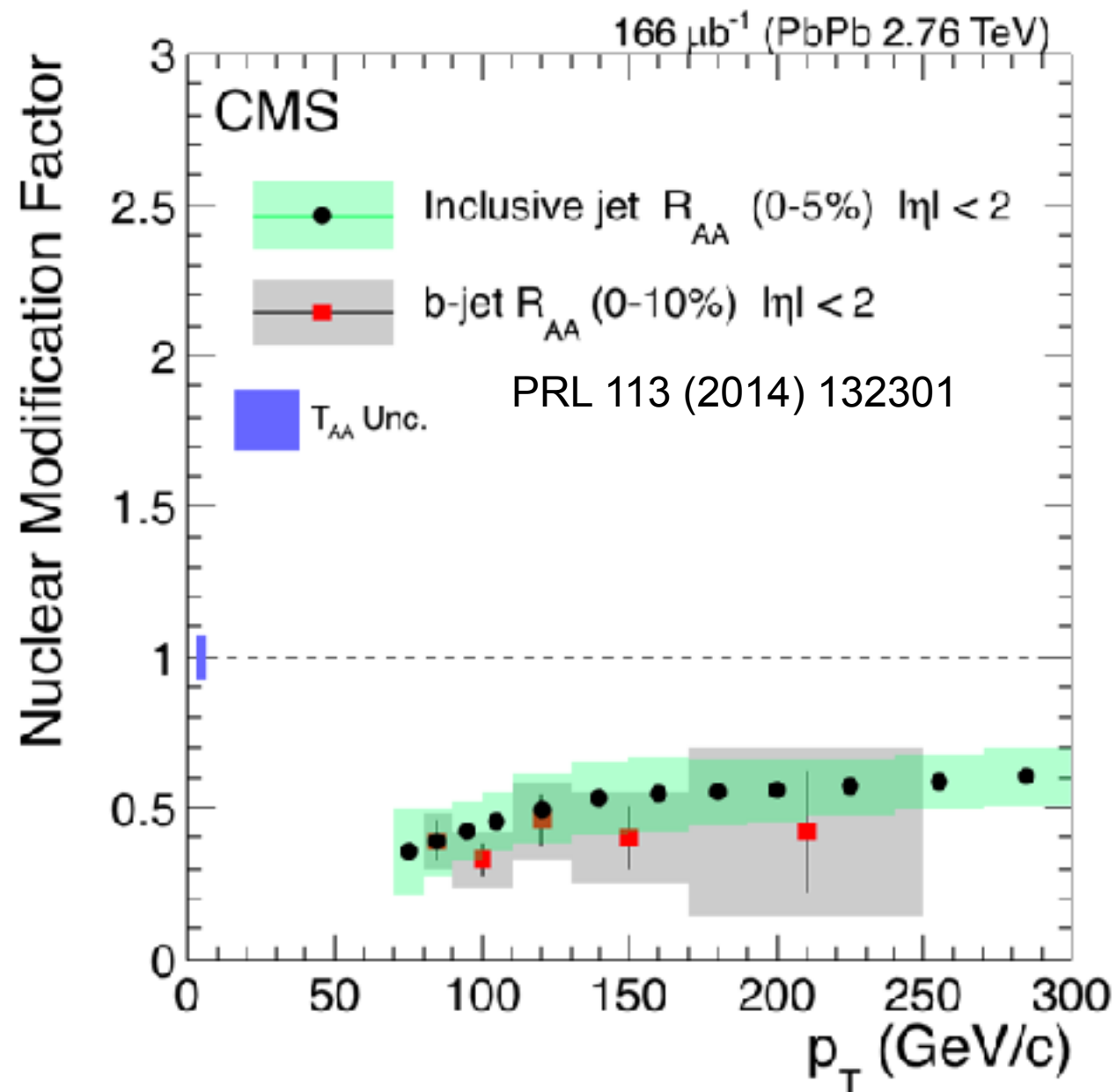
Flavour dependence at higher p_T



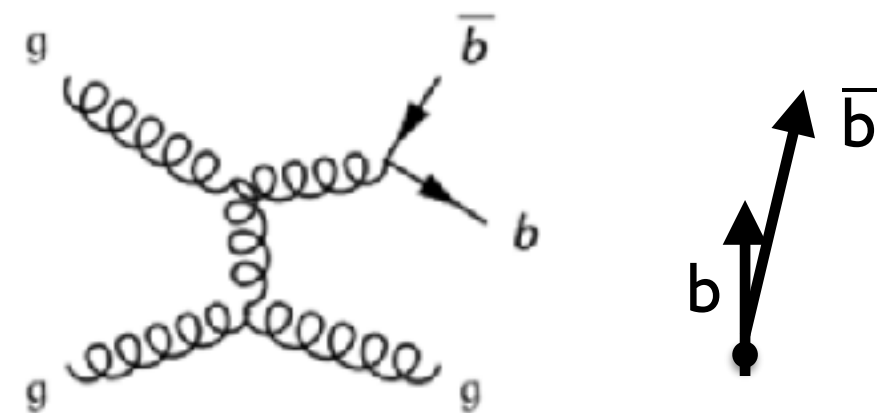
b-jet R_{AA}
inclusive jet R_{AA}

Same suppression for b-jets and inclusive jets at high p_T
Mass difference negligible at high p_T

Flavour dependence at higher p_T



b-jet R_{AA}
inclusive jet R_{AA}



NLO process: Gluon splitting
 → dominant at low opening angles

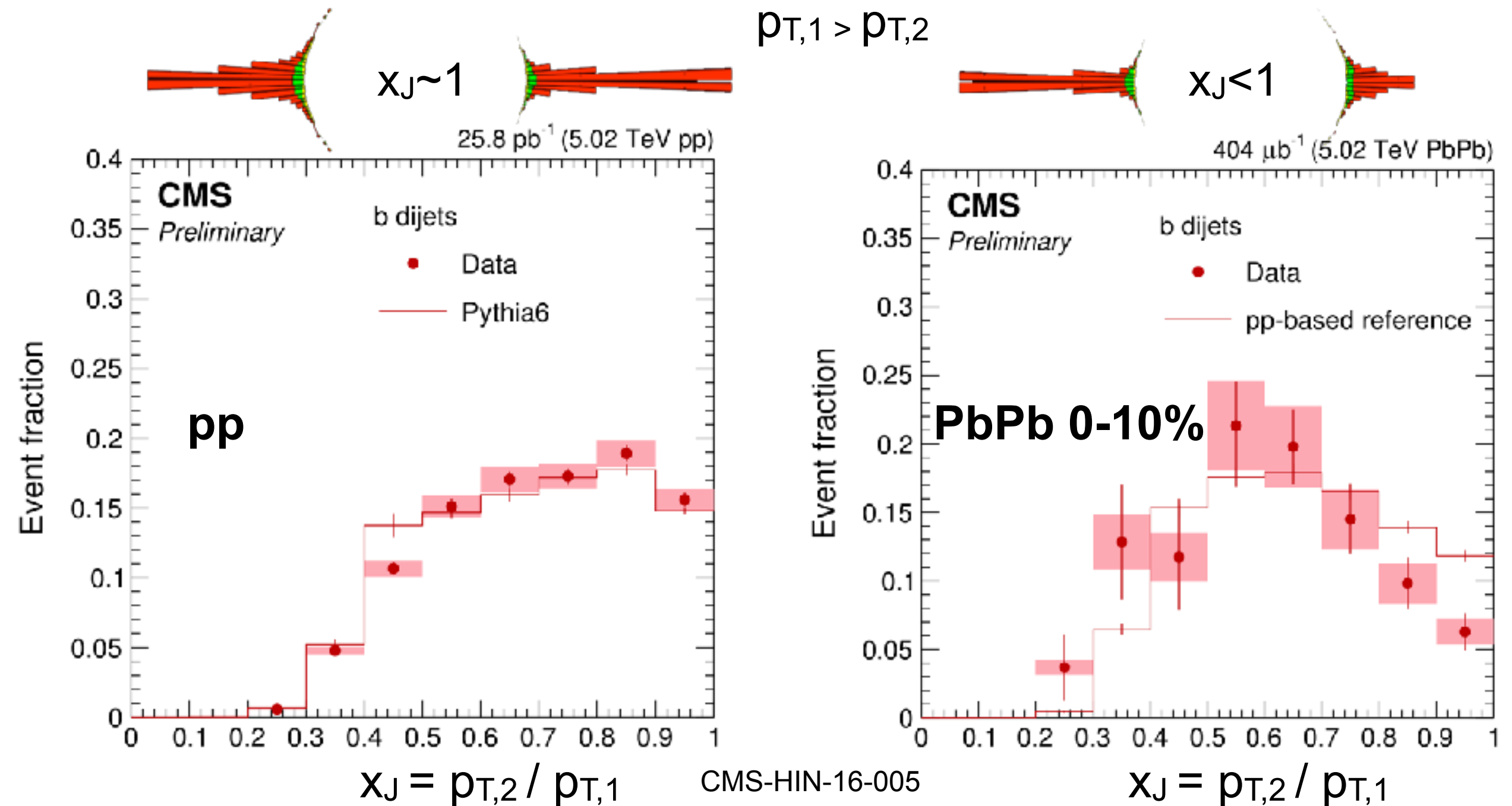
Same suppression for b-jets and inclusive jets at high p_T

Mass difference negligible at high p_T

Large contribution of gluon splitting processes? In GSP case, we are not measuring the b -quark E_{loss} but to some “fat” gluon E_{loss}

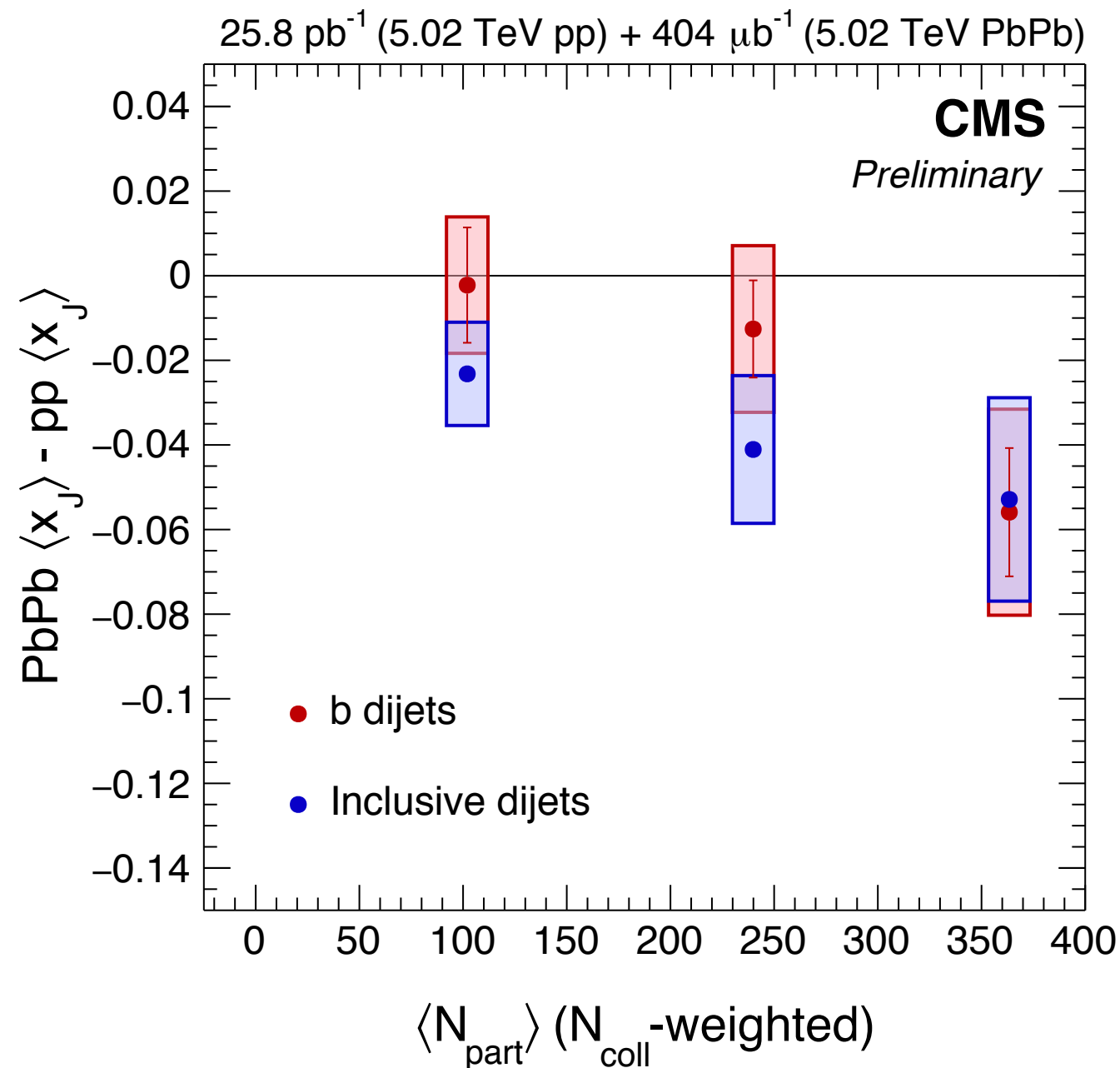
Di-b-jet measurement in PbPb at 5.02 TeV

→ In back-to-back events $b\bar{b}$ production via gluon splitting processes is negligible



x_J distributions of di-b-jets significantly modified in central PbPb collisions!

Di-b-jet measurement in PbPb at 5.02 TeV



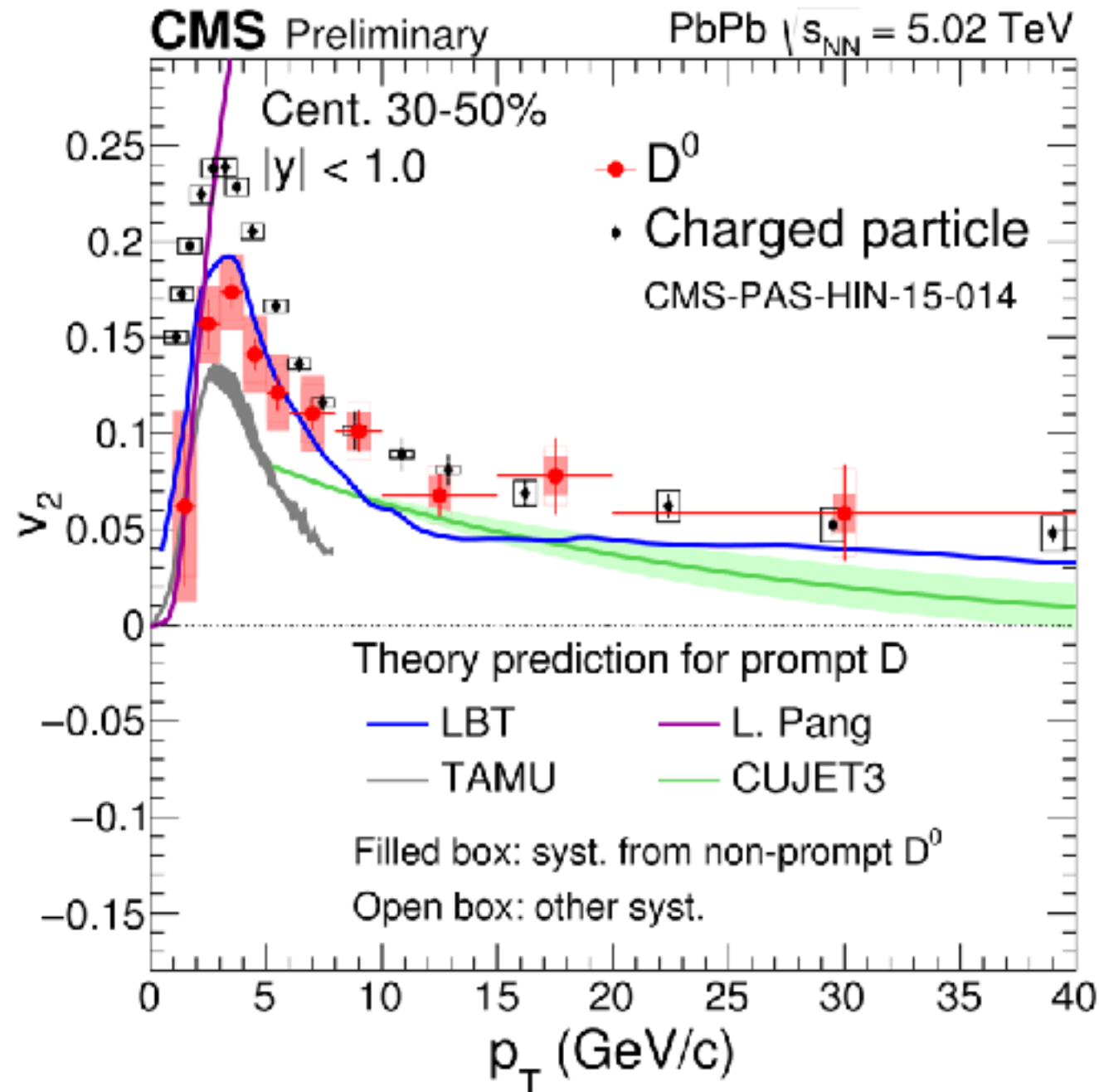
$$x_J = p_{T,2} / p_{T,1}$$

Same average asymmetry
observed for inclusive jets!

CMS-HIN-16-005

There is no significant difference in the suppression of inclusive and b-jets even after excluding the contribution of gluon splitting processes

D^0 elliptic flow in PbPb collisions



CMS-PAS-HIN-16-007

we need **charm quark diffusion** to describe the magnitude of the D meson v_2

Summary and conclusions

CMS experiment showed to be an excellent detector for HF measurements thanks to:

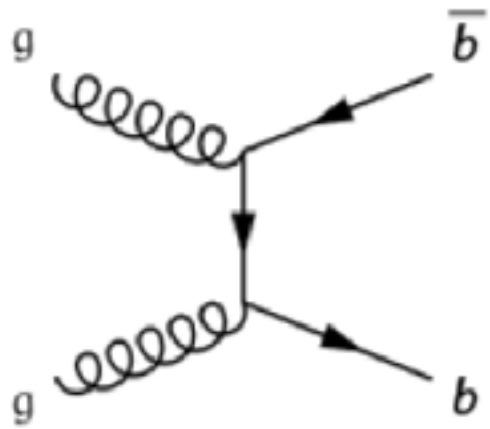
- excellent tracking and vertexing system
- efficiency di-muon trigger system and capability to trigger on HF probes
- possibility to run at very high luminosity and collect huge amount of MB
- PID (that is still in possible in CMS with pixel) is important but not fundamental

How can sPHENIX be complementary:

- both low p_T (flow, collectivity) and high p_T (jet quenching, flavour dependence) are still not completely understood.
- sPhenix could collect huge amount of data to help solving both the puzzles
- the possibility of having both data at LHC and RHIC energies is a unique opportunity to validate theoretical calculations that are now very well “tuned” at LHC.

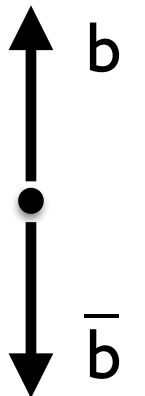
BACKUP

heavy quark production mechanism



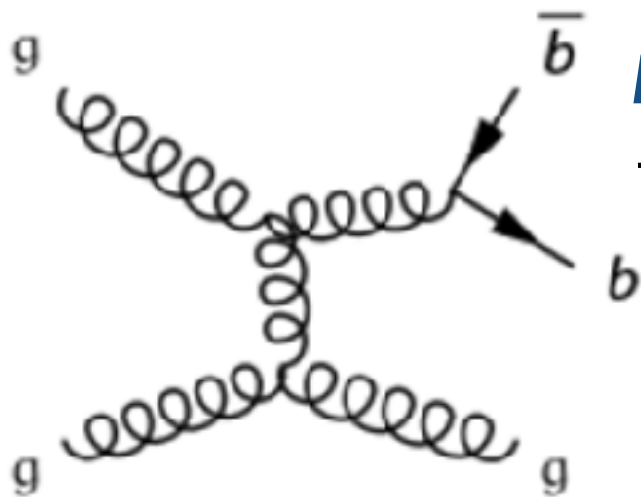
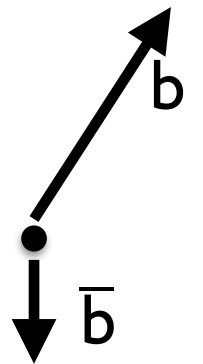
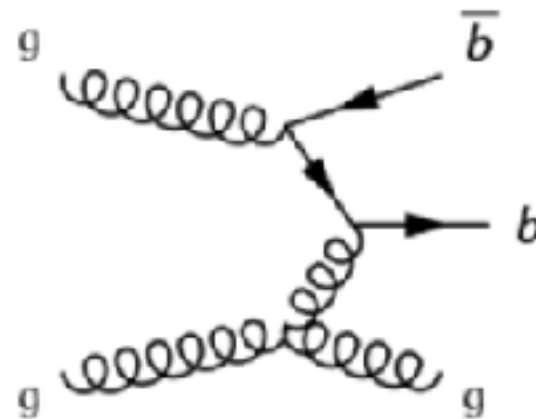
LO process: Flavour Creation (FCR)

→ $b\bar{b}$ produced back-to-back in azimuthal plane and symmetric in p_T



NLO process: Flavour Excitation (FEX)

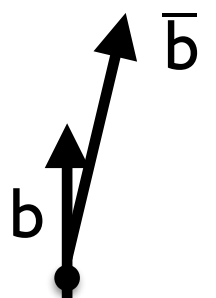
→ $b\bar{b}$ pairs produced asymmetric in p_T and with a broad opening angle



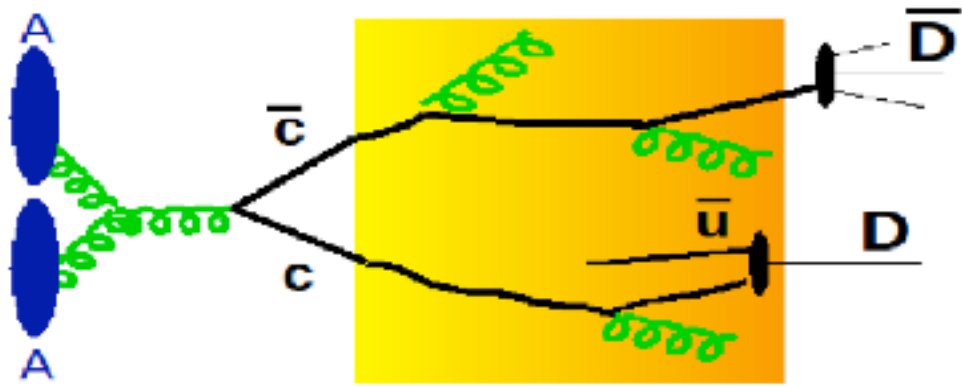
NLO process: Gluon splitting (GSP)

→ produced with small opening angles and asymmetric in p_T

→ $b\bar{b}$ are \bar{b} not involved in the hard scattering but produced later



Reminder on HF energy loss

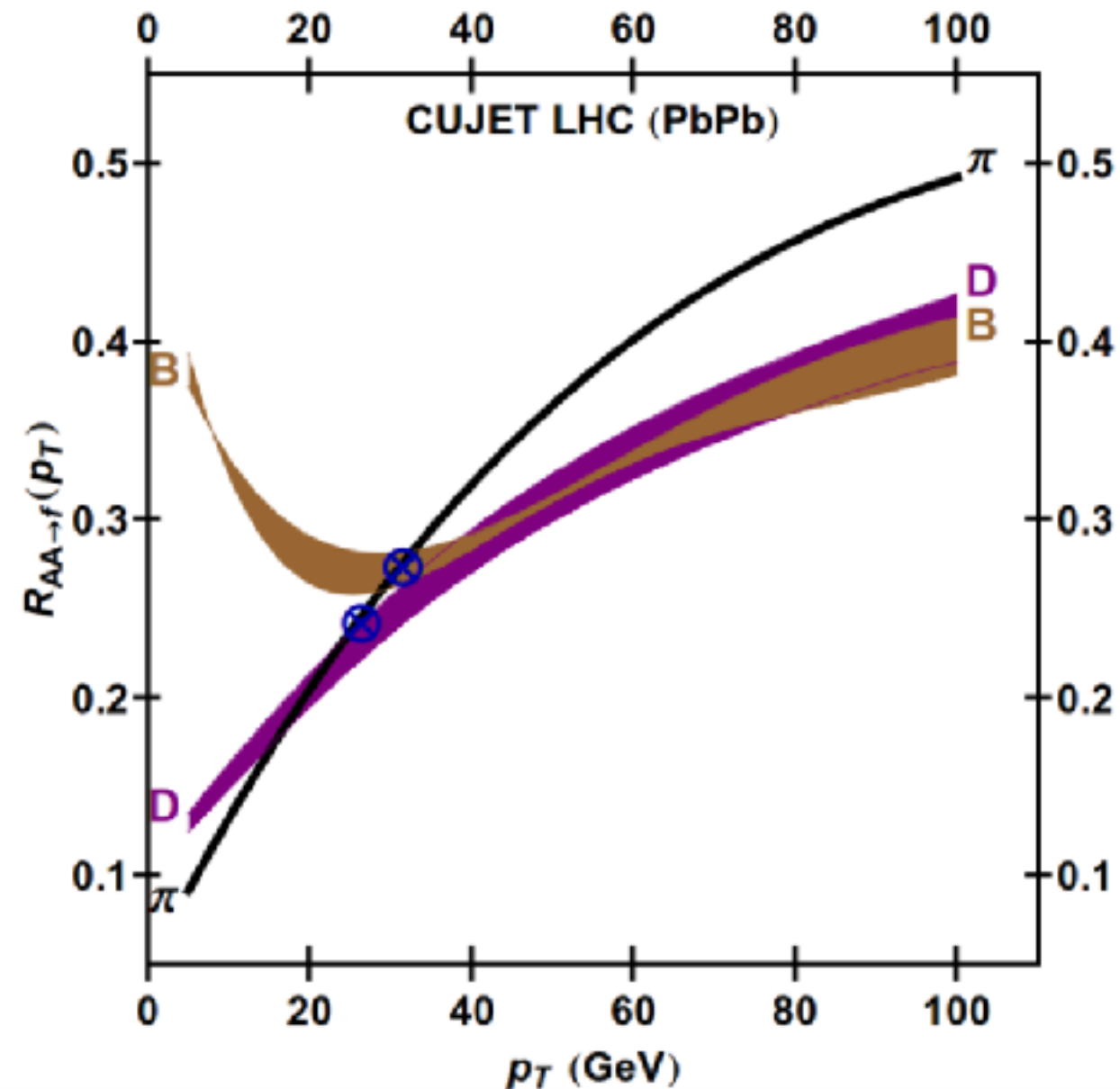


- produced early in the collision, they strongly interact with the deconfined medium

→ In-medium energy loss as a consequence of **radiative and collisional processes**.

Flavour-dependence of radiative energy loss:

- Larger for gluons than for quarks
E.g. in BDMPS model [1] $\langle \Delta E \rangle \propto \alpha_s C_R q L^2$
- Dead cone effect**: gluon radiation suppressed at small angles for massive quarks

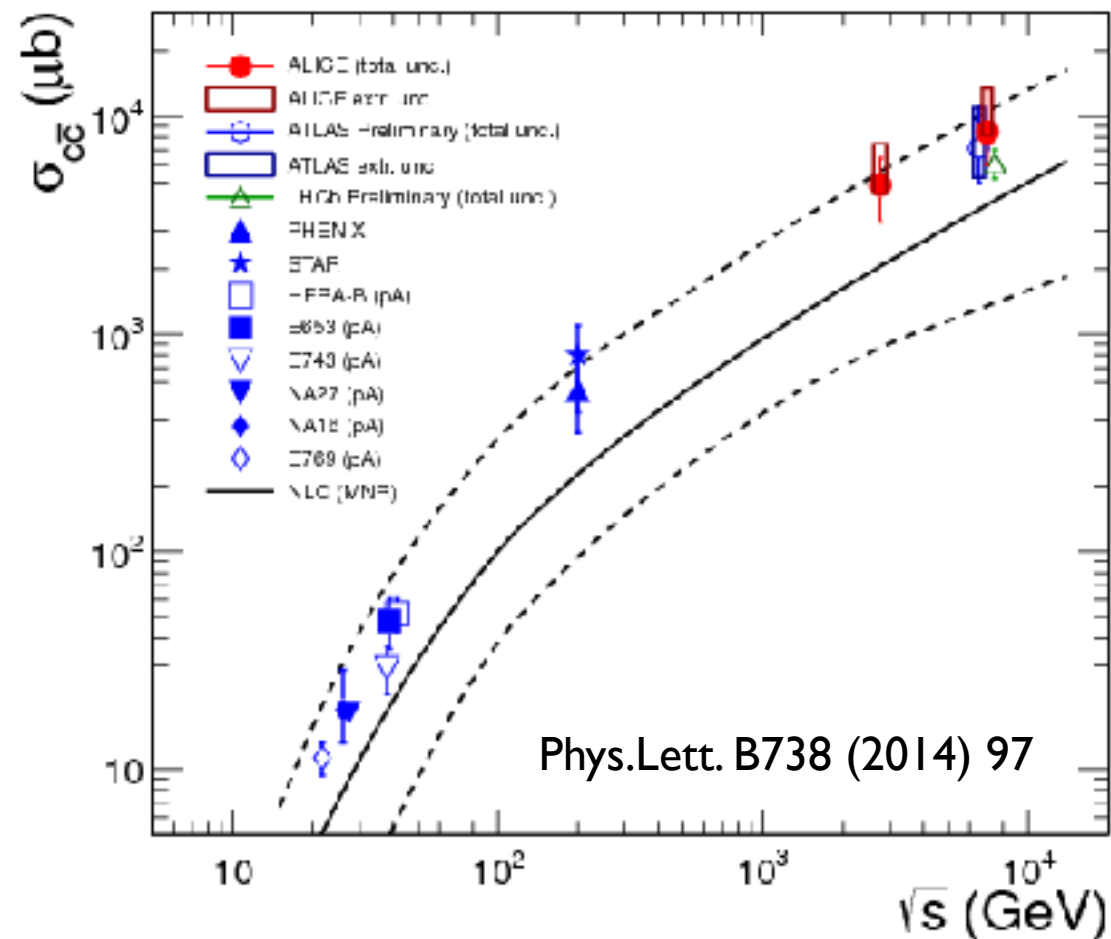


$$\Delta E_g > \Delta E_{u,d,s} > \Delta E_c > \Delta E_b$$

$$\rightarrow R_{AA}^B > R_{AA}^D > R_{AA}^{\text{light}} (??)$$

Heavy-flavours in pp and pPb collisions

Heavy quarks produced in high Q^2 processes at early stages of the collisions



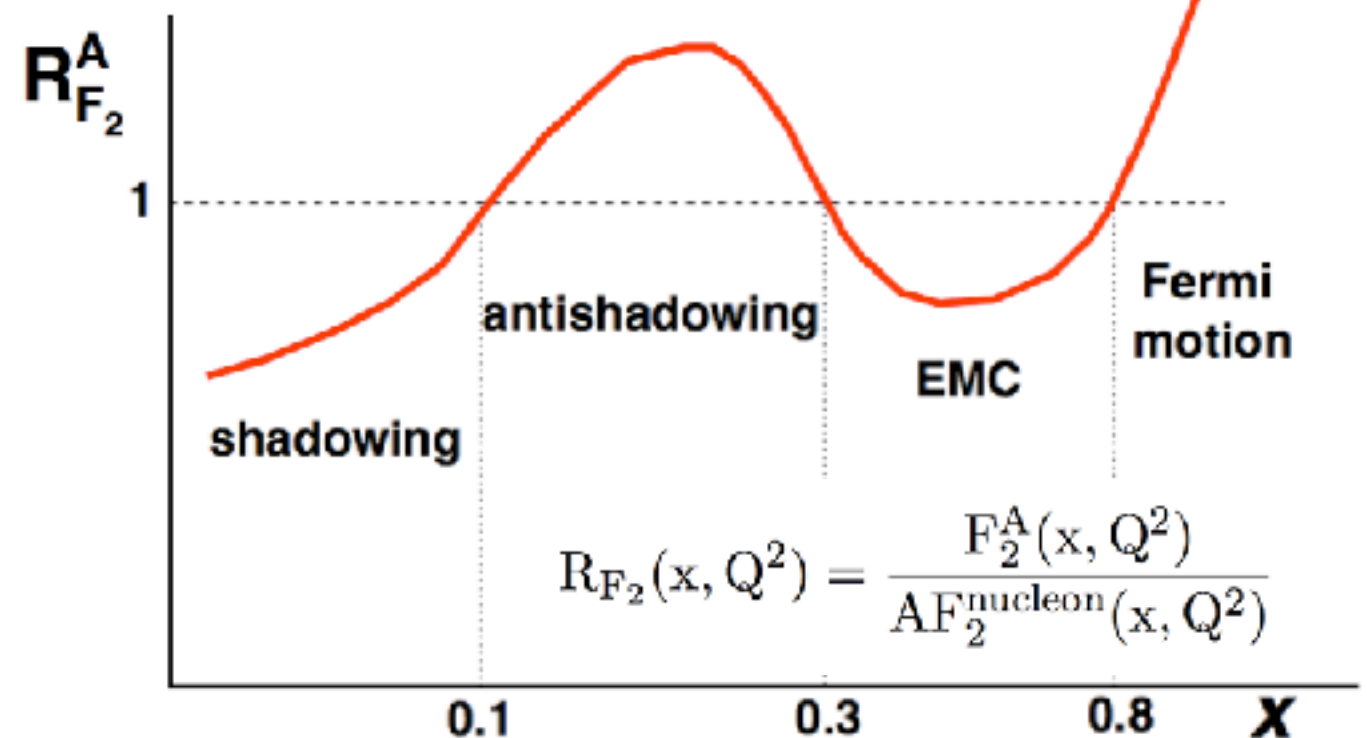
pp:

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pPb:

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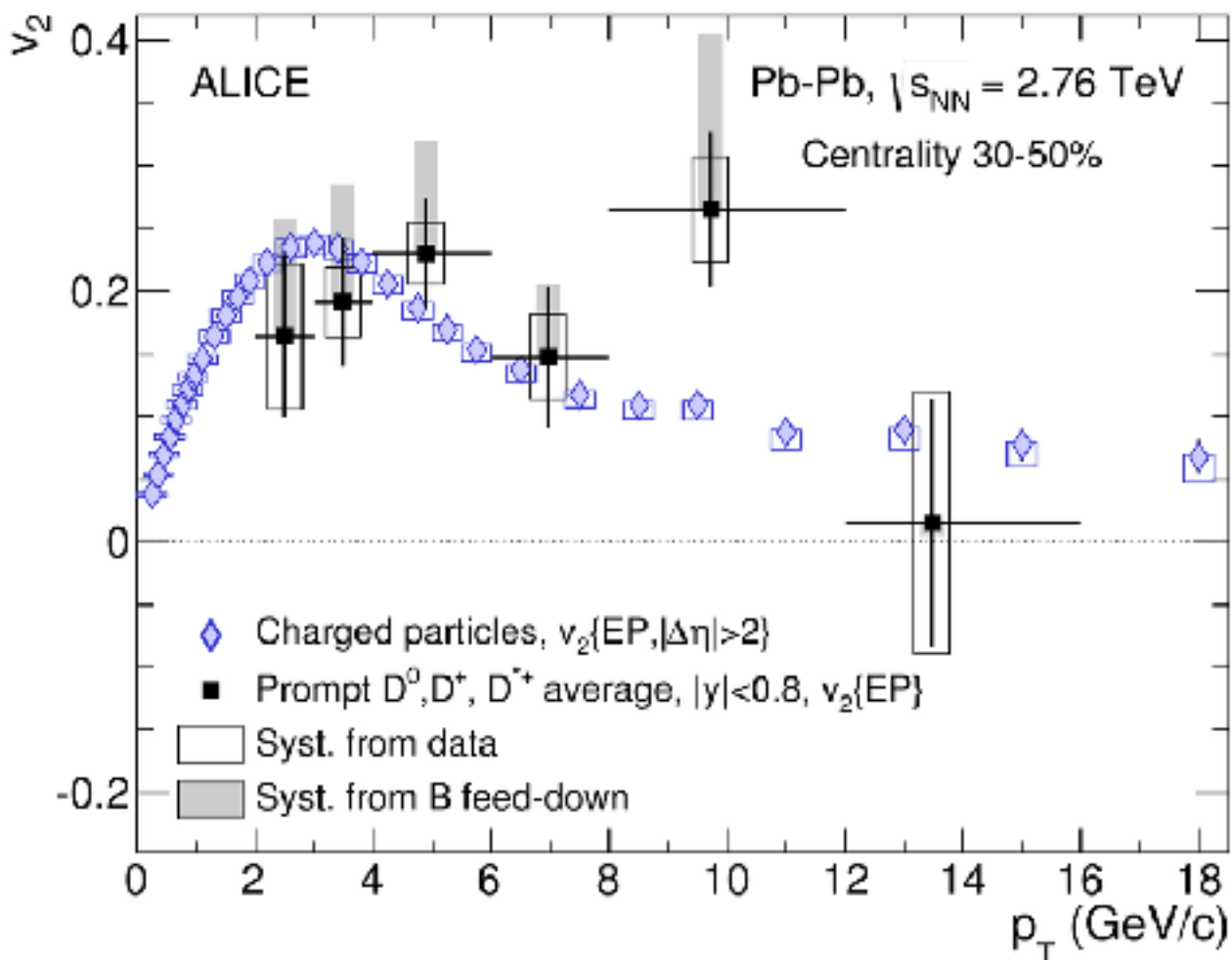
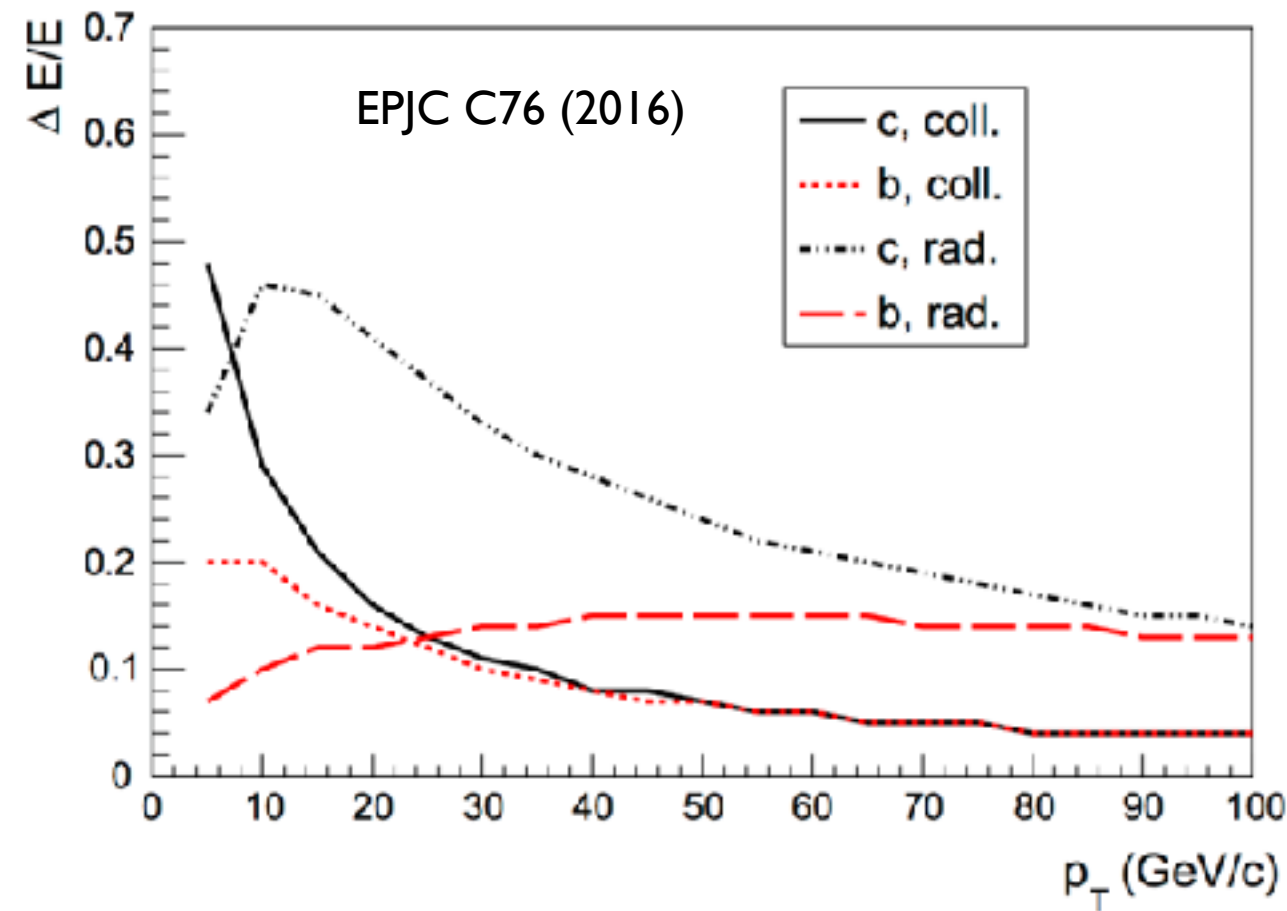
Heavy-flavours in PbPb collisions

Heavy quark energy loss in PbPb:

- collisional vs **radiative** component

Flavour dependence energy loss:

- $\langle \Delta E \rangle \propto \alpha_s C_R q L^2$
- **Dead cone effect**: gluon radiation suppressed at small angles for massive quarks



Collective behaviour:

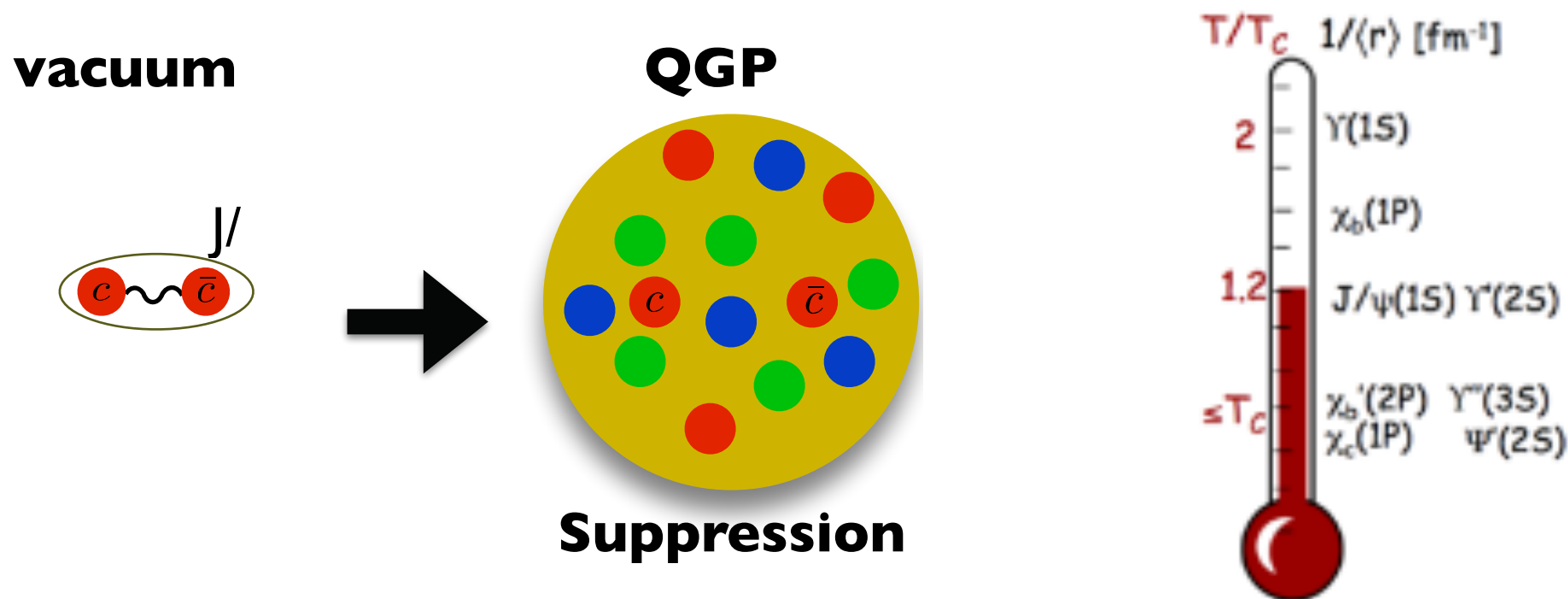
- v_n measurements to study collective behaviour of heavy quarks
- charm recombination in medium?

Charmonia in heavy-ion collisions

Charmonia are bound states of $c\bar{c}$

$$\tau_{\text{formation}}^{c\bar{c}} \lesssim \tau_{\text{formation}}^{QGP} < \tau_{\text{life}}^{QGP} < \tau_{\text{decay}}^{\text{quarkonium}}$$

The presence of QGP should affect charmonia production (yield and kinematics)



Less bounded states melts at lower temperature

$$T^{\text{diss}}(1S) > T^{\text{diss}}(2S) > \dots$$

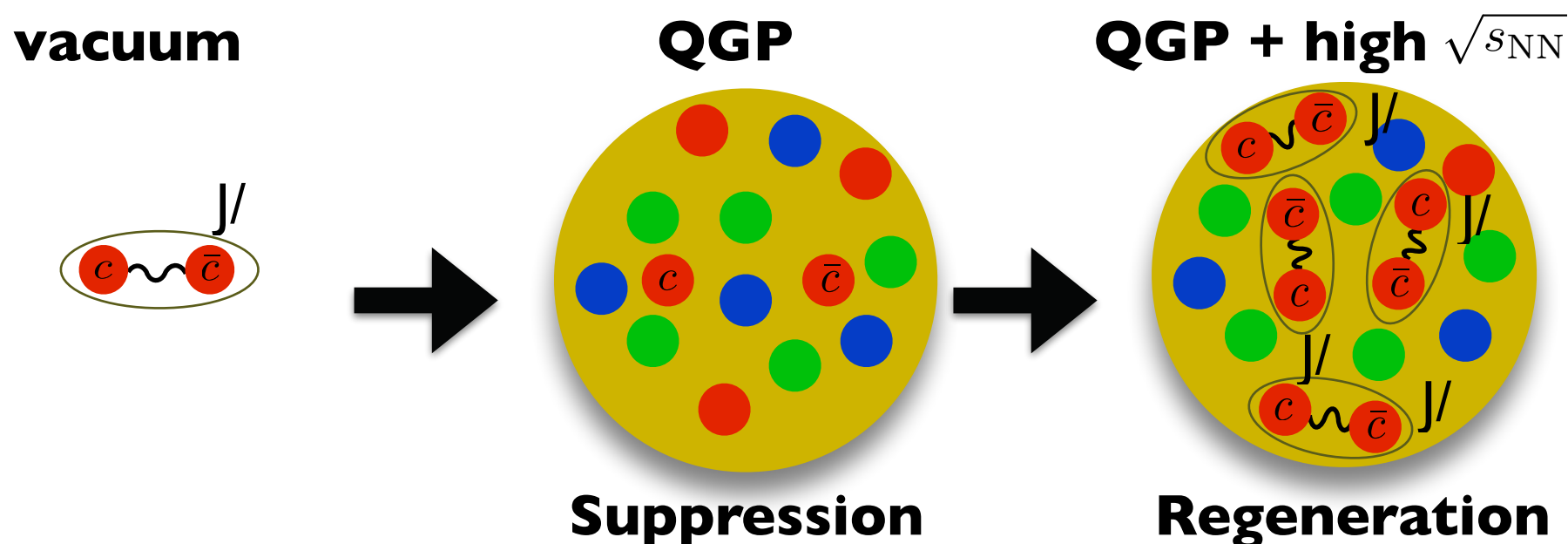
We should observe a hierarchy in the dissociation of different quarkonia states depending on their binding energies

Charmonia in heavy-ion collisions

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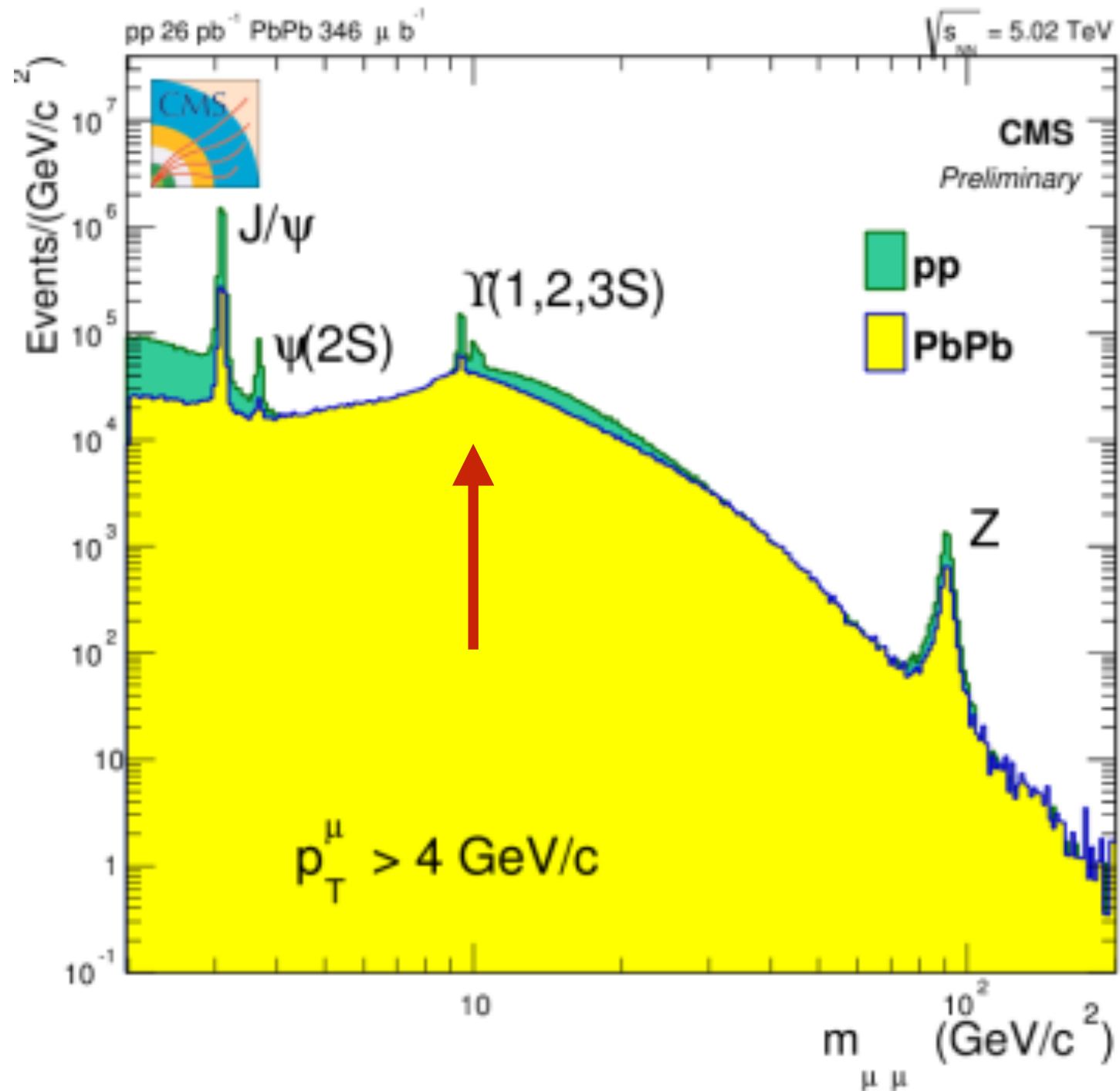


Less bounded states melts at lower temperature

$$T^{\text{diss}}(1S) > T^{\text{diss}}(2S) > \dots$$

But life is not that easy, charm (re)combination in the medium is expected to play a significant role at LHC!

Bottomonia, a cleaner probe for the QGP



$\Upsilon(nS)$ states are less likely to be created via recombination!

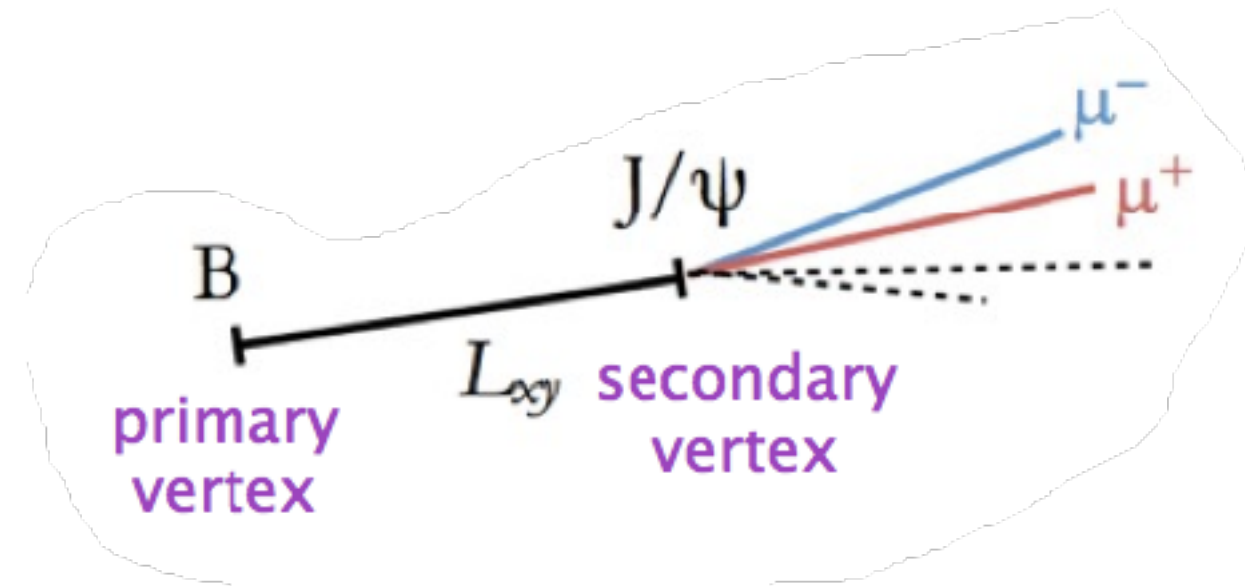
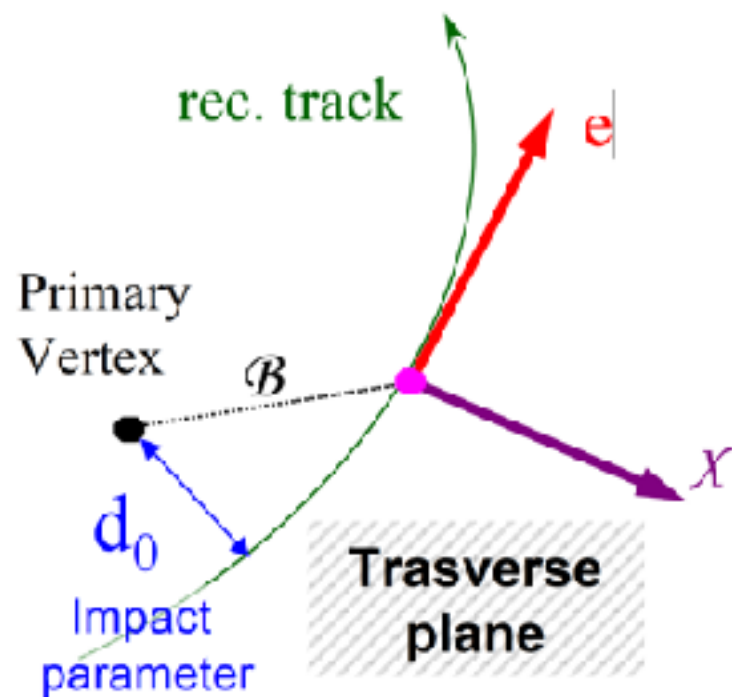
HF models overview

Table 11: Comparative overview of the models for heavy-quark energy loss or transport in the medium described in the previous sections.

<i>Model</i>	<i>Heavy-quark production</i>	<i>Medium modelling</i>	<i>Quark-medium interactions</i>	<i>Heavy-quark hadronisation</i>	<i>Tuning of medium-coupling (or density) parameter(s)</i>
Djordjevic <i>et al.</i> [511–515]	FONLL no PDF shadowing	Glauber model nuclear overlap no fl. dyn. evolution	rad. + coll. energy loss finite magnetic mass	fragmentation	Medium temperature fixed separately at RHIC and LHC
WHDG [459, 519]	FONLL no PDF shadowing	Glauber model nuclear overlap no fl. dyn. evolution	rad. + coll. energy loss	fragmentation	RHIC (then scaled with $dN_{\text{ch}}/d\eta$)
Vitev <i>et al.</i> [422, 460]	non-zero-mass VFNS no PDF shadowing	Glauber model nuclear overlap ideal fl. dyn. 1+1d Bjorken expansion	radiative energy loss in-medium meson dissociation	fragmentation	RHIC (then scaled with $dN_{\text{ch}}/d\eta$)
AdS/CFT (HG) [624, 625]	FONLL no PDF shadowing	Glauber model nuclear overlap no fl. dyn. evolution	AdS/CFT drag	fragmentation	RHIC (then scaled with $dN_{\text{ch}}/d\eta$)
POWLANG [507–509, 585, 586]	POWHEG (NLO) EPS09 (NLO) PDF shadowing	2+1d expansion with viscous fl. dyn. evolution	transport with Langevin eq. collisional energy loss	fragmentation recombination	assume pQCD (or l-QCD U potential)
MC@_sHQ+EPOS2 [528–530]	FONLL EPS09 (LO) PDF shadowing	3+1d expansion (EPOS model)	transport with Boltzmann eq. rad. + coll. energy loss	fragmentation recombination	QGP transport coefficient fixed at LHC, slightly adapted for RHIC
BAMPS [537–540]	MC@NLO no PDF shadowing	3+1d expansion parton cascade	transport with Boltzmann eq. rad. + coll. energy loss	fragmentation	RHIC (then scaled with $dN_{\text{ch}}/d\eta$)
TAMU [491, 565, 606]	FONLL EPS09 (NLO) PDF shadowing	2+1d expansion ideal fl. dyn.	transport with Langevin eq. collisional energy loss diffusion in hadronic phase	fragmentation recombination	assume l-QCD U potential
UrQMD [608–610]	PYTHIA no PDF shadowing	3+1d expansion ideal fl. dyn.	transport with Langevin eq. collisional energy loss	fragmentation recombination	assume l-QCD U potential
Duke [587, 628]	PYTHIA EPS09 (LO) PDF shadowing	2+1d expansion viscous fl. dyn.	transport with Langevin eq. rad. + coll. energy loss	fragmentation recombination	QGP transport coefficient fixed at RHIC and LHC (same value)

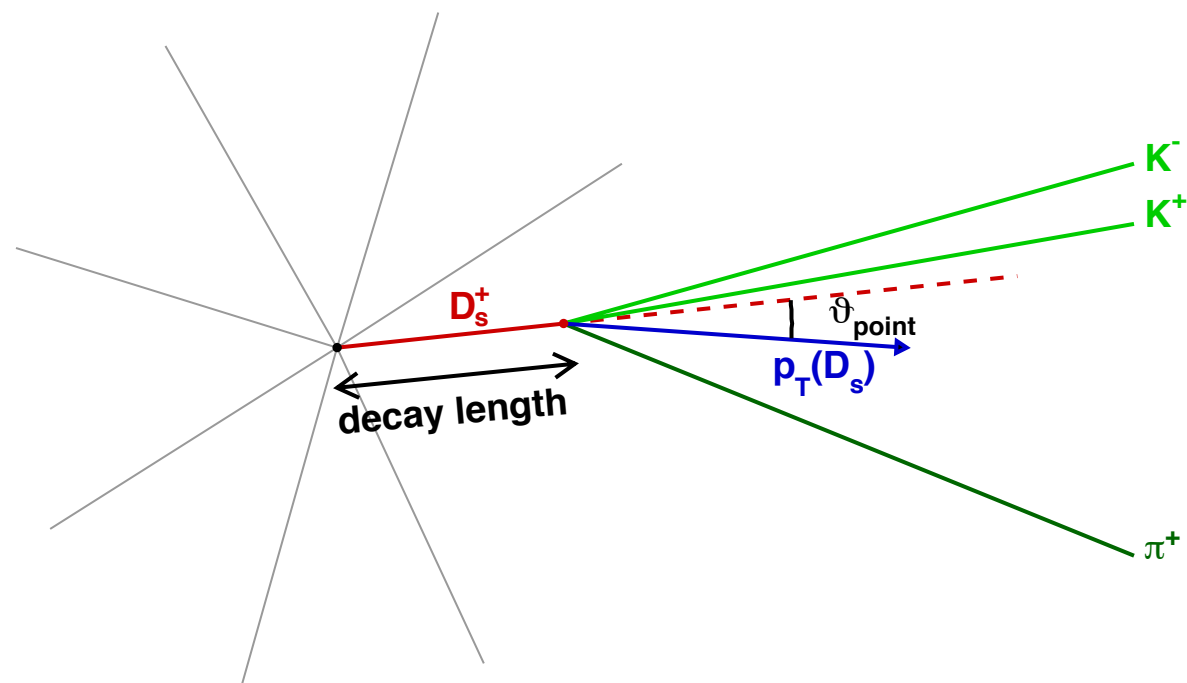
[1506.03981]

Our experimental tools



Displayed J/ψ from B decays

Semi-leptonic electrons and muons from c and b quarks



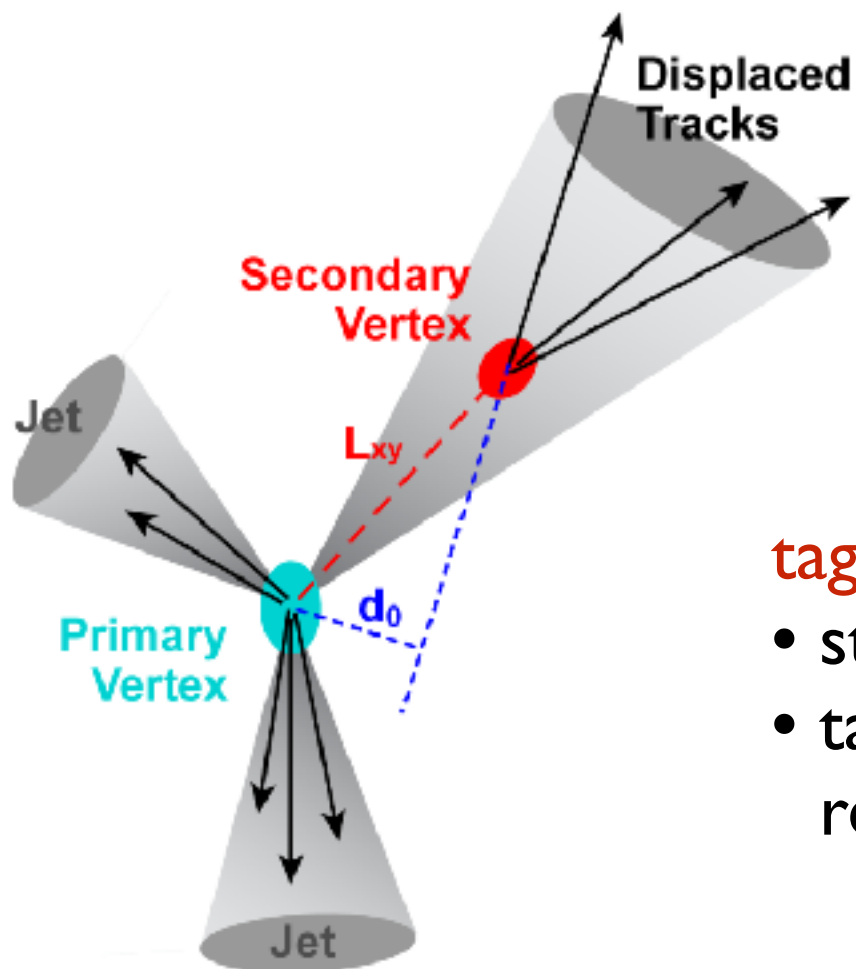
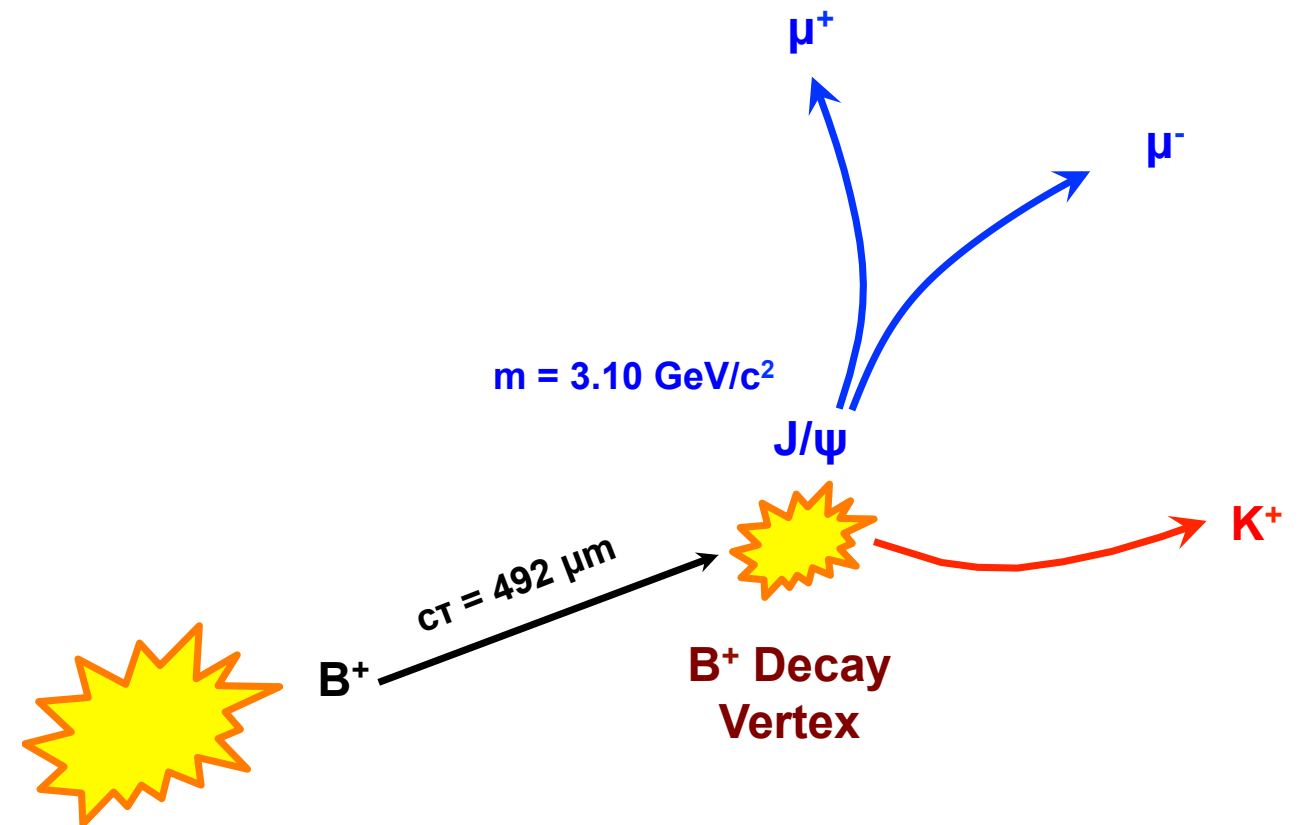
Fully reconstructed D meson decays:

- $D^0 \rightarrow K^- + \pi^+$
- $D^+ \rightarrow K^- + \pi^+ + \pi^+$
- $D^{*+} \rightarrow D^0 + \pi^+$
- $D_s^+ \rightarrow \phi + \pi^+$

Our experimental tools

Fully reconstructed B meson decays:

- $B^+ \rightarrow J/\psi K^+ \rightarrow \mu^+ \mu^- K^+$
- $B^0 \rightarrow J/\psi K^{0*} \rightarrow \mu^+ \mu^- K^+ \pi^-$
- $B_s \rightarrow J/\psi \phi \rightarrow \mu^+ \mu^- K^+ K^-$



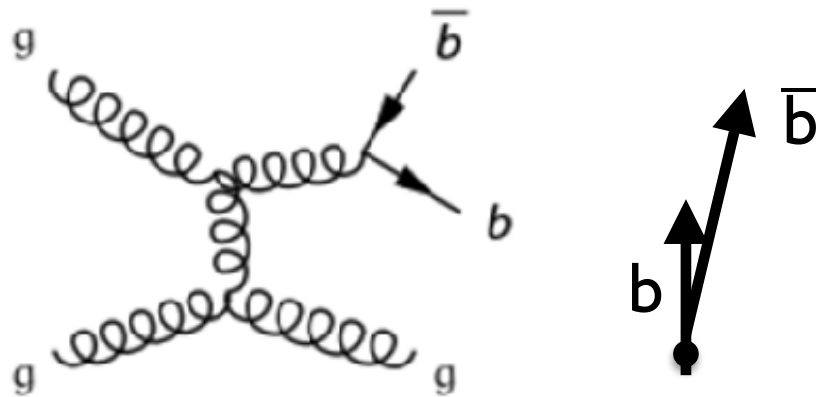
tagged c- and b-jets

- standard jet reconstruction
- tagging based on the displacement with respect to the primary vertex

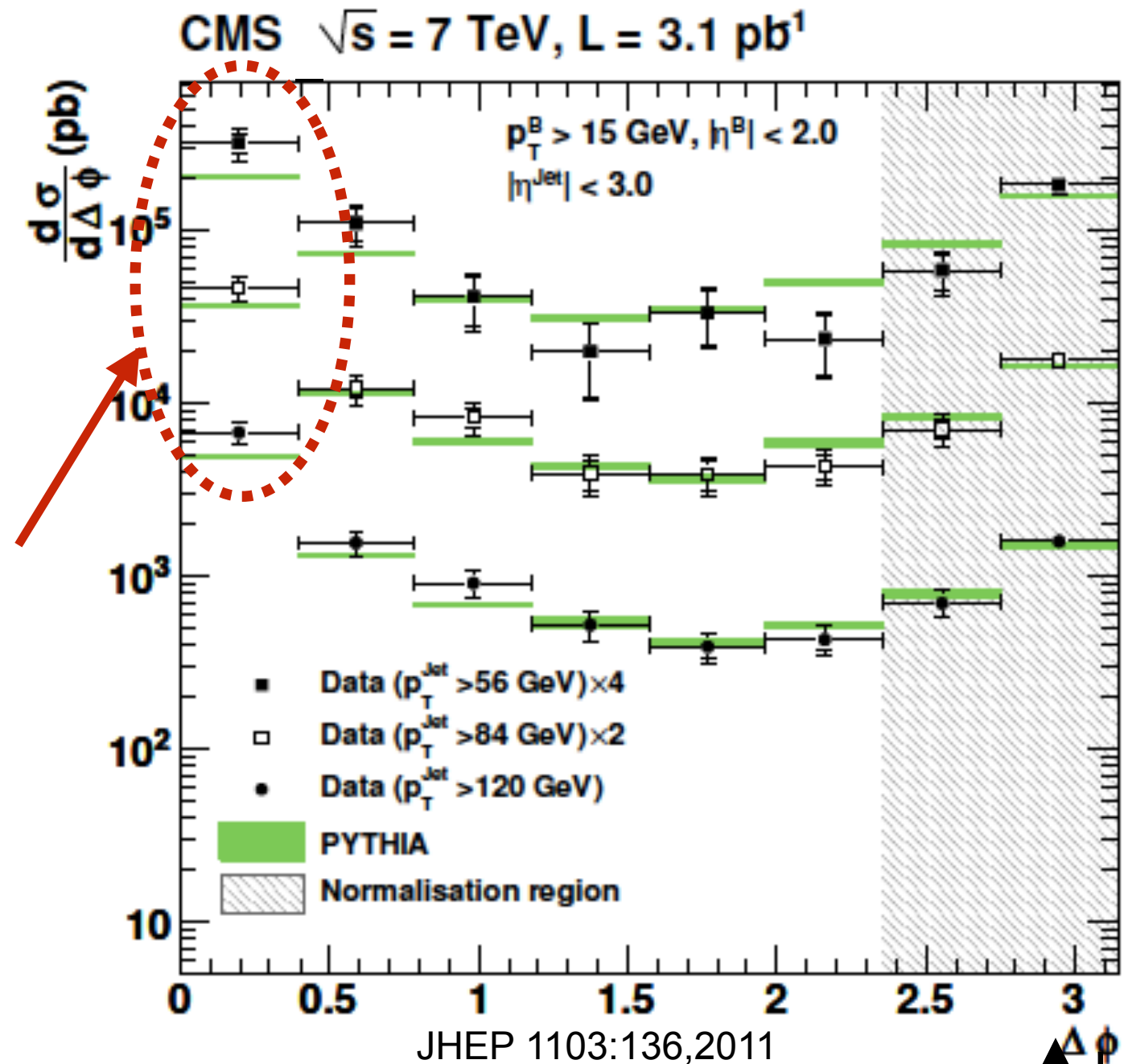
$B\bar{B}$ $\Delta\phi$ correlations

NLO process: Gluon splitting (GSP)

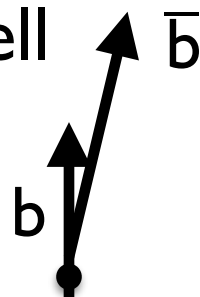
→ produced with small opening angles and asymmetric in p_T



$B\bar{B}$ correlations strongly affected by gluon splitting processes at low $\Delta\phi$



Gluon splitting (GS) contribution not well modelled by most of the calculations
 → *GS contribution underestimated by models*



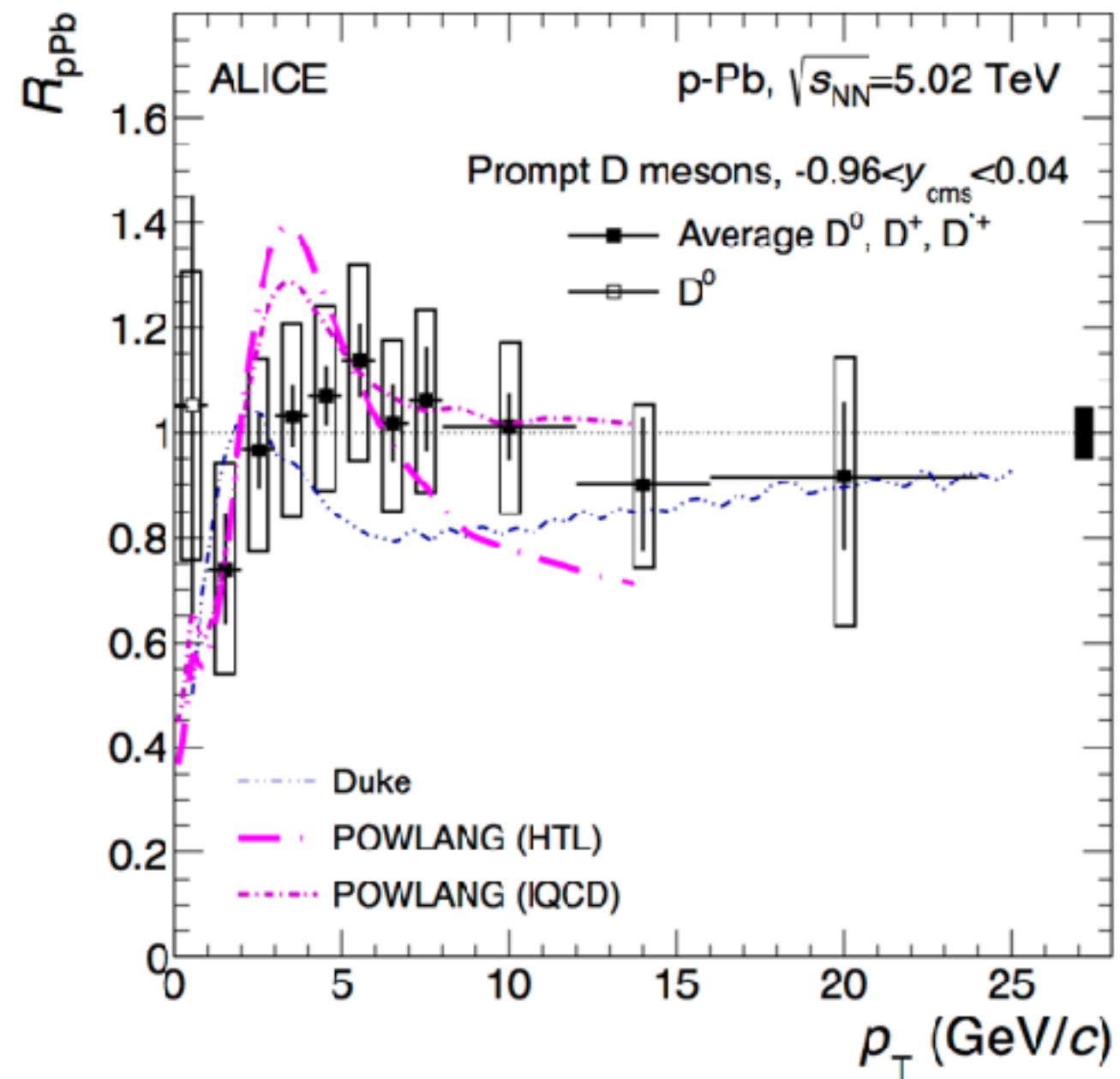
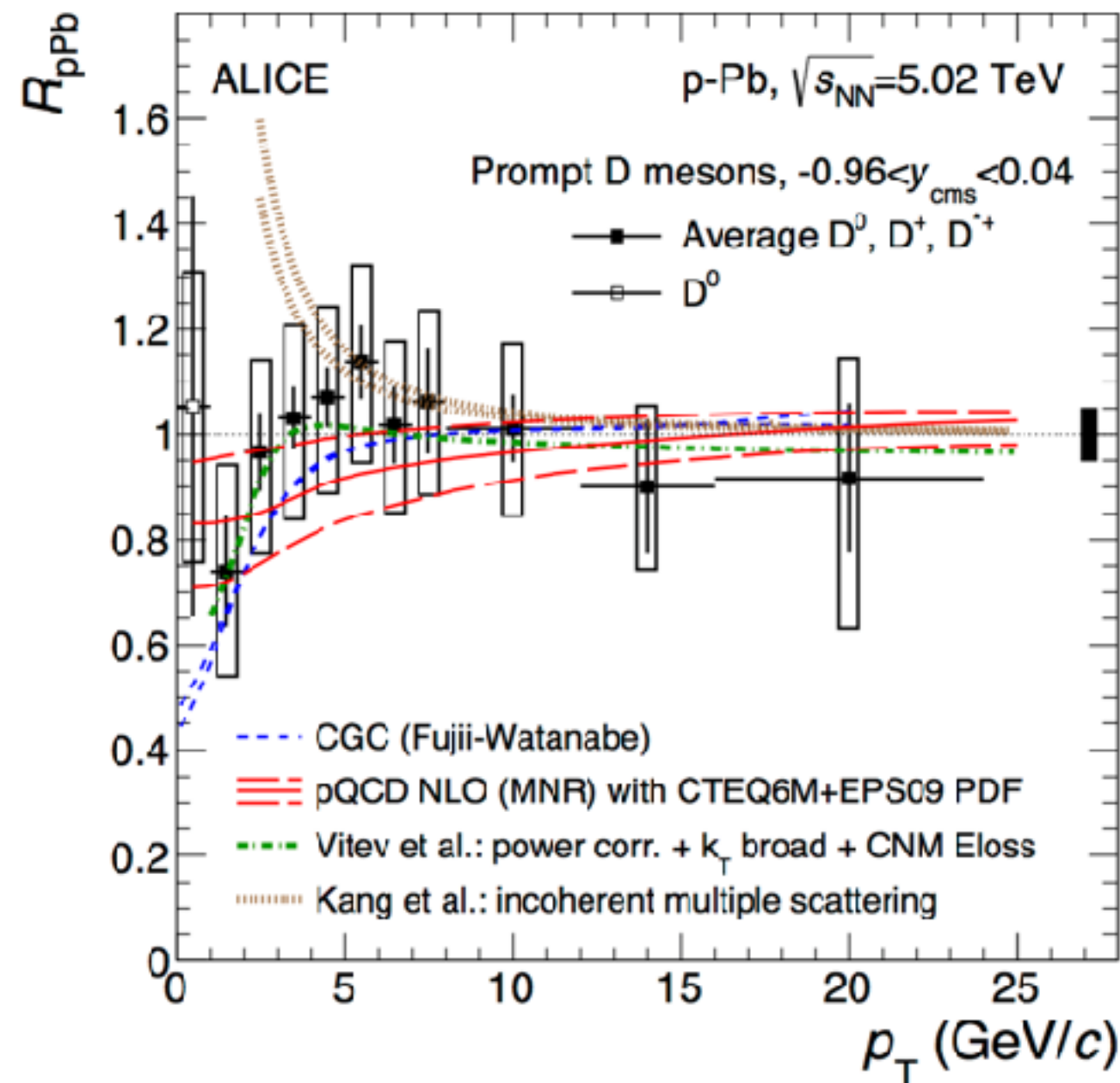
$\Delta\phi(B-\bar{B})$



D^0 production in pPb collisions

ALICE D measurements at 5.02 TeV, $|y| < 0.5$

ALICE, arXiv:1605.07569

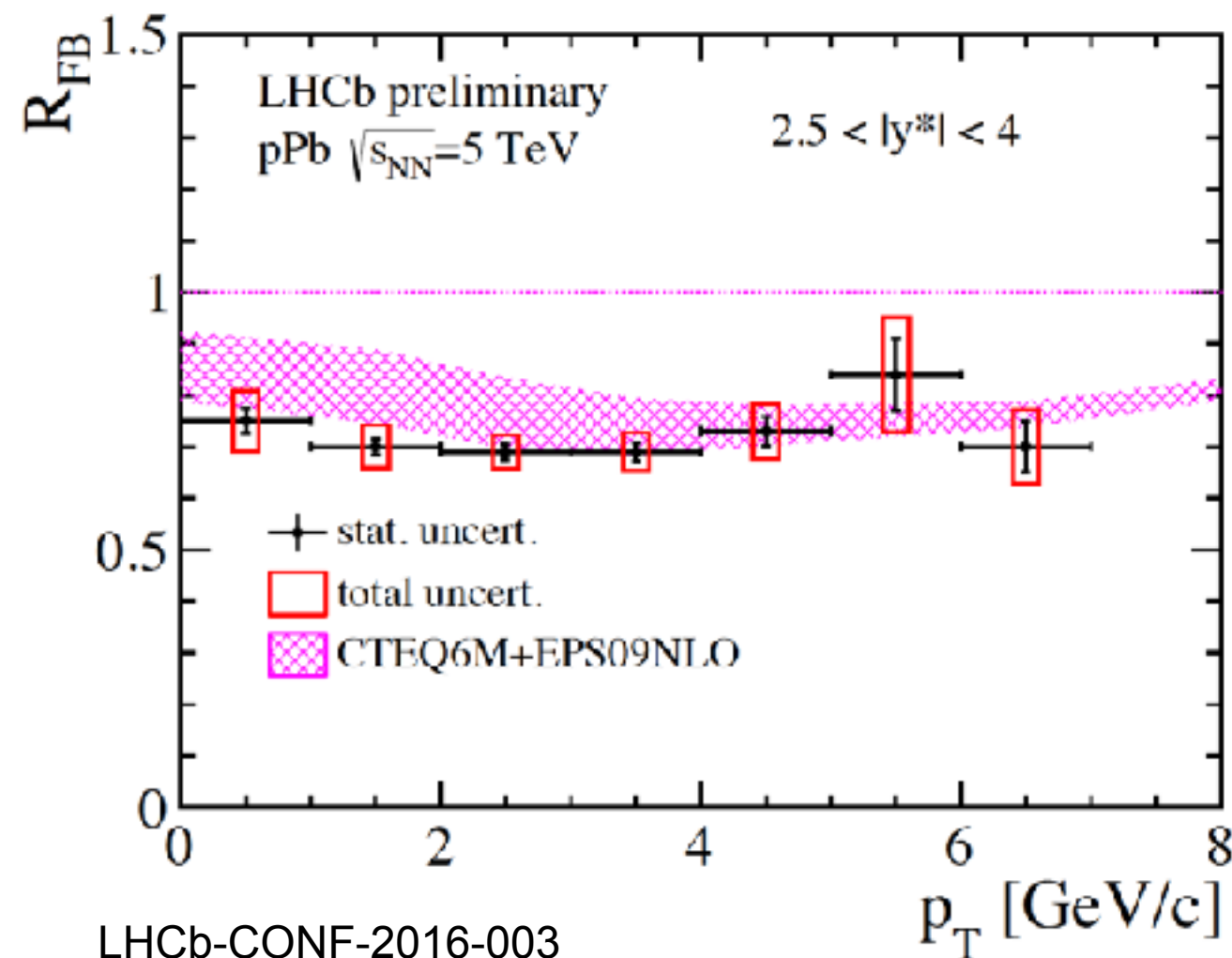


R_{pA} well described by Cold Nuclear Matter (CNR) models and consistent with unity at high p_T !

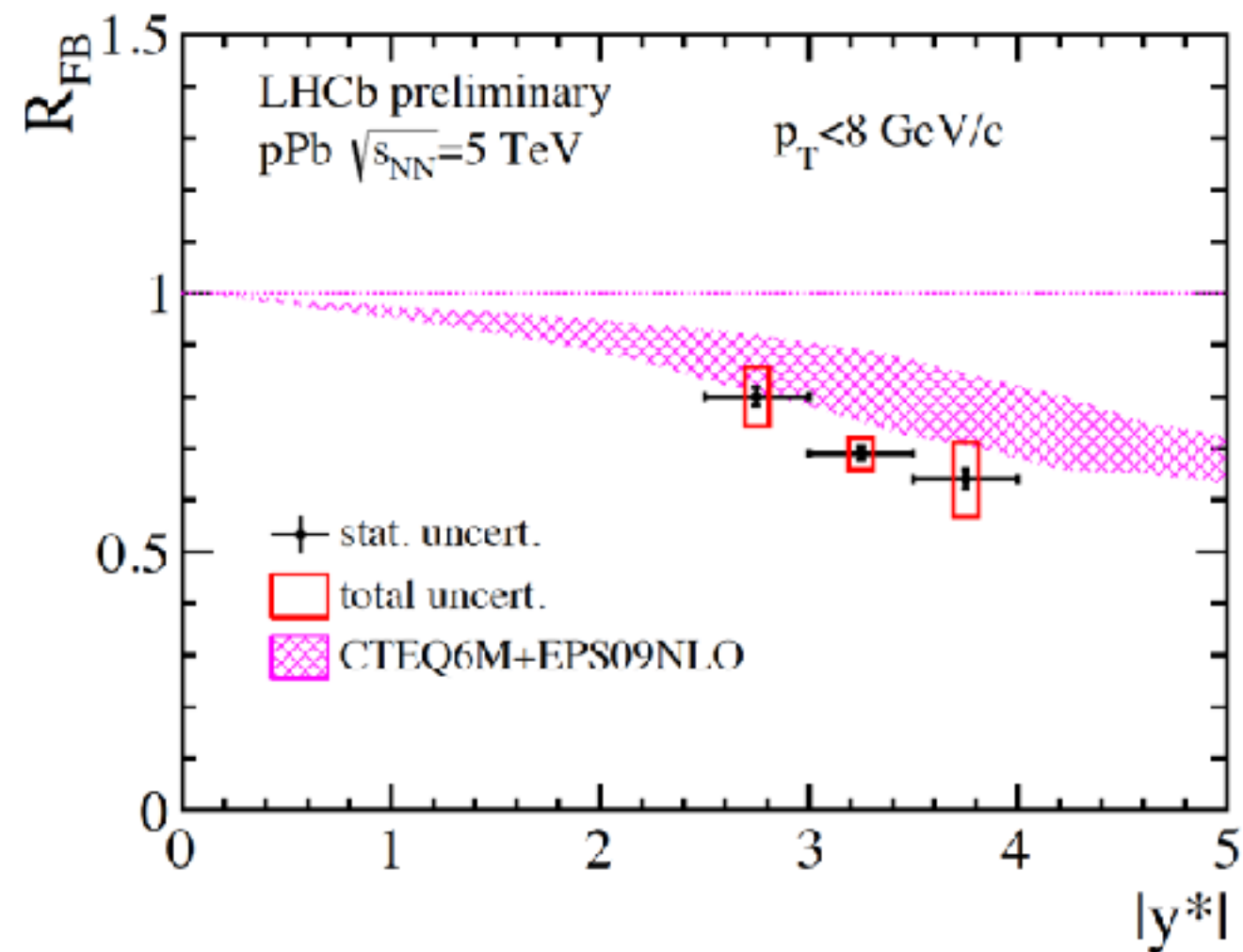
Not possible to discriminate between various models with current uncertainties

D^0 meson R_{pA} at 5.02 TeV

LHCb D^0 measurement at 5.02 TeV in forward (F) and backward (B) region as a function of transverse momentum and rapidity

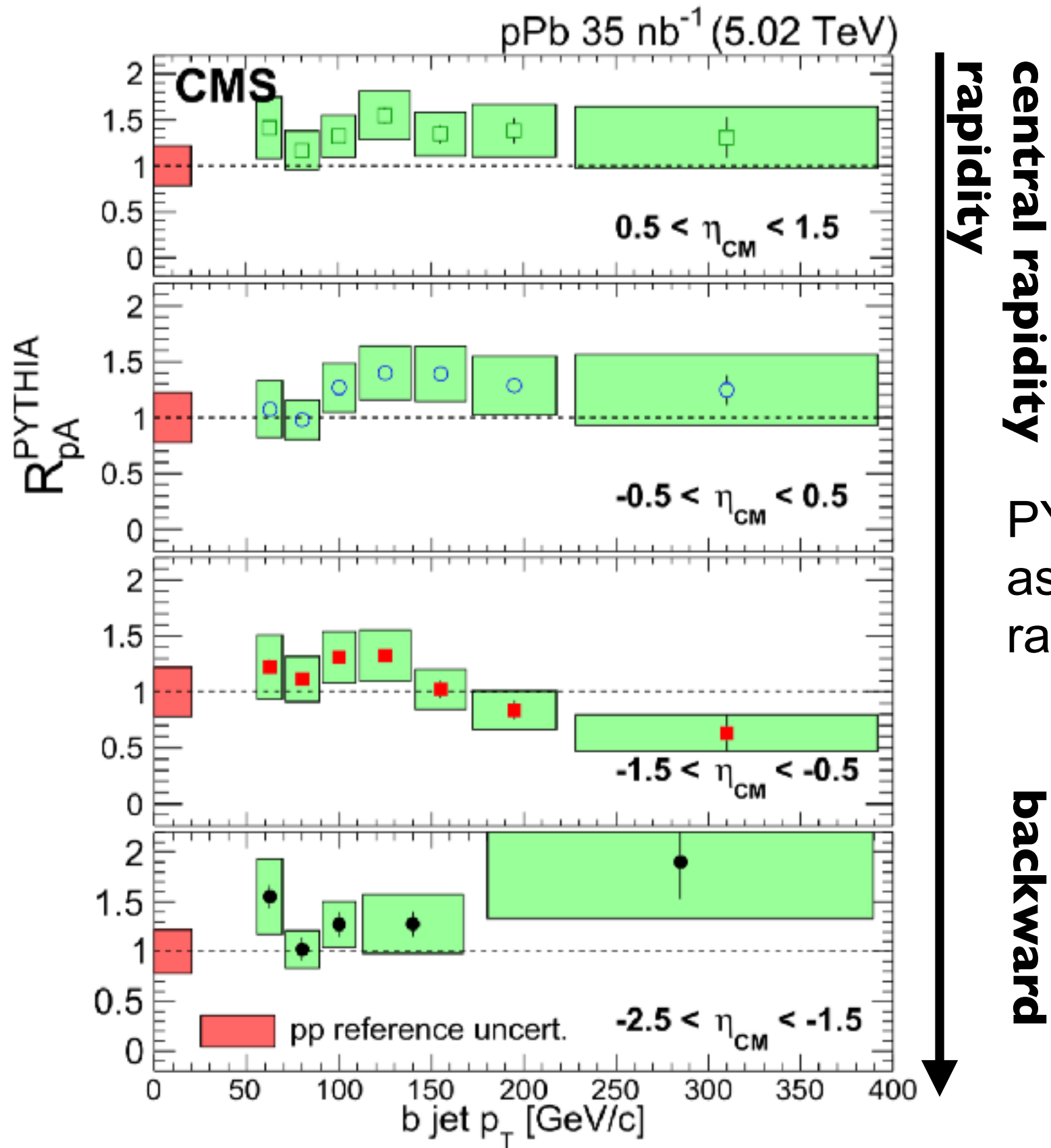


LHCb-CONF-2016-003



R_{pA} and R_{FB} described by to NLO prediction that include EPS09 parametrisation of the nuclear PDFs

b-jet nuclear modification factor in pPb



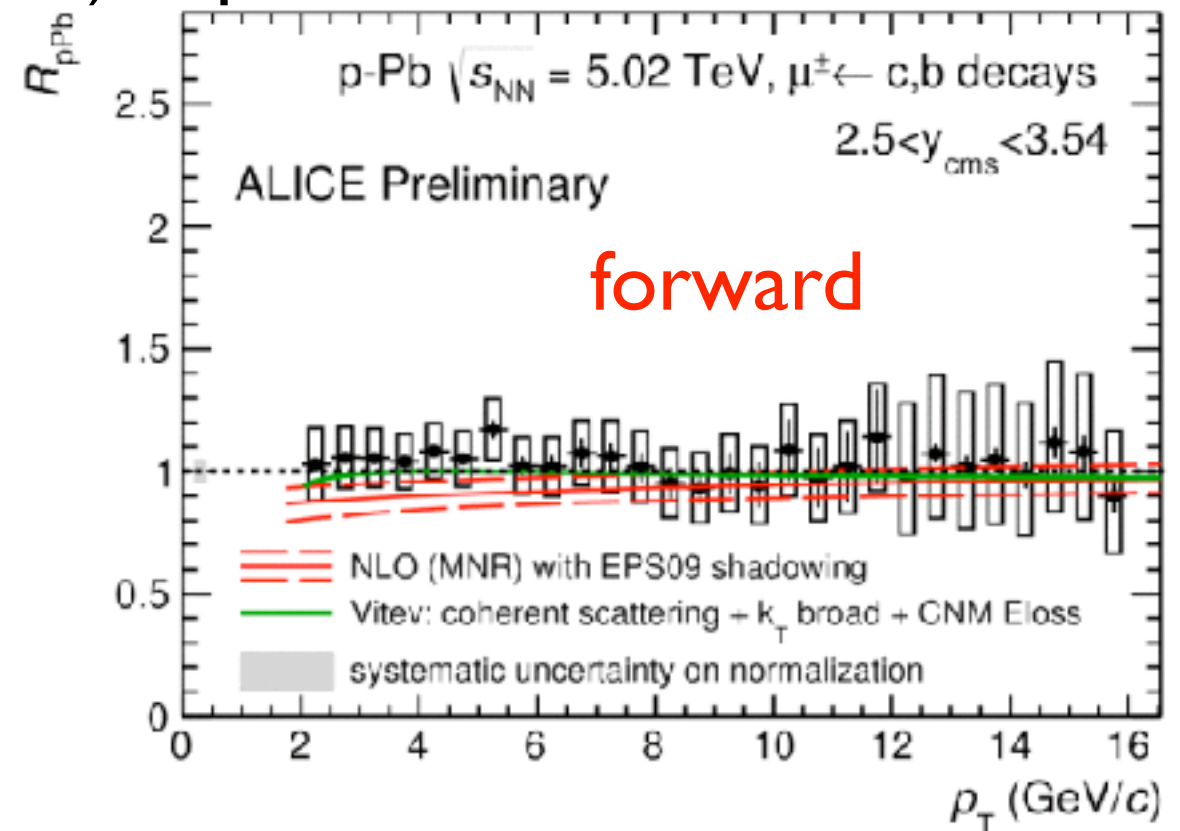
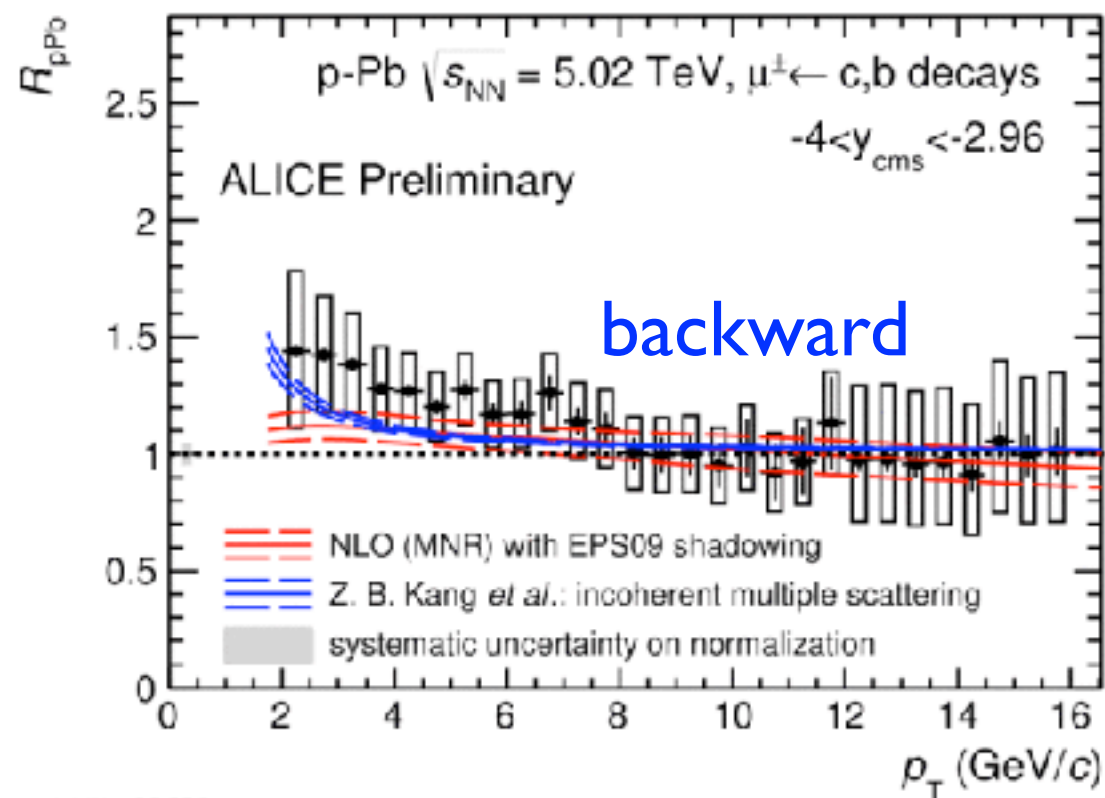
CMS b-jet R_{pA} in bins of transverse momentum and pseudo-rapidity

PYTHIA R_{pA} consistent with unity as a function of p_T and pseudo-rapidity

CMS, PLB 754 (2016) 59

Heavy flavour leptons: LHC vs. RHIC

ALICE heavy flavour muons ($c, b \rightarrow \mu$) in pPb collisions at 5.02 TeV



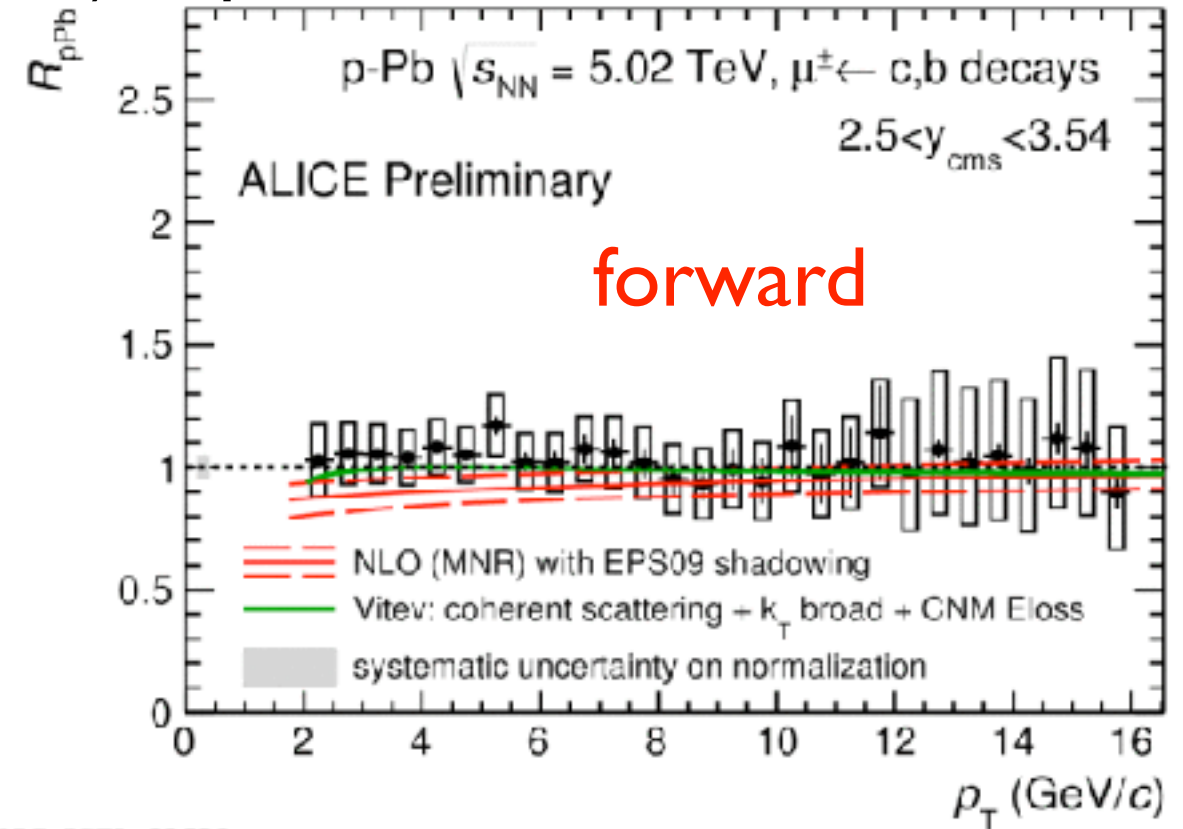
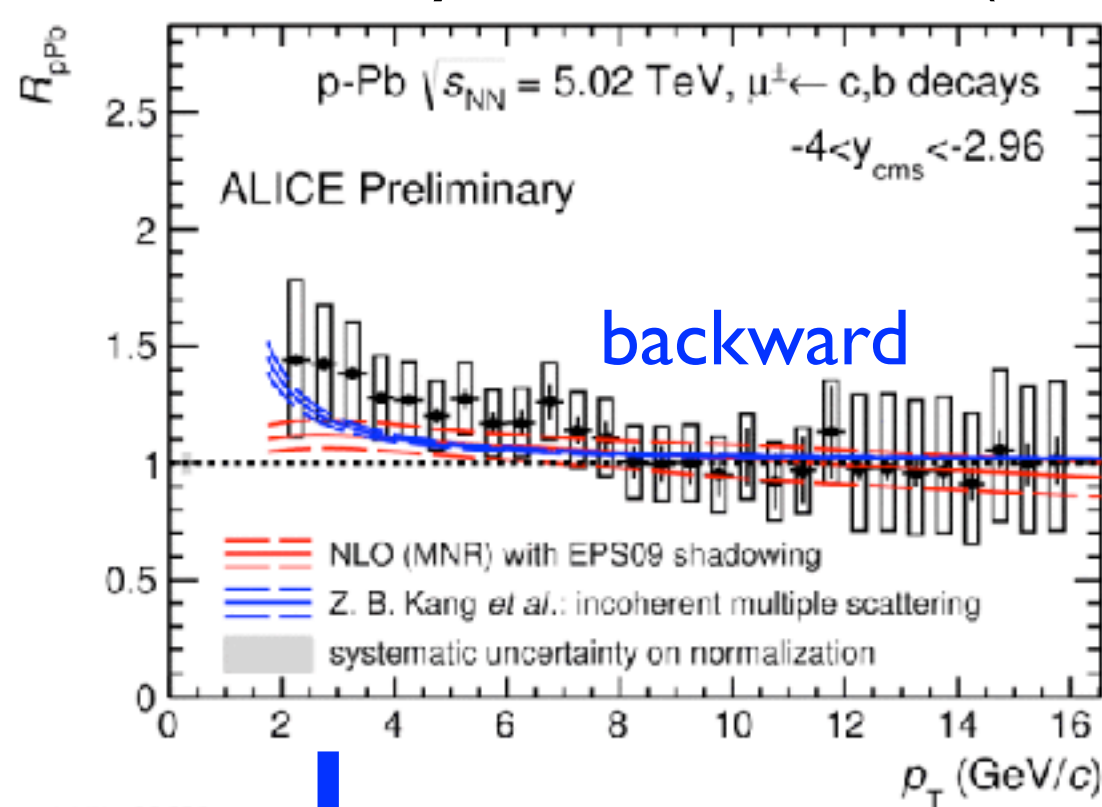
forward (shadowing)

backward (anti-shadowing)

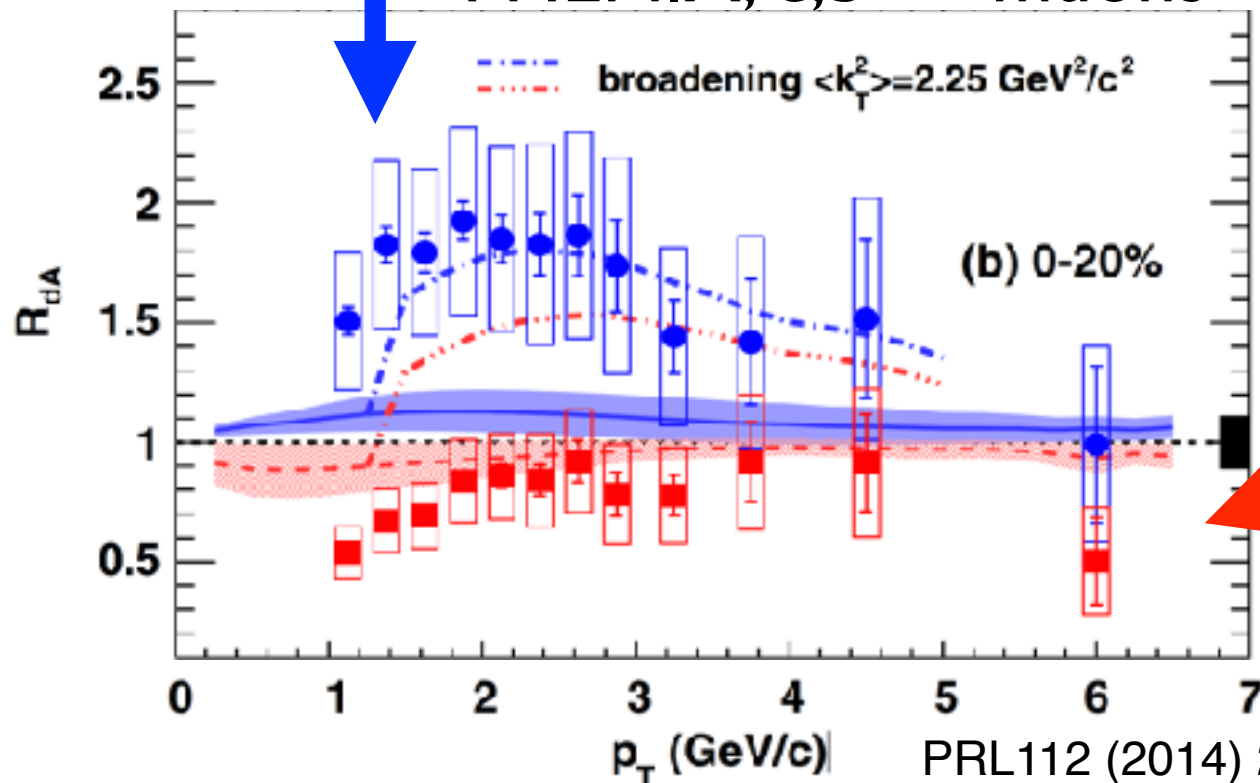
Models with CNM describe forward/backward rapidity at LHC

Heavy flavour leptons: LHC vs. RHIC

ALICE heavy flavour muons ($c, b \rightarrow \text{muons}$) in pPb collisions at 5.02 TeV



PHENIX, $c, b \rightarrow \text{muons}$



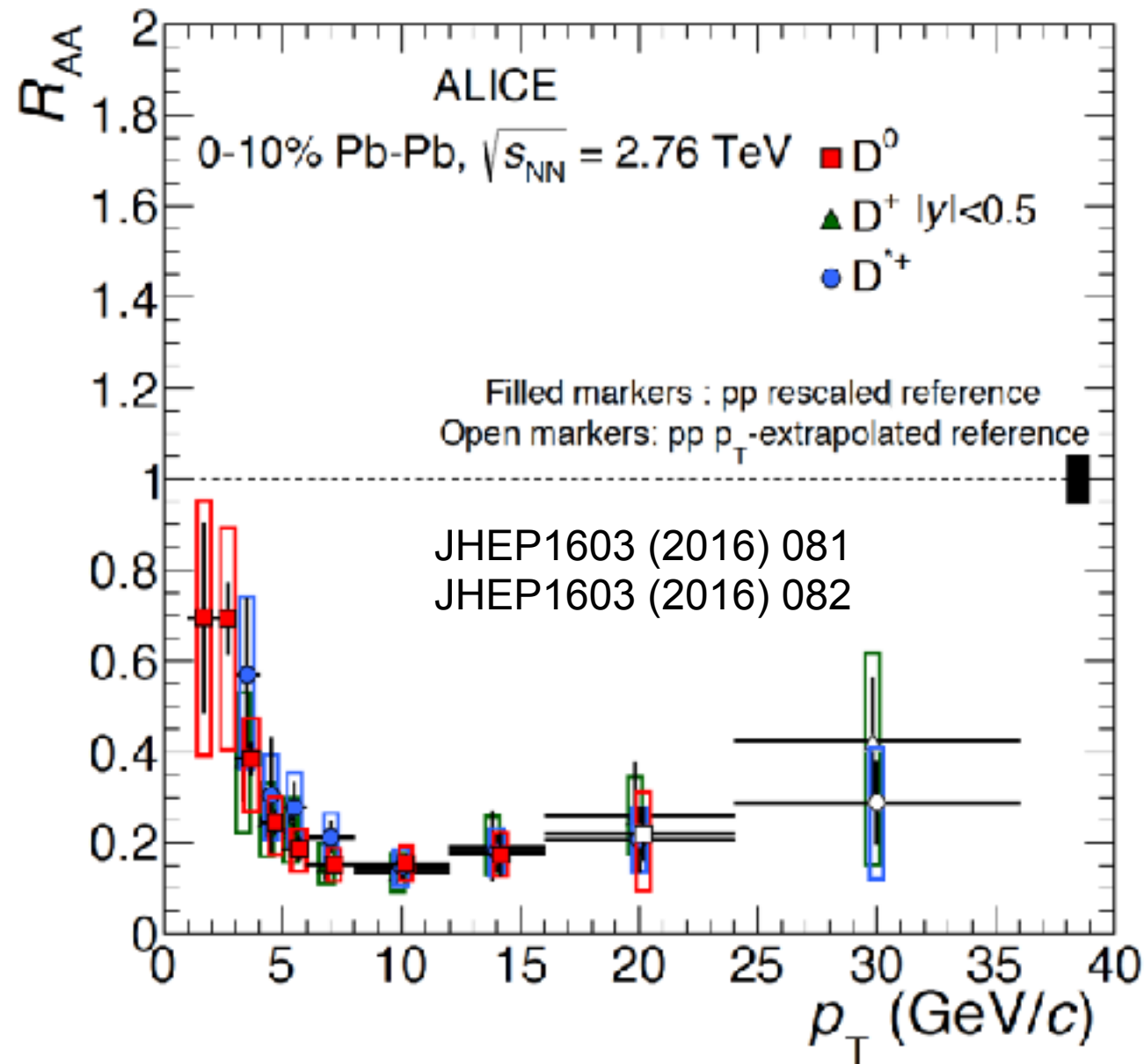
forward (shadowing)
 backward (anti-shadowing)

Models with CNM describe forward/backward rapidity at LHC
→ Not possible at RHIC

PRL112 (2014) 252301

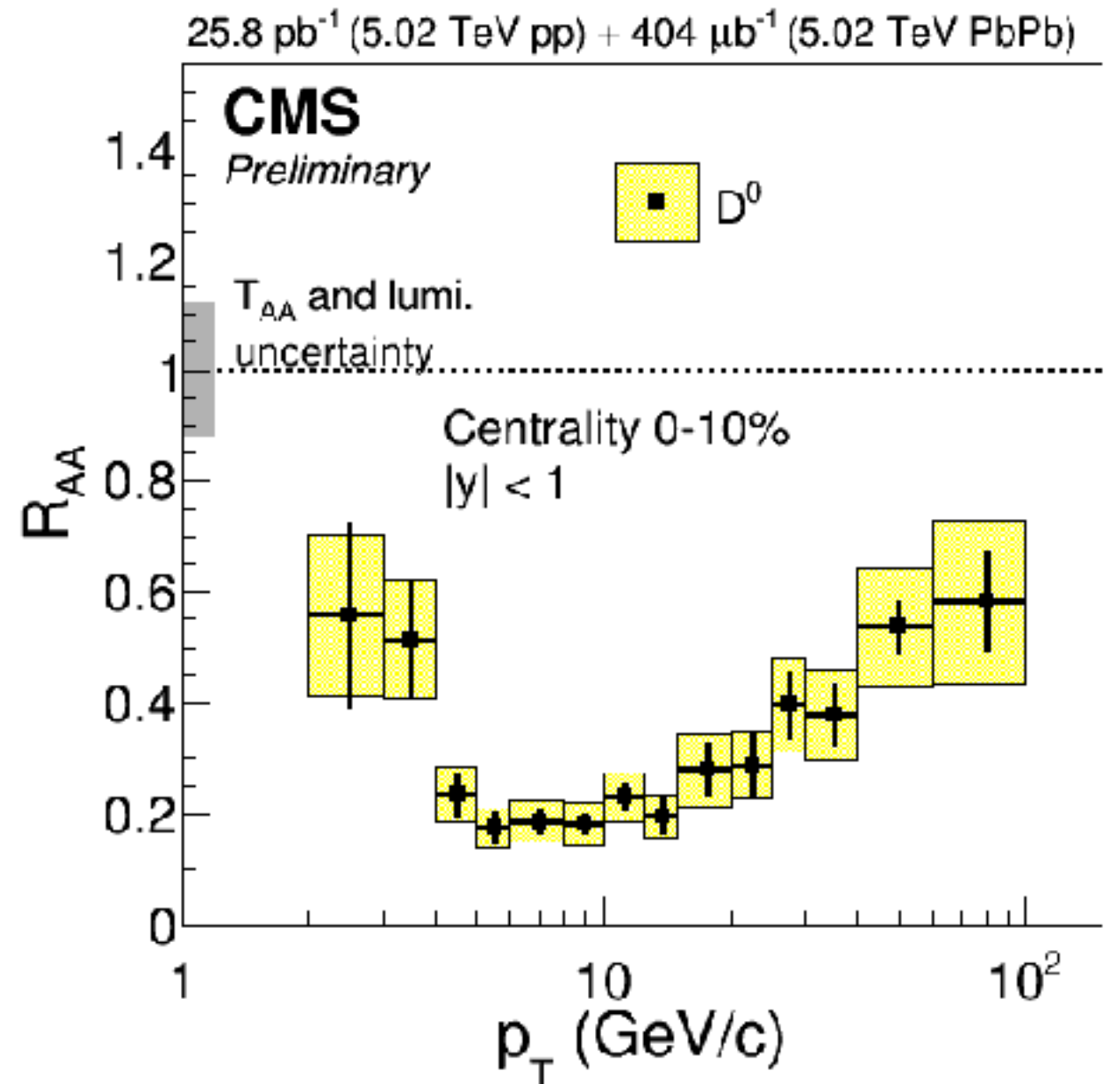
D meson R_{AA} in 0-10%

ALICE D^0 R_{AA} $|y| < 0.5$ at **2.76 TeV**



Strong suppression at 2.76 TeV:
same suppression for D^0, D^+, D^{+}*

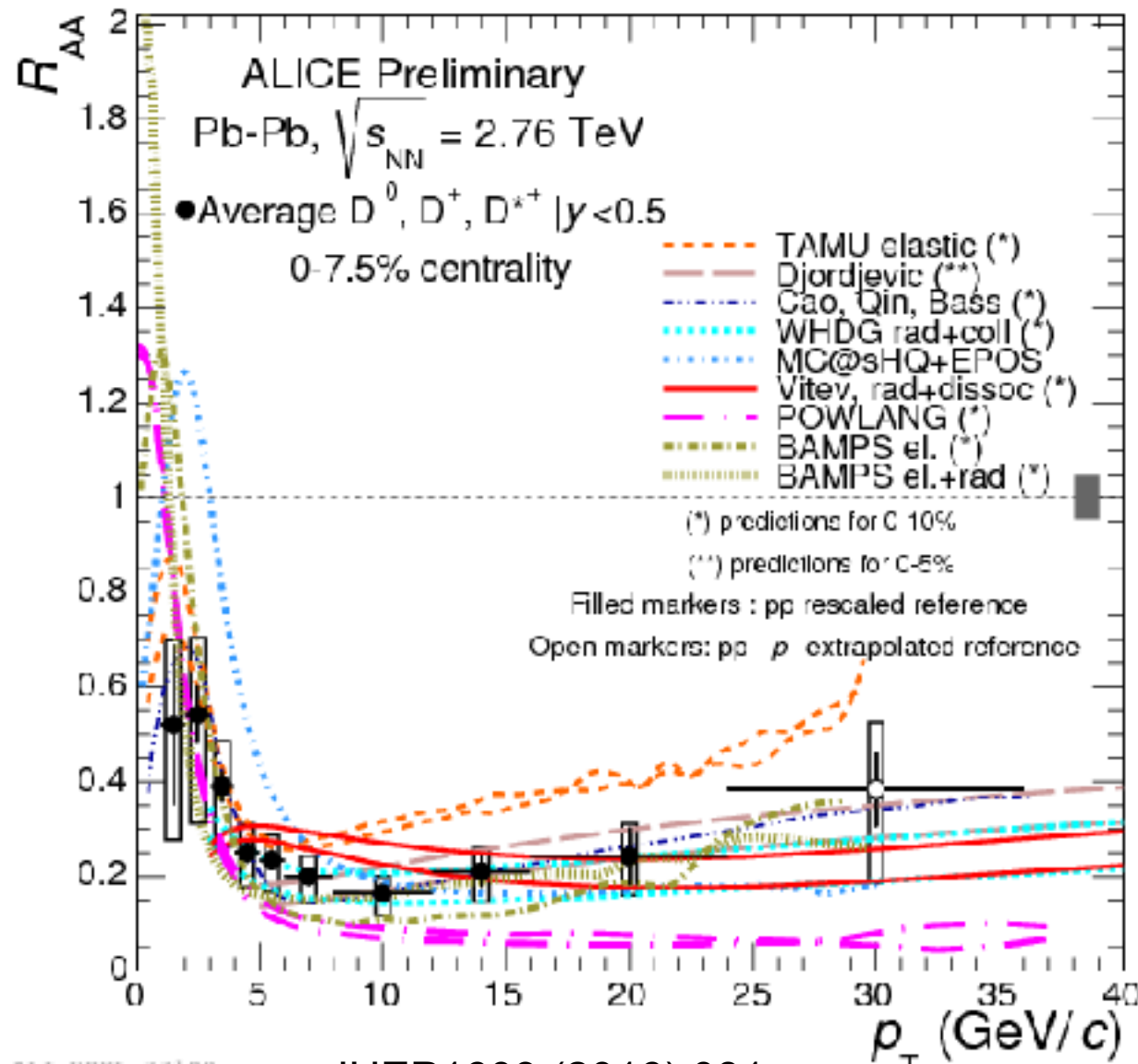
CMS D^0 R_{AA} $|y| < 1.0$ at **5.02 TeV**



CMS-PAS-HIN-16-001

Similar suppression at 5.02 TeV:
Rising trend observed when going to high p_T

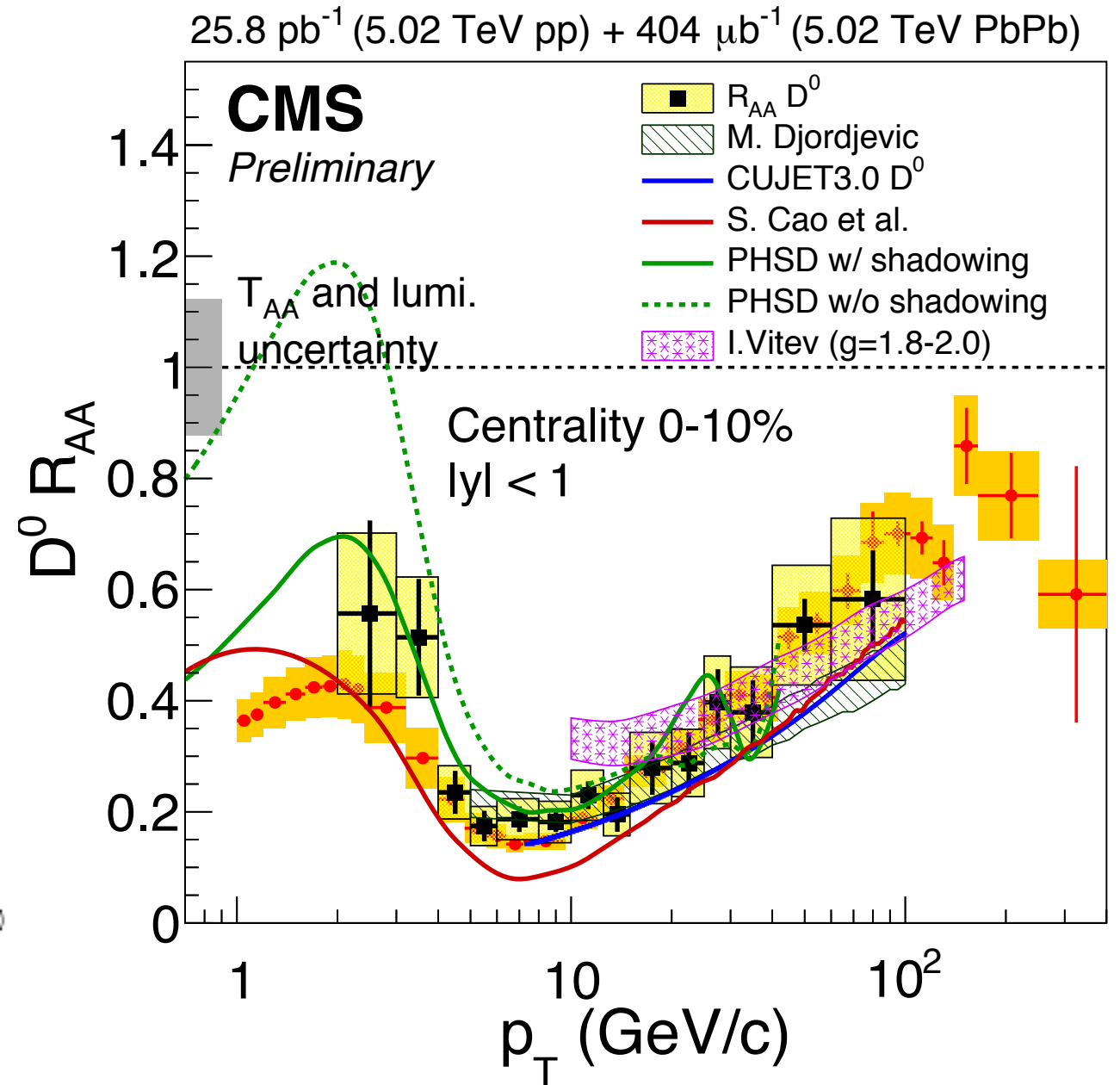
Comparison to theoretical calculations



ALI-PRKL-77100

JHEP1603 (2016) 081

JHEP1603 (2016) 082

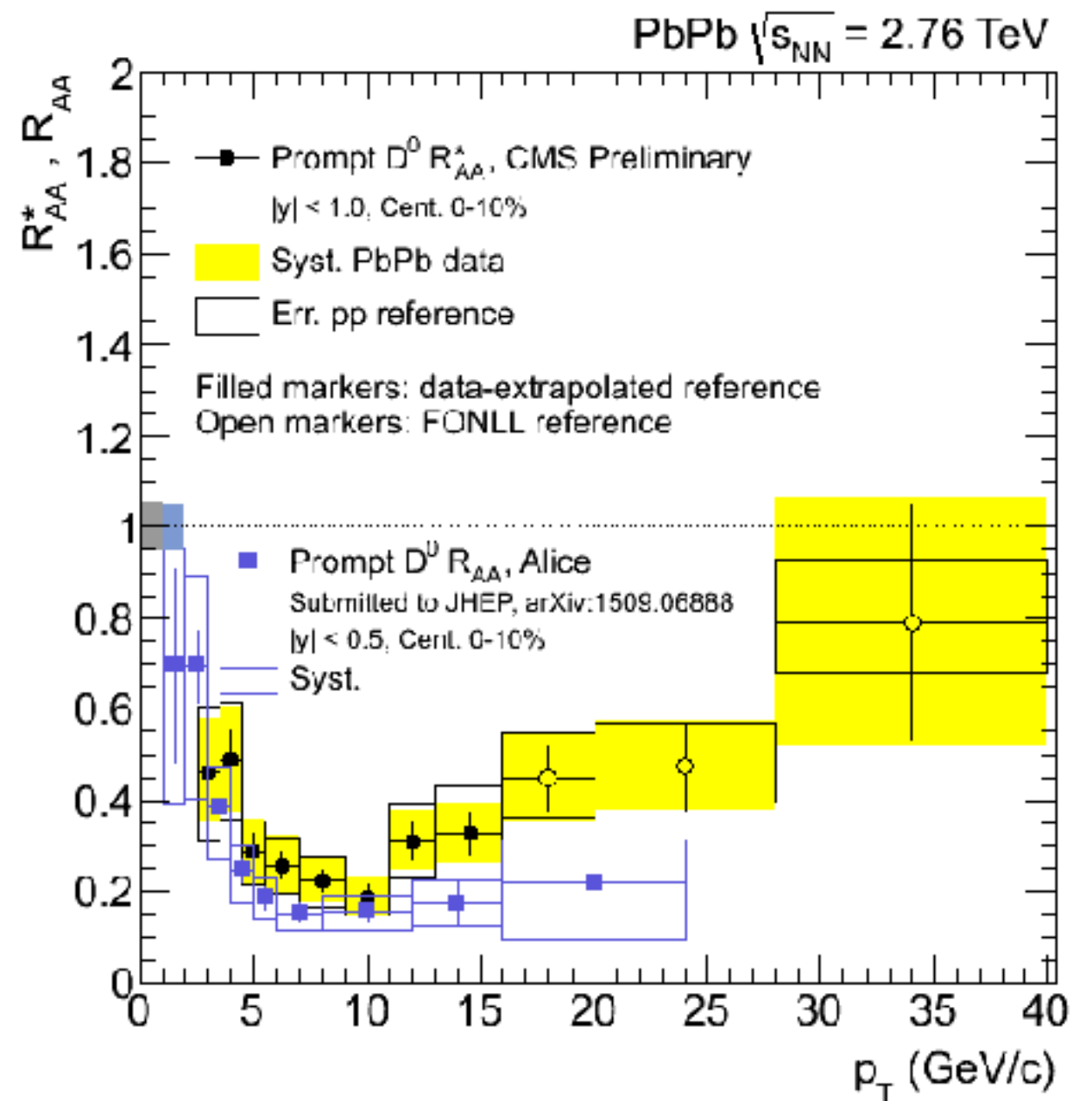
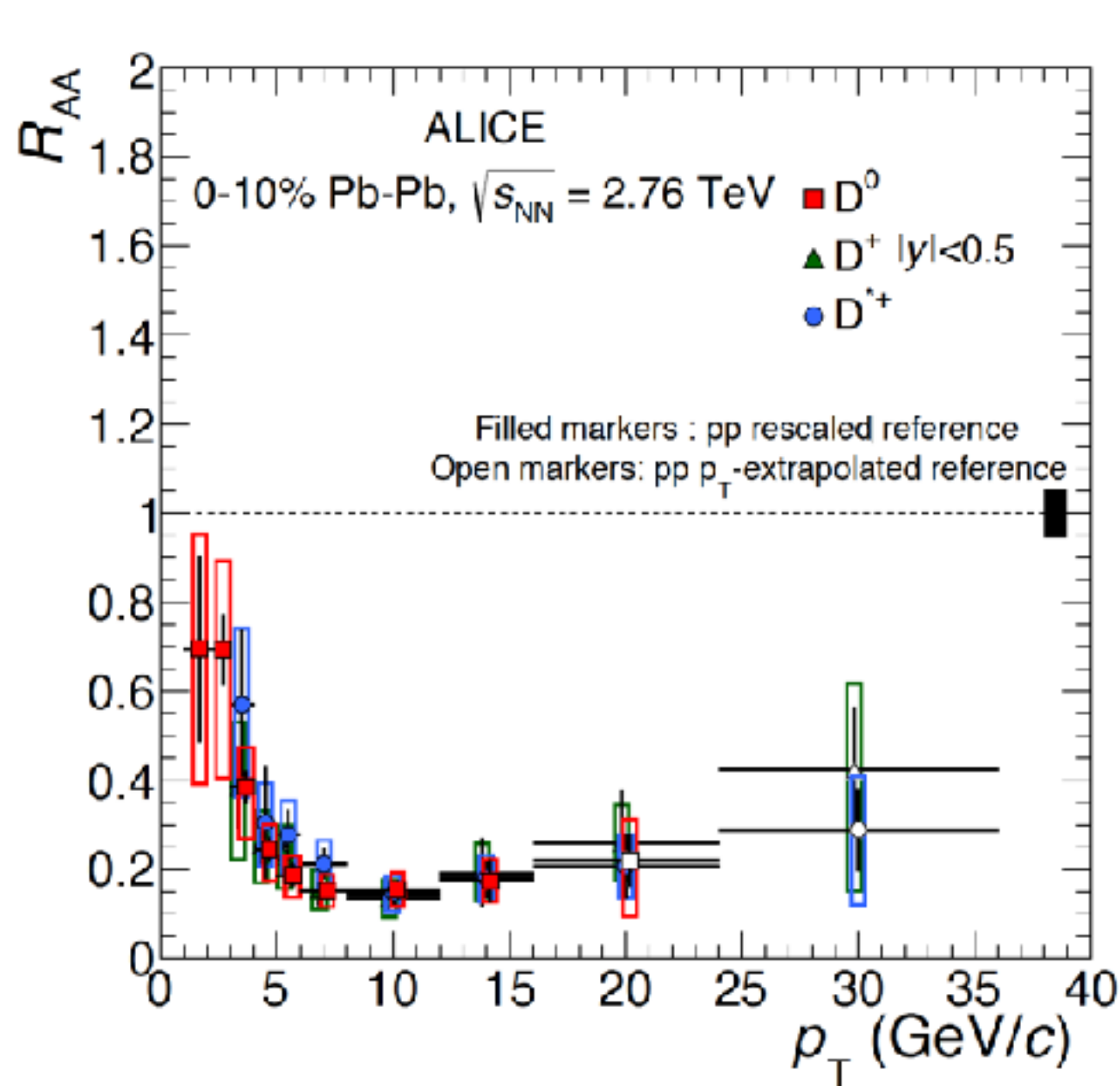


CMS-PAS-HIN-16-001

Several models describe the data within uncertainties:

- hints at low p_T that collisional energy loss is non negligible
- pure collisional models can describe the R_{AA} up to high p_T (??)
- shadowing improve description of the data at low p_T

D meson R_{AA} at 2.76 TeV



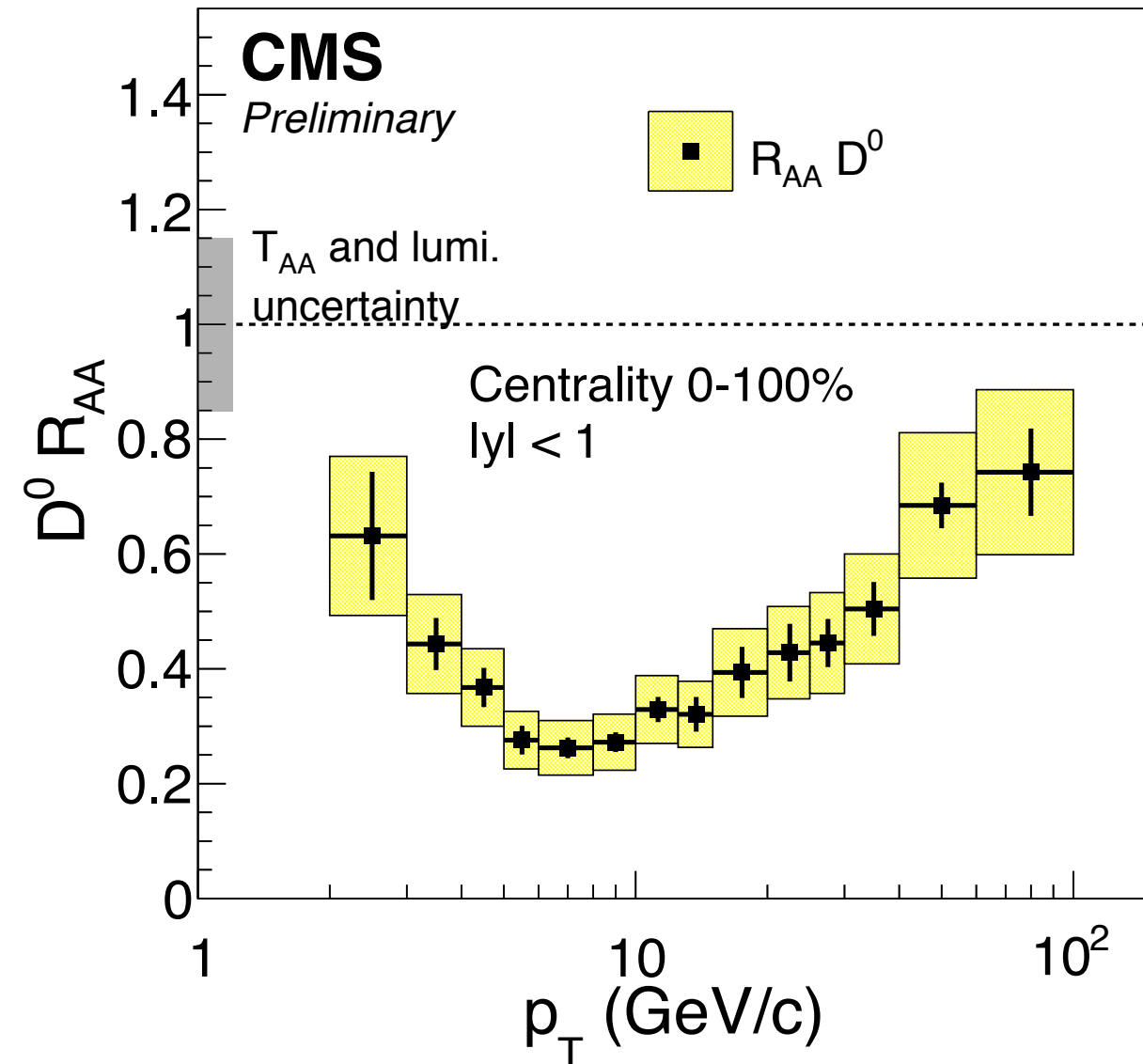
Strong suppression in central PbPb events:
 same suppression for D^0, D^+, D^{*+} indicate
 independence from fragmentation

ALICE and **CMS** in good agreement
 Differences at higher p_T due to different
 pp references

D^0 meson R_{AA} at 5.02 TeV

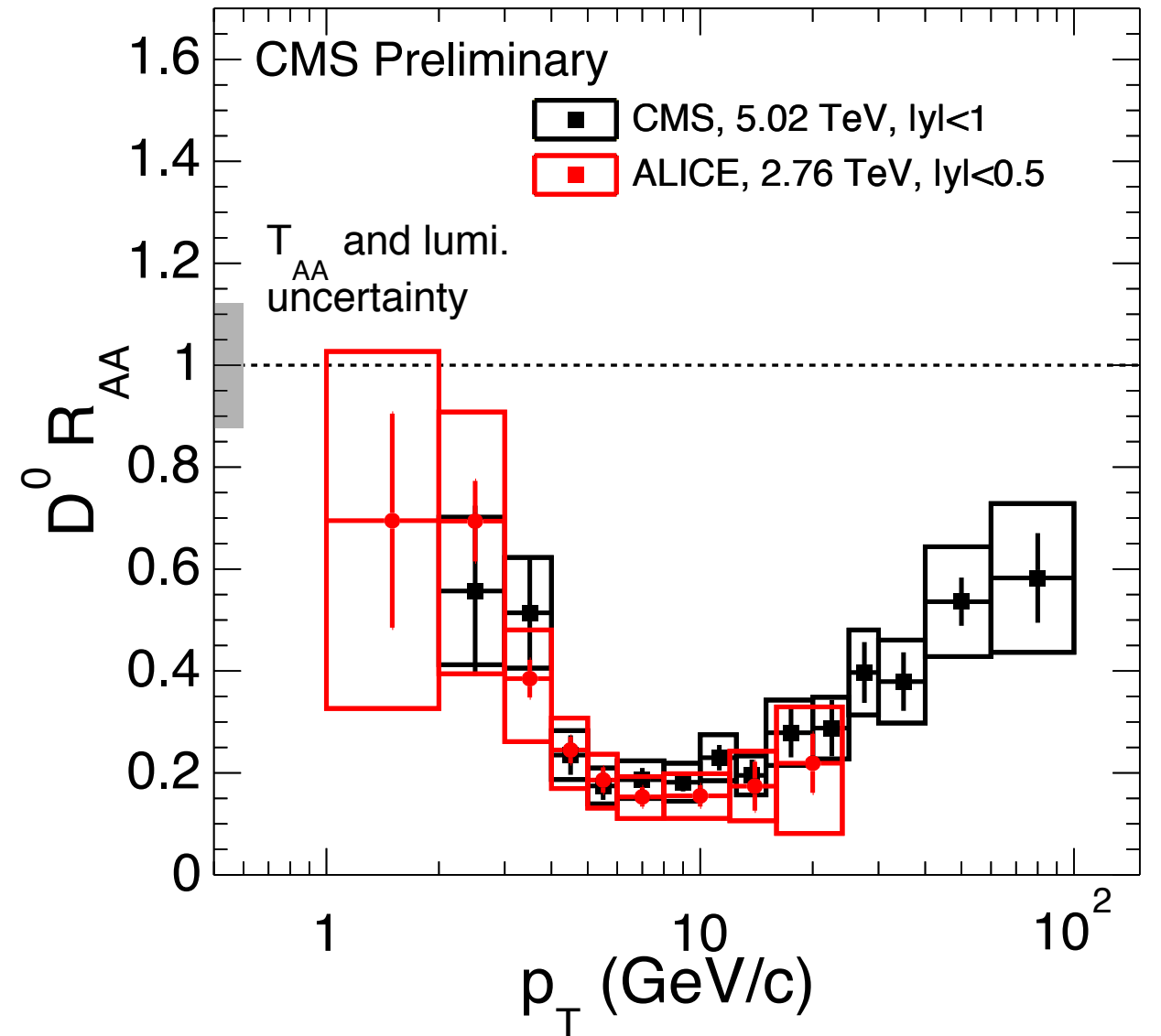
CMS D^0 R_{AA} $|y| < 1.0$ at 5.02 TeV

25.8 pb⁻¹ (5.02 TeV pp) + 404 μ b⁻¹ (5.02 TeV PbPb)



Strong suppression observed at 5.02 TeV
Rising trend observed when going to high p_T

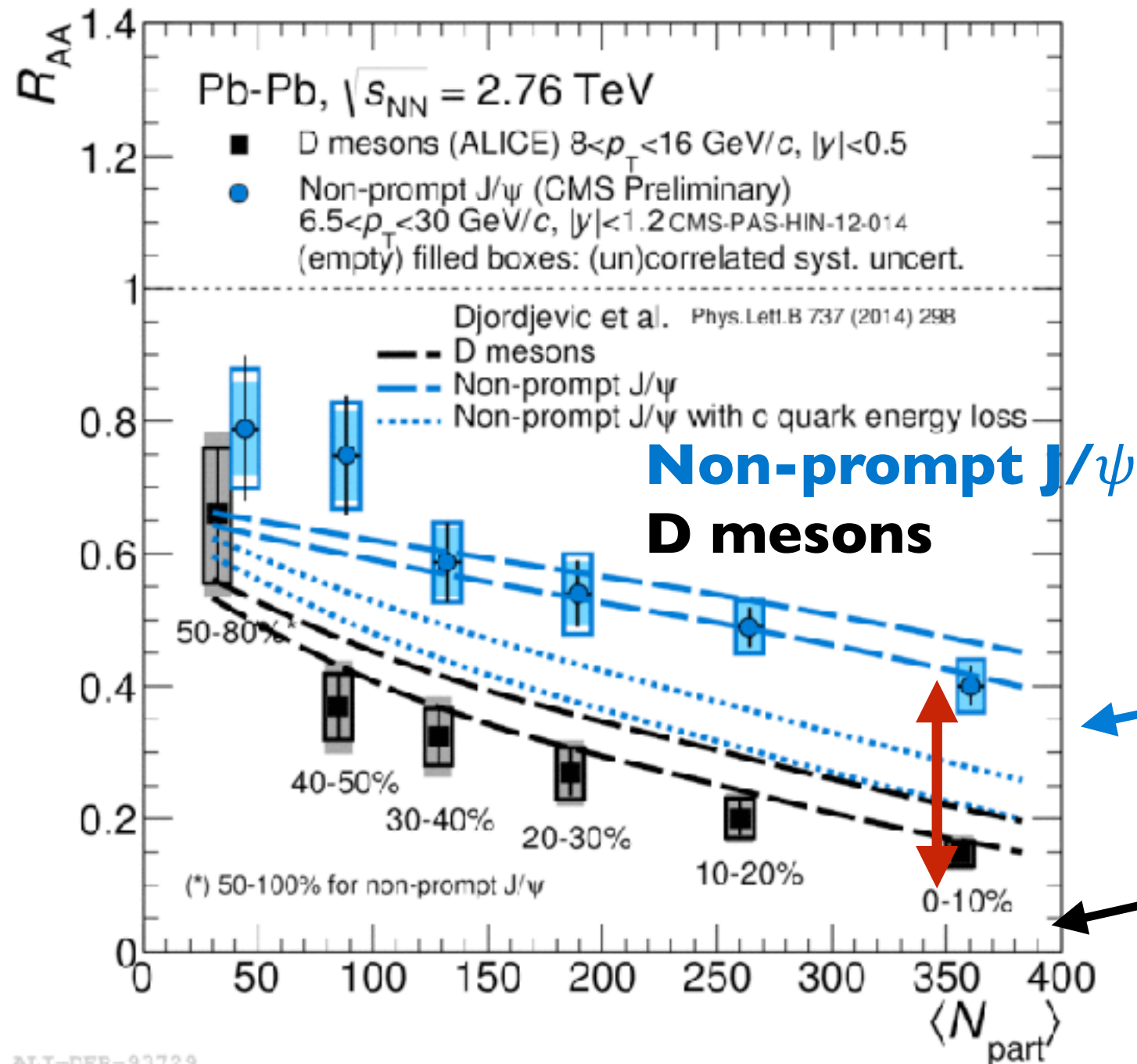
25.8 pb⁻¹ (5.02 TeV pp) + 404 μ b⁻¹ (5.02 TeV PbPb)



Similar suppression observed at 2.76 and 5.02 TeV by CMS and ALICE
Caveat: different rapidities

Flavour dependence of E_{loss} at 2.76 TeV

ALICE, JHEP 1511 (2015) 205



pQCD model (M.Djordjevic) that assumes two different mass hypotheses for non prompt J/ ψ

M.Djordjevic, PRL 112, 042302 (2014)

b-quark E_{loss}

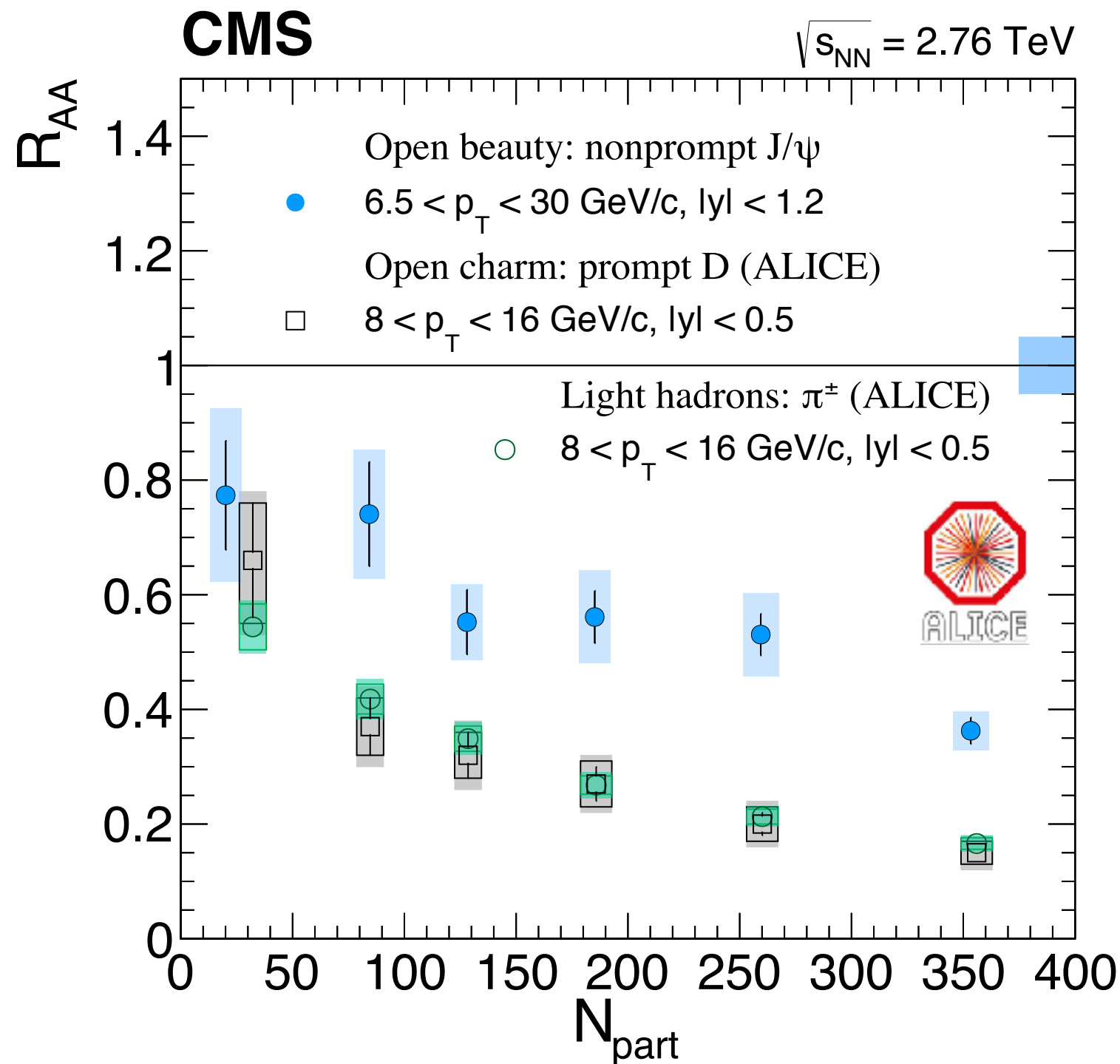
c-quark E_{loss}

ALI-DEP-93729

According to this model, the difference R_{AA} for non prompt J/ ψ and B can be attributed to a difference in the E_{loss} of charm and beauty quarks

Flavour dependence of E_{loss} at 2.76 TeV

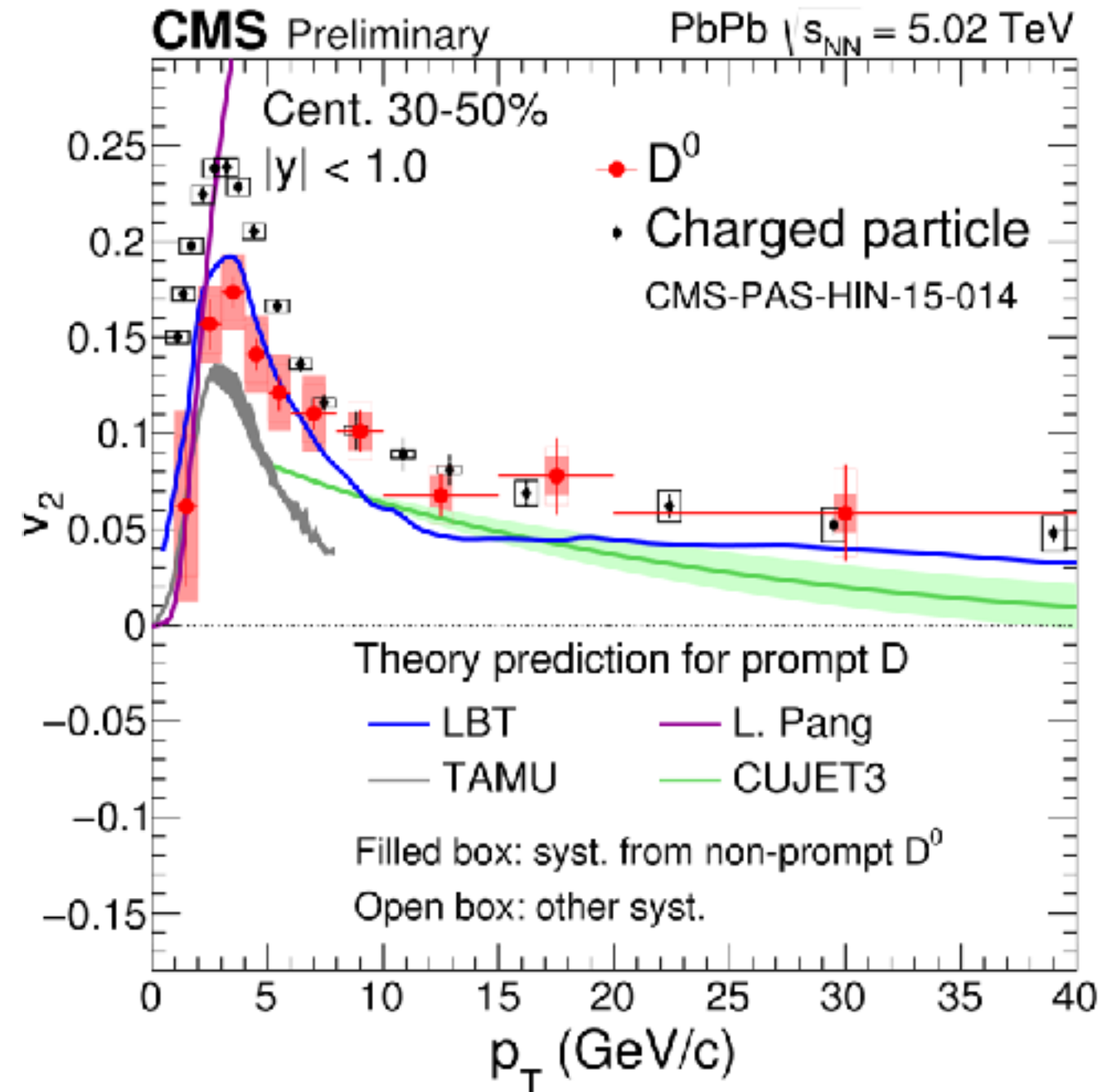
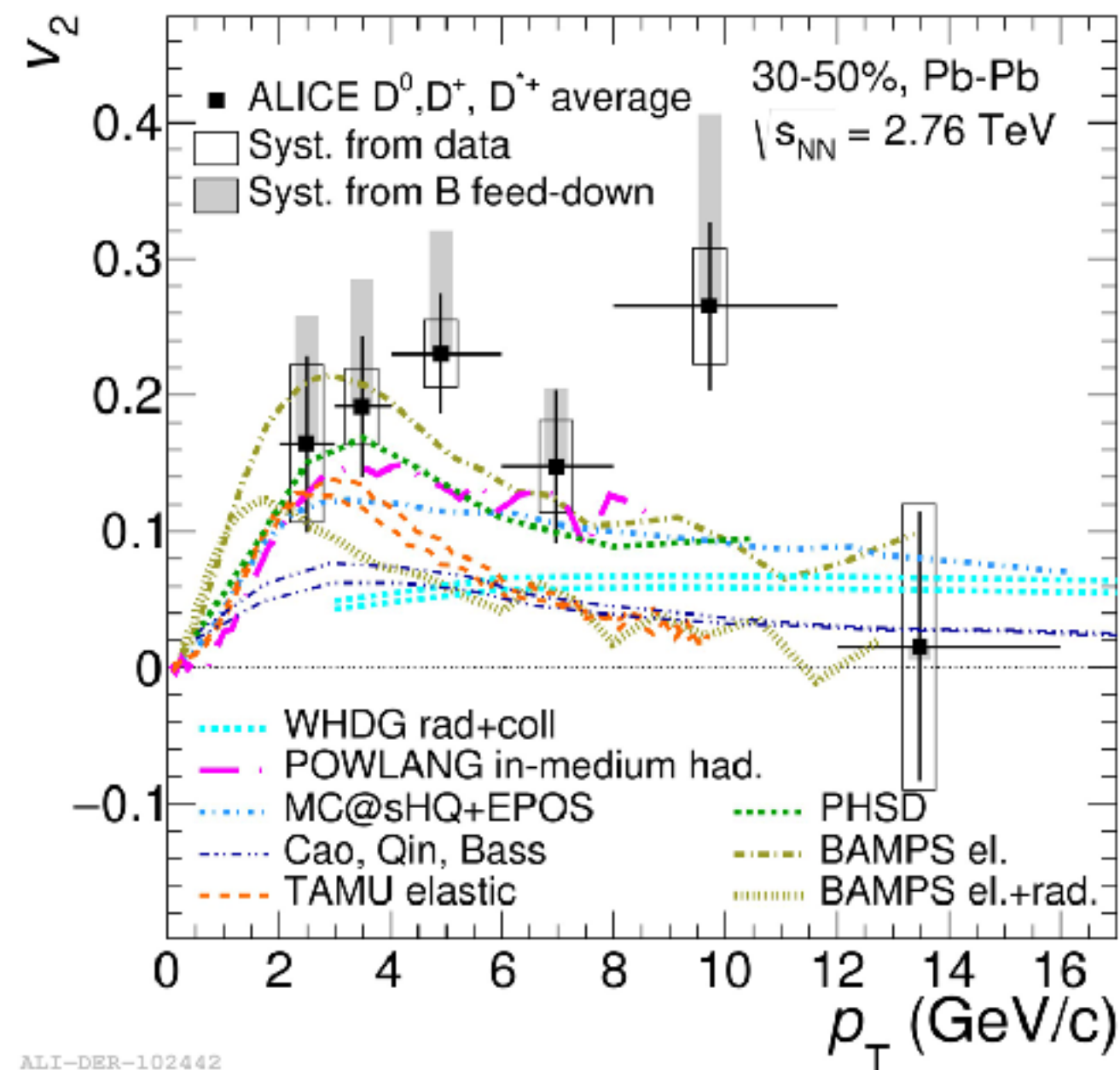
CMS-PAS-HIN-15-005



Non-prompt J/ψ
D mesons
 π^{+-}

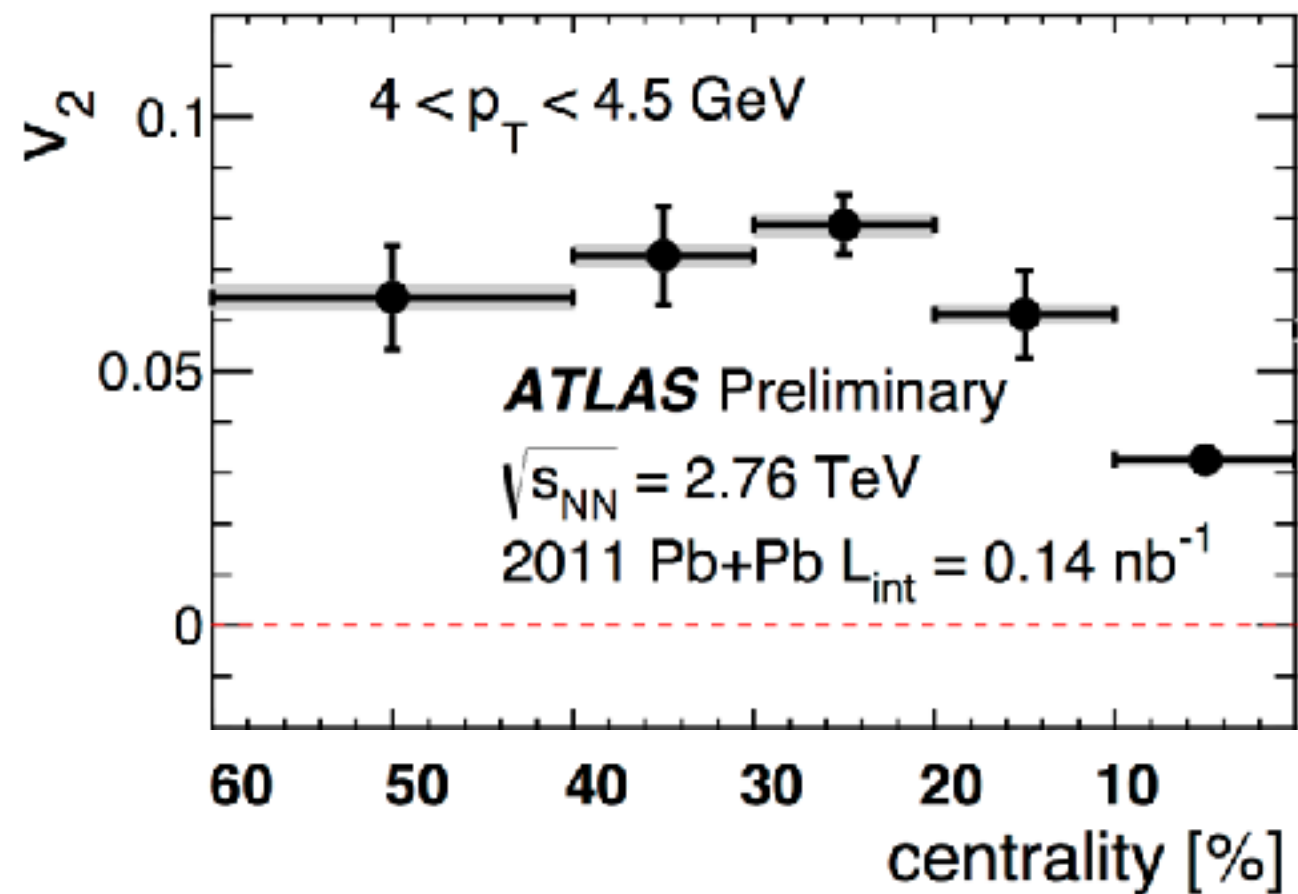
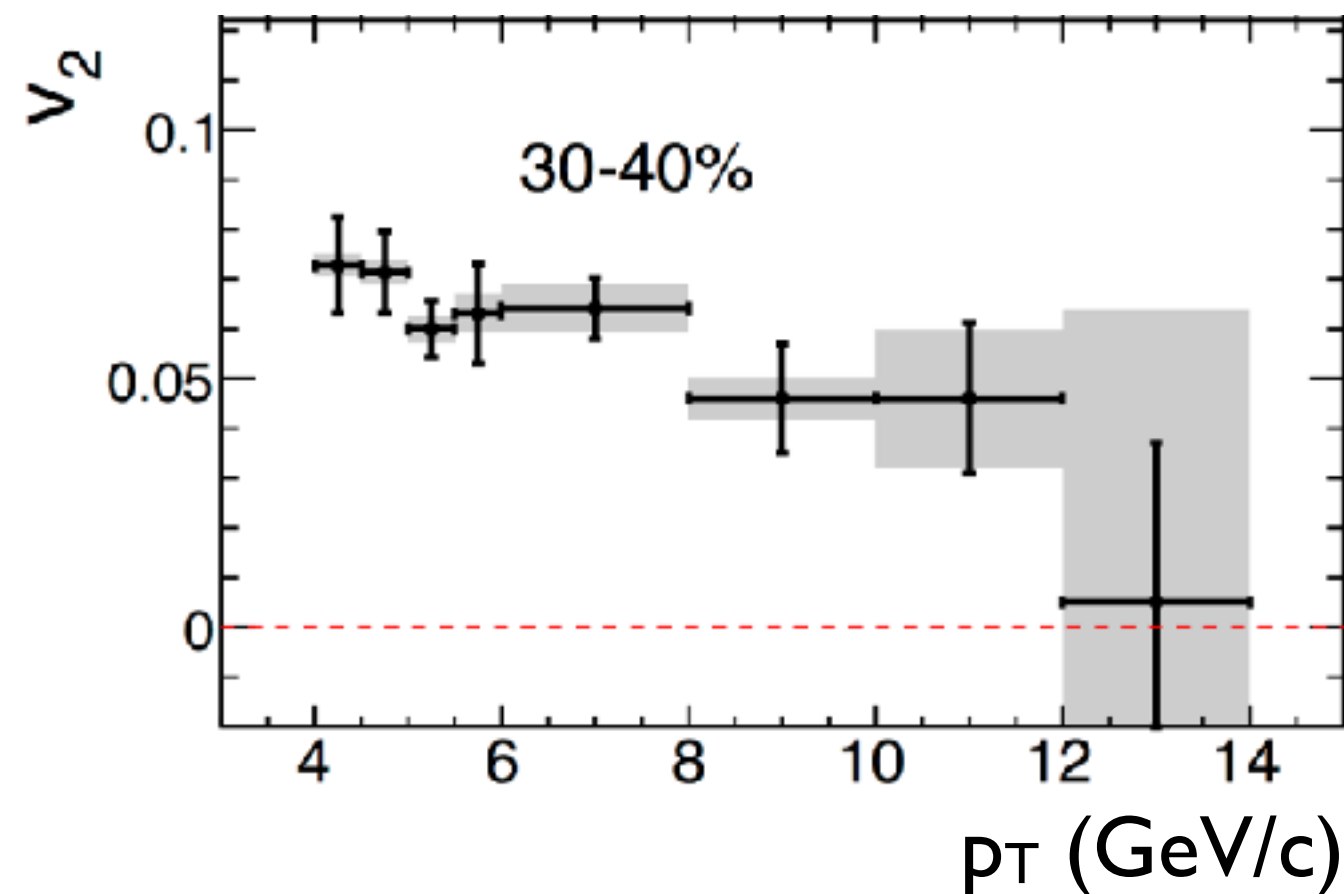
No change in the physics message when comparing to the final result of non prompt J/ψ R_{AA} from CMS

Comparison with models



we need **charm quark diffusion** to describe the magnitude of the D meson v_2 !

Heavy-flavour muons at 2.76 TeV

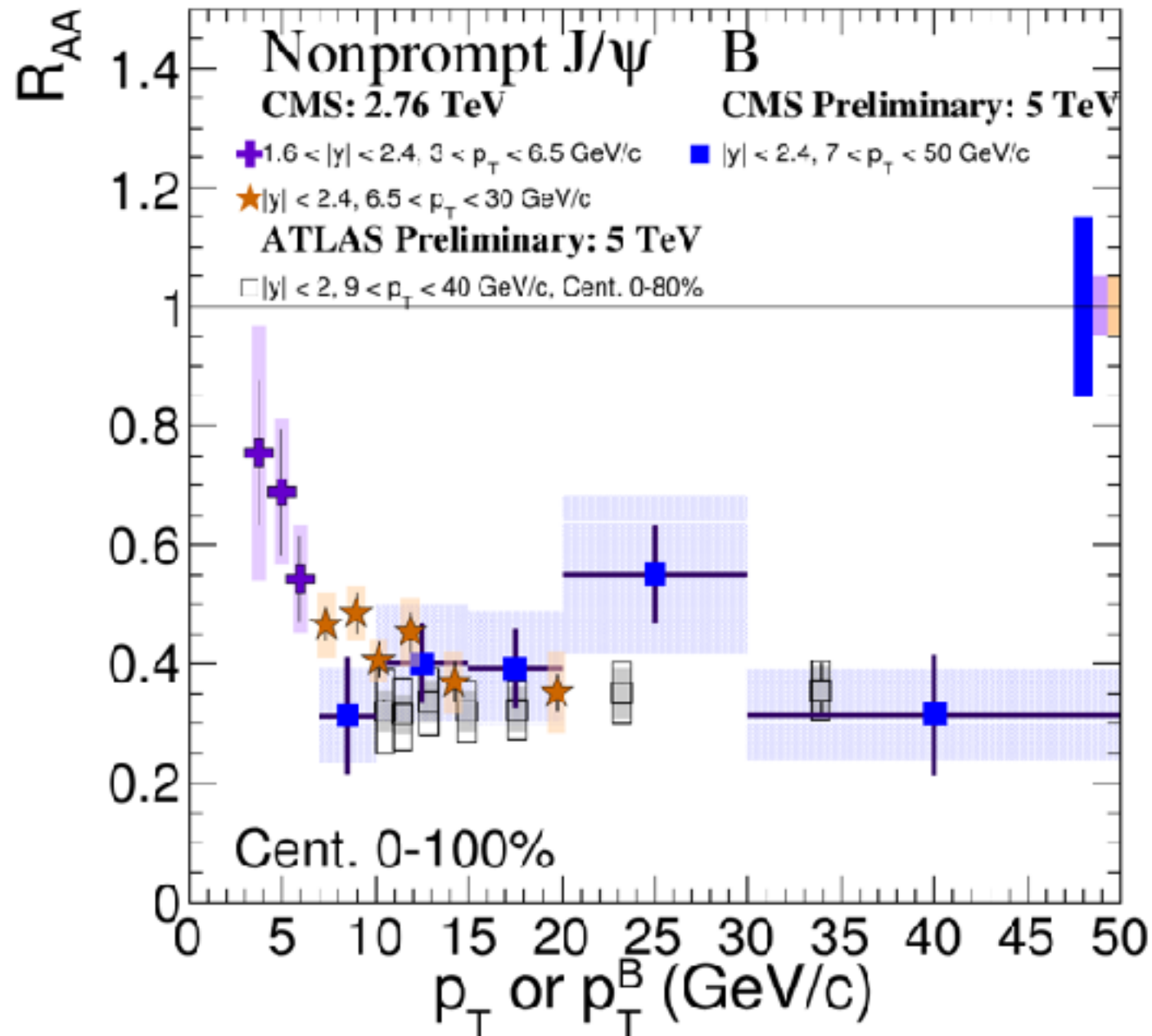


ATLAS-CONF-2015-053

Positive v_2 for muons from heavy-flavour decays (b+c) at LHC:

- include the contributions of beauty to v_2 that is currently unknown
- **v_2 of heavy flavour muons $< v_2$ (D^0) from ALICE**

Non-prompt J/ψ at 2.76 TeV vs B^+ at 5.02 TeV

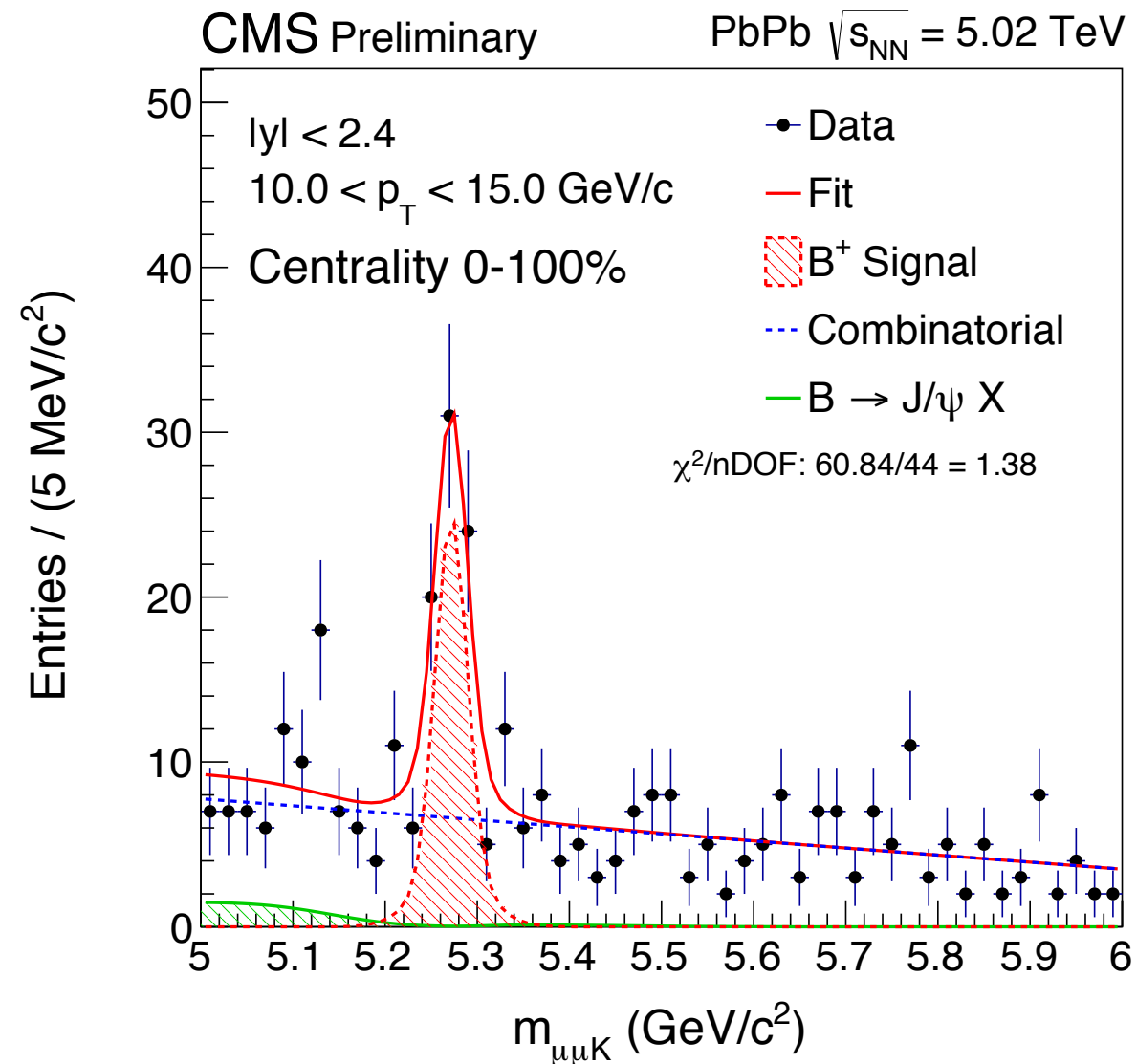
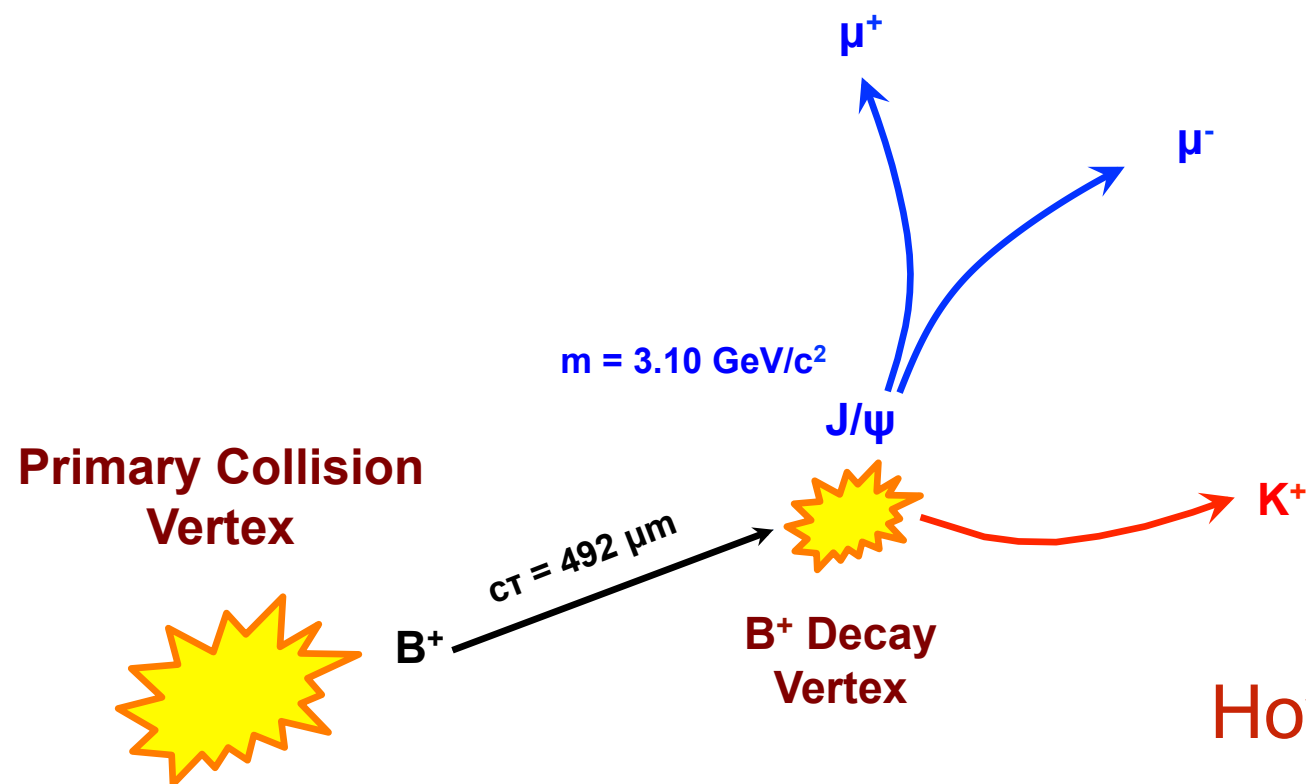


The B^+ R_{AA} at 5.02 TeV and non-prompt J/ψ at 2.76 fully compatible within uncertainties!
BIG CAVEAT: different energies!

Exclusive B-meson measurements

Exclusive B mesons can span the full range and get closer to the parton kinematic!

$$B^+ \rightarrow J/\psi K^+ \rightarrow \mu^+ \mu^- K^+$$

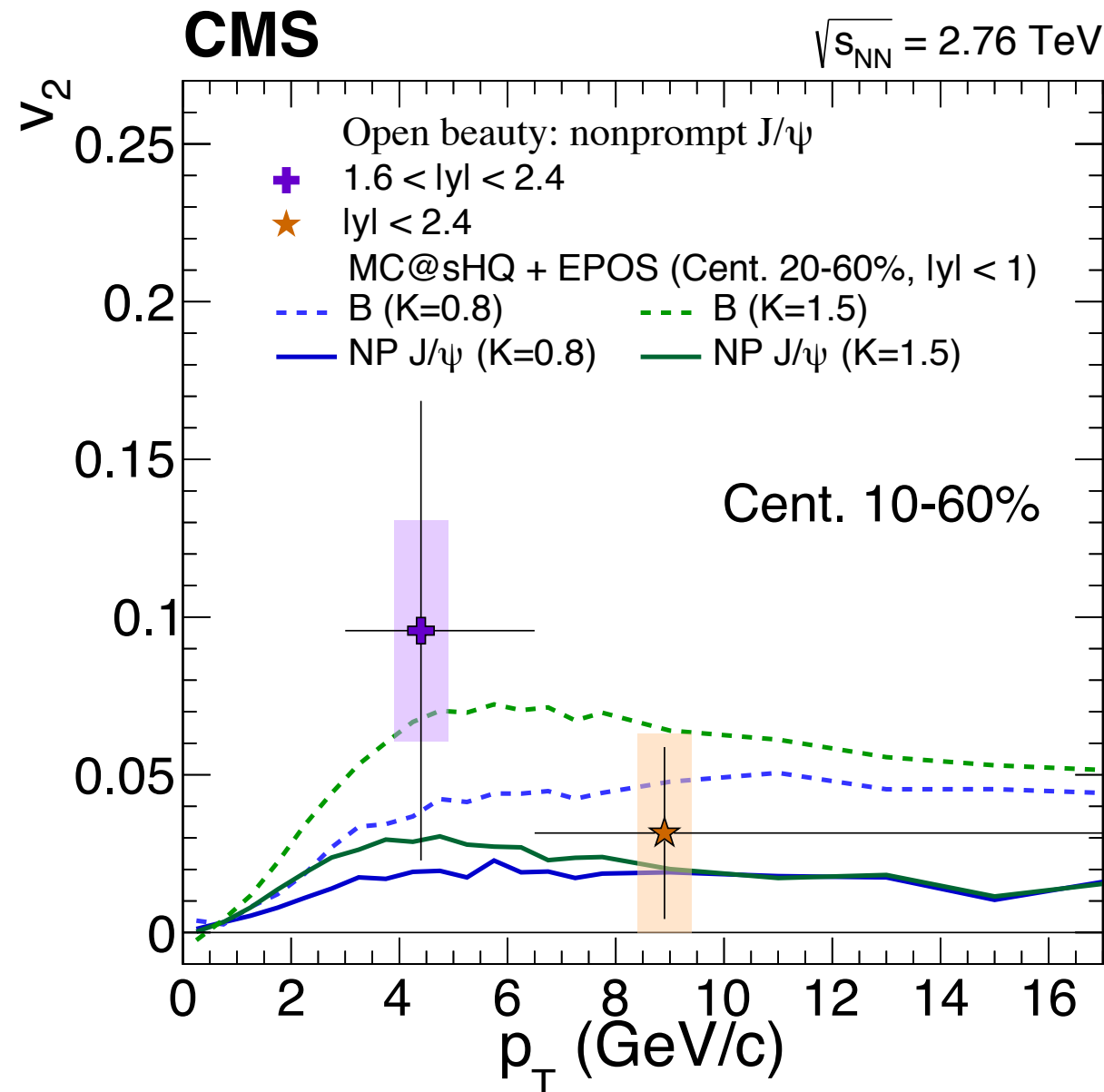


How do we reconstruct B mesons in CMS?

→ Clean and high statistics sample collected by triggering on muons!

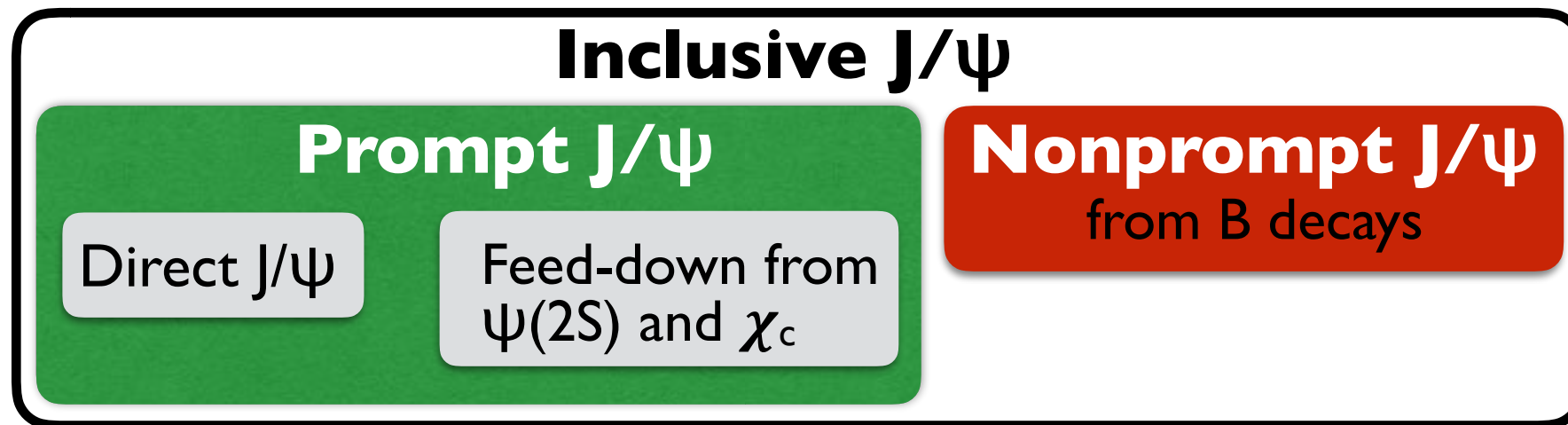
v_2 of non prompt J/ψ

v_2 of non prompt J/ψ in PbPb collisions at 2.76 TeV



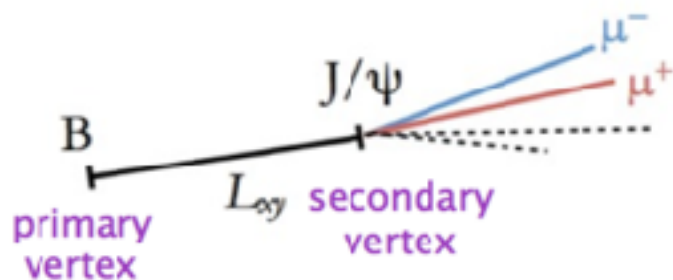
→ Compatible within uncertainties with theoretical calculations
Looking to see the new measurement with Run2 data with higher statistics!

prompt and non prompt charmonia



- **Prompt component:**
affected by color screening and regeneration in the QGP
- **Nonprompt component:**
reflects E_{loss} of b quarks in the medium

Separation of components based on **pseudo-proper decay length** ($\ell_{J/\psi}$):



$$\ell_{J/\psi}^{3D} = L_{xyz} \cdot \frac{m_{J/\psi}}{p_{\mu\mu}}$$

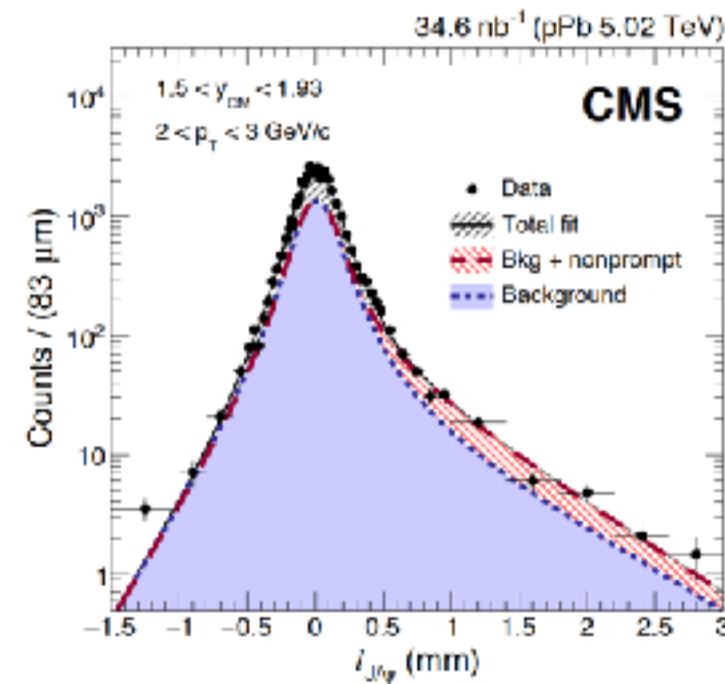
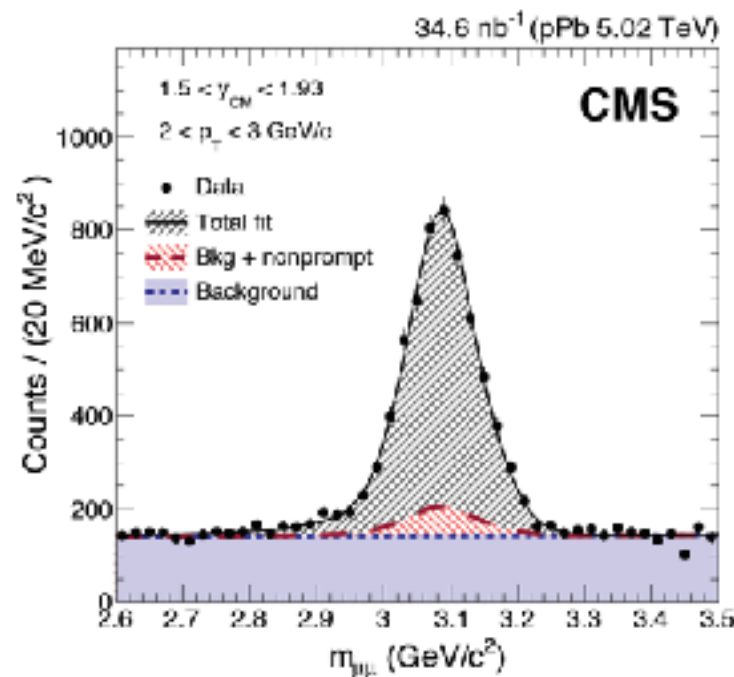
This talk:
prompt charmonia

Ta-Wei Wang's talk (8th):
nonprompt charmonia
& **full B reconstruction**

prompt and non prompt charmonia

Two techniques to separate components:

1. 2D fits of dimuon mass and pseudo-proper decay length



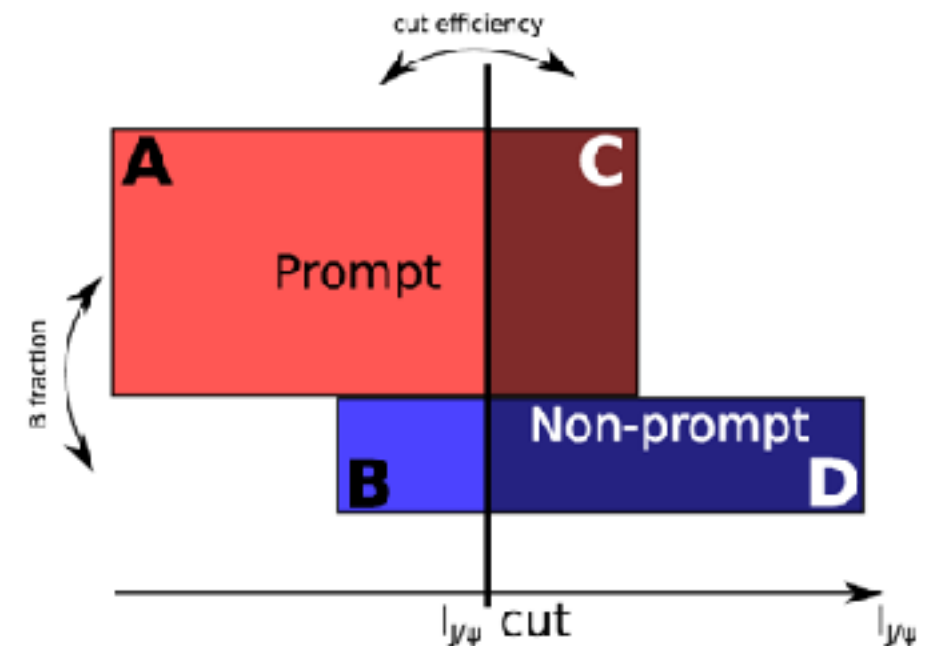
arXiv:1702.01462

2. Rejecting nonprompt using a cut on $l_{J/\psi}$

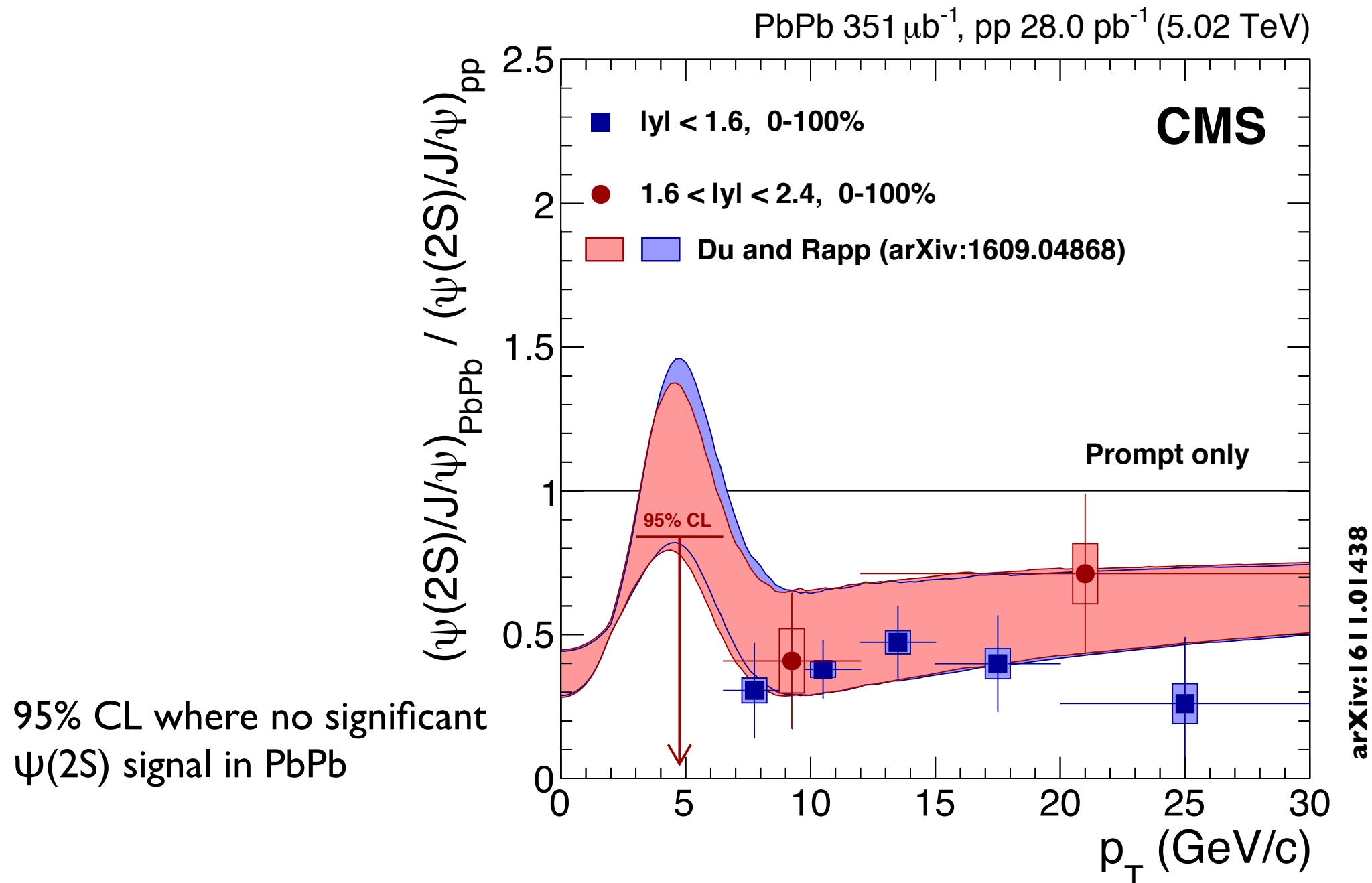
(can be used with low stats: $\psi(2S)$ analyses)

Correction (from data) to account for remaining nonprompt contamination:

- Using reverted $l_{J/\psi}$ cut
- MC efficiency of $l_{J/\psi}$ cut



$\psi(2S) / J/\psi$ vs p_T



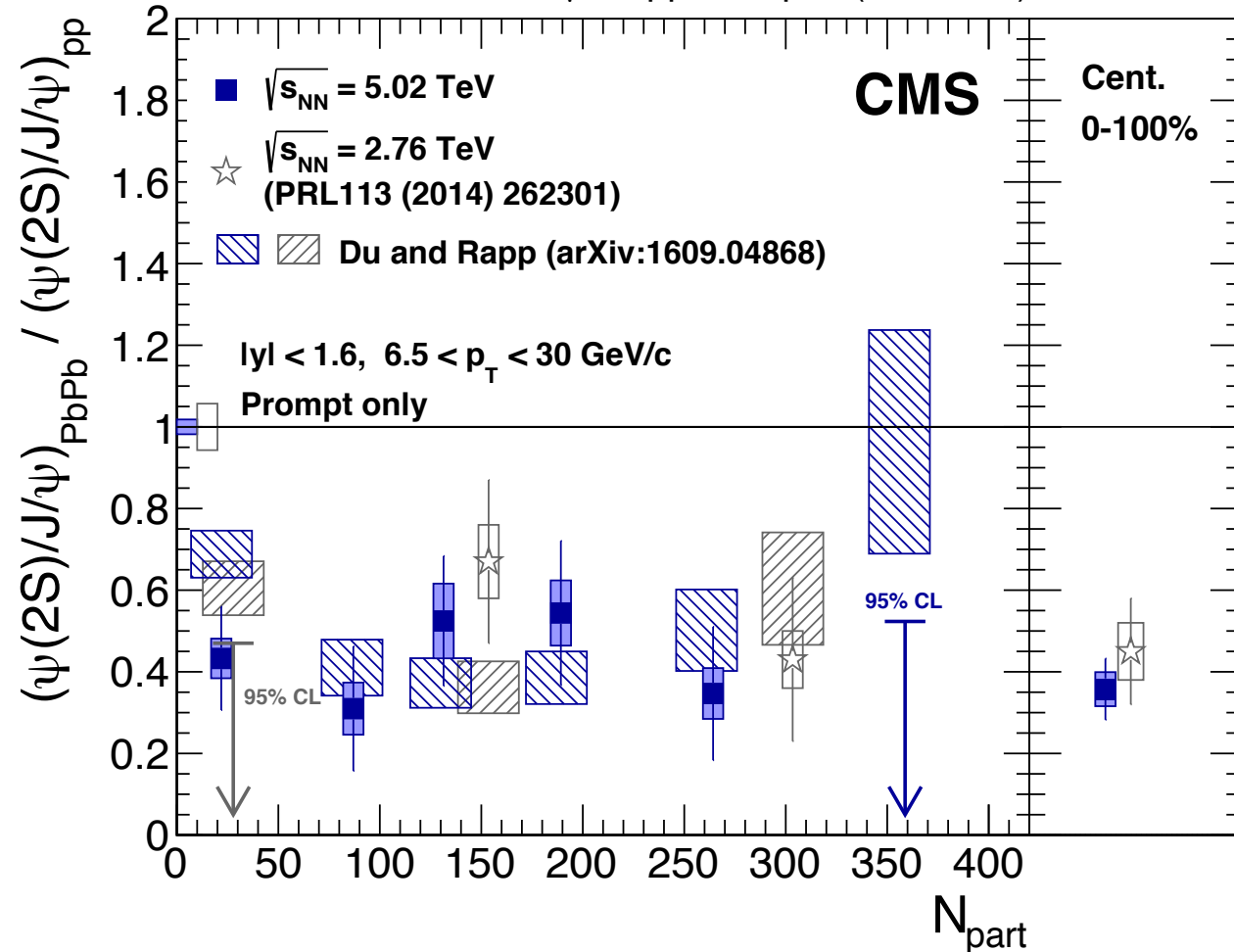
- $R_{AA}(\psi(2S))/R_{AA}(J/\psi) < 1$ in all bins \rightarrow **$\psi(2S)$ is more suppressed than J/ψ**
- No p_T dependence within uncertainties
- X. Du and R. Rapp: transport model with temperature dependent reaction rates
 \rightarrow $\psi(2S)$ regenerated later than J/ψ in the fireball evolution?

$\psi(2S) / J/\psi$ vs centrality

2.76 vs. 5.02 TeV

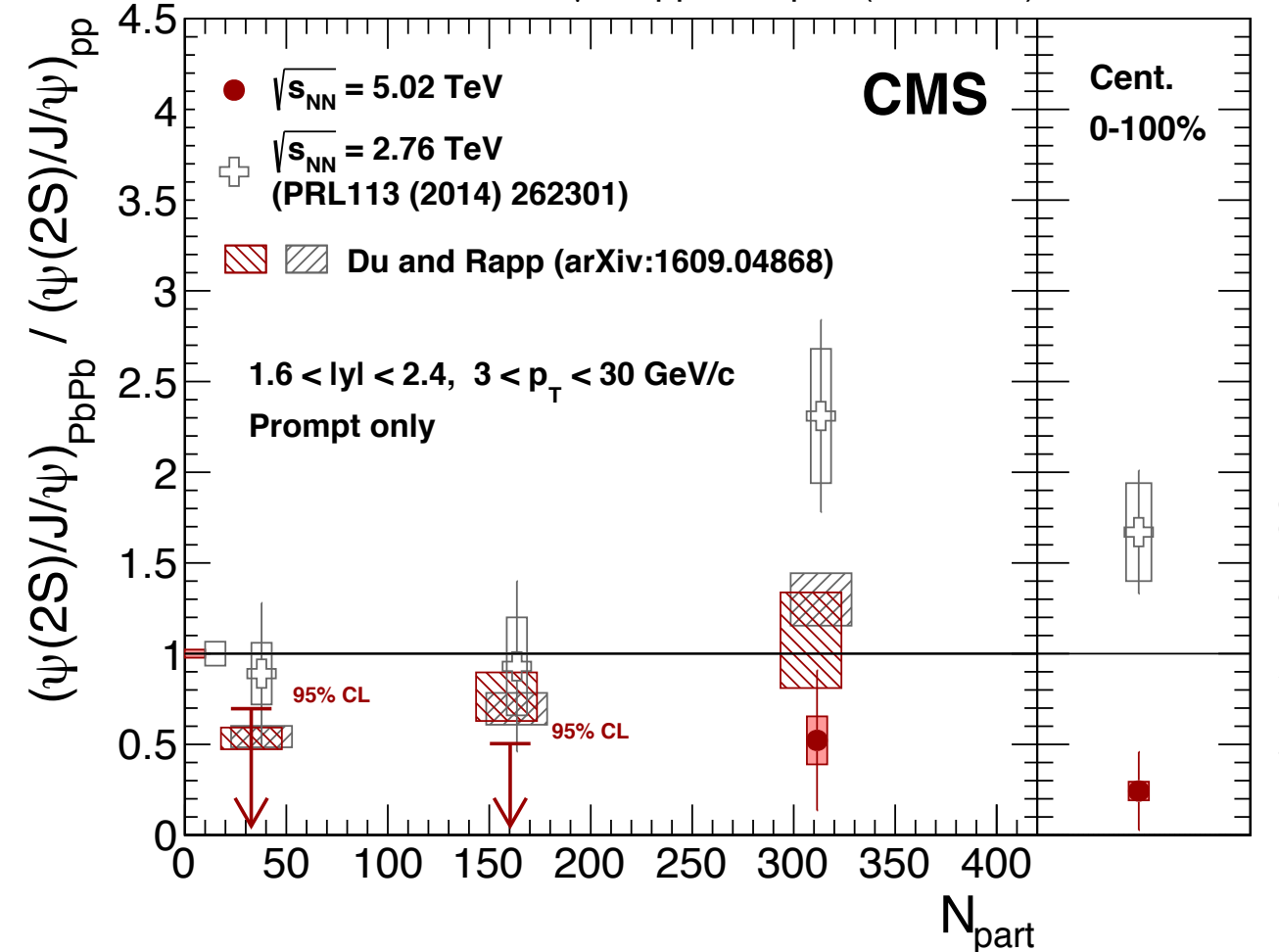
$|y| < 1.6 ; 6.5 < p_T < 10 \text{ GeV}/c$

PbPb $351 \mu\text{b}^{-1}$, pp 28.0 pb^{-1} (5.02 TeV)



$1.2 < |y| < 2.4 ; 3 < p_T < 30 \text{ GeV}/c$

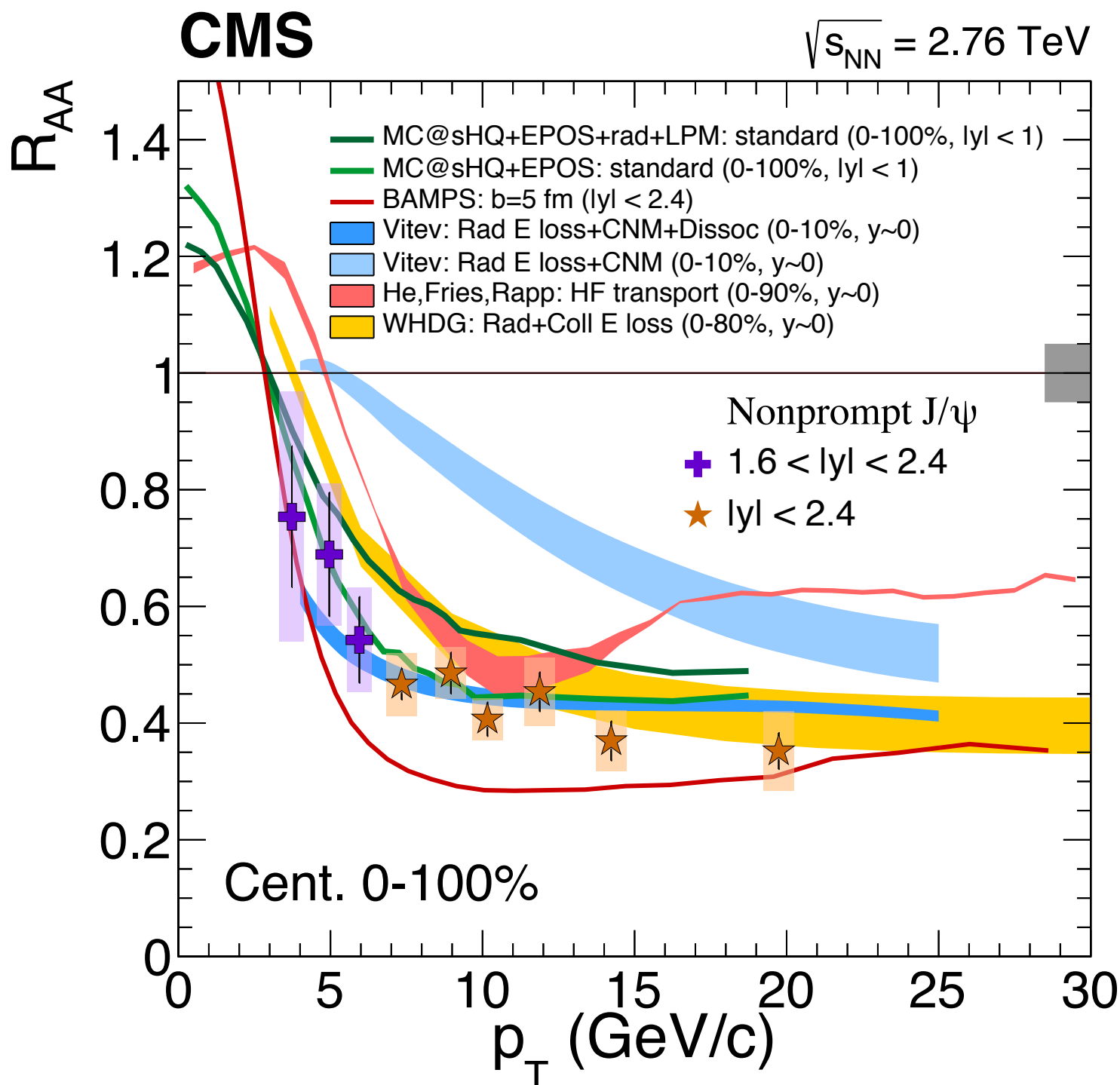
PbPb $351 \mu\text{b}^{-1}$, pp 28.0 pb^{-1} (5.02 TeV)



CMS results vs centrality, p_T and rapidity can help to constrain the model:

- Relative contribution of primordial and regenerated charmonia
- Dissociation and regeneration rates
- Temperatures at which J/ψ and $\psi(2S)$ regenerate
- ...

Comparison to theoretical calculations



CMS non prompt $1.6 < |y| < 2.4$
CMS non prompt $|y| < 2.4$

arXiv:1610.00613, Submitted to Eur. Phys. J. C

Strong suppression observed for non prompt J/ψ in PbPb collisions
Clear suppression as a function of p_T